

# M3D-C1 ZOOM Meeting

08/09/2021

## Announcements

### CS Issues

1. Intel-MPI on stellar
2. Mesh adaptation update
3. NERSC Time
4. Changes to github on b master since last meeting
5. Regression tests
6. Progress on optimizing the matrix assembling on GPUs – Chang Liu

### Physics Studies

1. Energy conservation with  $itemp=0,1$ ,  $ipres=0,1$  -- Lyons
2. Resistive wall mode in a periodic Cylinder – Strauss
3. Sawteeth with 50% runaways – Chen Zhao

**Note:** [meeting minutes posted on m3dc1.pppl.gov](https://m3dc1.pppl.gov)

## In attendance

Steve Jardin

Hank Strauss

Patrick Kim

Mark Shephard

Jin Chen

Adelle Wright

Nate Ferraro

Andreas Kleiner

Chen Zhao

Brendan Lyons

Chang Liu

Seegyoung Seol

Priyanjana Sinha

# Announcements

- Virtual Sherwood Meeting August 16-18
  - Registration until August 9 (Today)
  - No m3dc1 meeting August 16
- APS Nov 8-12
  - Meeting will be IN PERSON with virtual option
  - M3D-C1 Invited talks by C. Liu, A. Wingen
  - CTTS meeting? Will DOE allow travel?
- EPS 6/27 – 7/1 2022 in Maastricht, Netherlands
  - Nominate invited speakers by 29 October 2021

# stellar.princeton.edu

**From Adelle Wright 8/5/2021:**

I tested intelmpi on the stellarator version for a 3D fixed boundary case (24 planes, 6 nodes). There was no noticeable speedup but the memory usage was roughly halved.

While it's probably just that I did not configure everything correctly, I did have issues with post-processing (trace and IDL) when the modules needed for intelmpi were loaded.

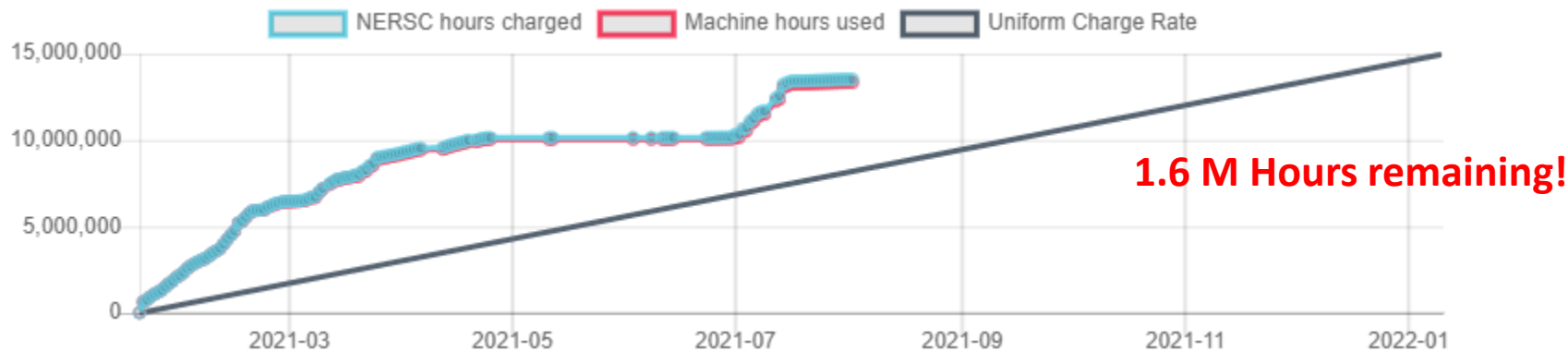
Update?

# Mesh Adaptation Update

RPI?

# NERSC Time

mp288



- mp288 received 10M Hrs for CY 2021
- Initial allocation exhausted by May 1
- John Mandrekas (DOE) added 5M Hrs additional
- More time may be possible if this is exhausted
- Pearlmutter time will not be charged for this FY
- We are NESAP Tier 2. Machine not yet ready. Phase-I w GPUs

# Changes to github master since 08/02/21

No Changes!

## Local Systems

- PPPL centos7(08/02/21)
  - 6 regression tests **PASSED** on centos7:
- PPPL greene (08/02/21)
  - 5 regression tests **PASSED on greene (m3dc1)**
- STELLAR (08/02/21)
  - 6 regression tests **PASSED** on stellar
- TRAVERSE(03/29/21)
  - Code compiles
  - Regression test failed: split\_smb not found in PATH



## Other Systems

- Cori-KNL (2/08/2021)
  - 6 regression tests passed on KNL
- Cori-Haswell (6/29/2021)
  - 6 regression tests passed
- PERSEUS
  - All 6 regression tests PASSED on perseus (J. Chen, 9/04/20)
- MARCONI
  - All regression tests PASSED on MARCONI (J. Chen, 9/04/20)
- CORI GPU (10/26)
  - ??

# Progress on optimizing the matrix assembling on GPUs

Chang Liu

# Discrepancies in Heat Flux Diagnostic

by  
**Brendan C. Lyons**

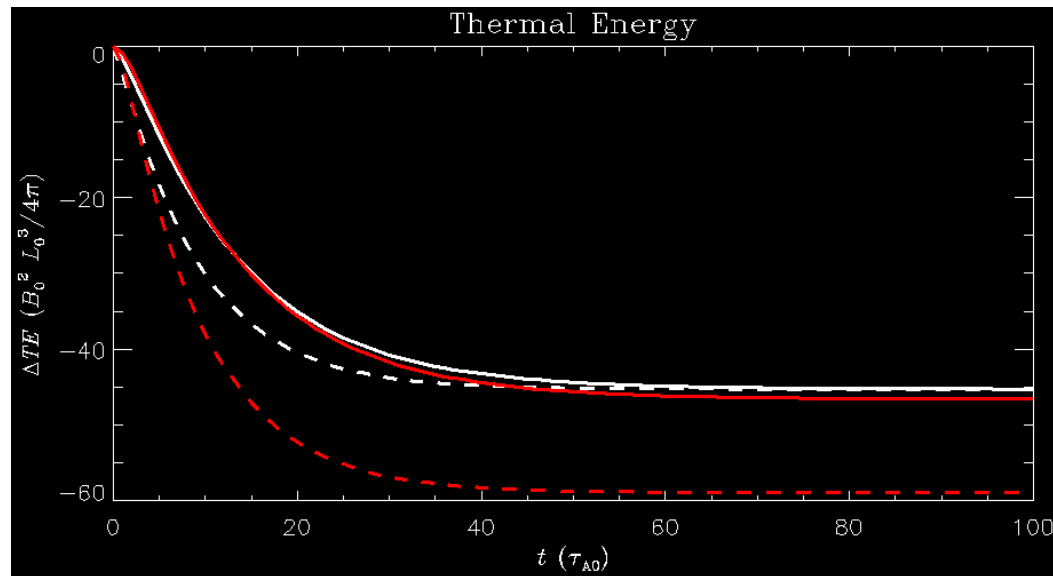
**August 2<sup>nd</sup>, 2021**

# General Comments

- **Based on a 2D ITER L-mode simulation**
- **Heat conduction should be the dominant source/sink**
  - No ohmic heating: `iohmic_heating=0`
  - No impurities
  - Kinetic energy  $\ll$  Thermal energy
- **In figures, white is the change thermal energy, red is the integrated `flux_thermal`**
- **All folders in `/pfs/nobackup/blyons/kappa_test/` on portal**
- **"Perpendicular" dominant case: `kappat=1e-1, kappar=1e-8`**
- **"Parallel" dominant case: `kappat=1e-1, kappar=1e+4`**
- **Boundary conditions**
  - `itemp=1` cases use `iconst_t=1` with `tebound=tibound=2e-5`
  - `itemp=0` cases use `iconst_p=1`, since `iconst_t=1` doesn't work for pressure equations
- **Lowering time step doesn't seem to make a qualitative difference in my tests**

# Two Temperature Equations: ipres=1, itemp=1

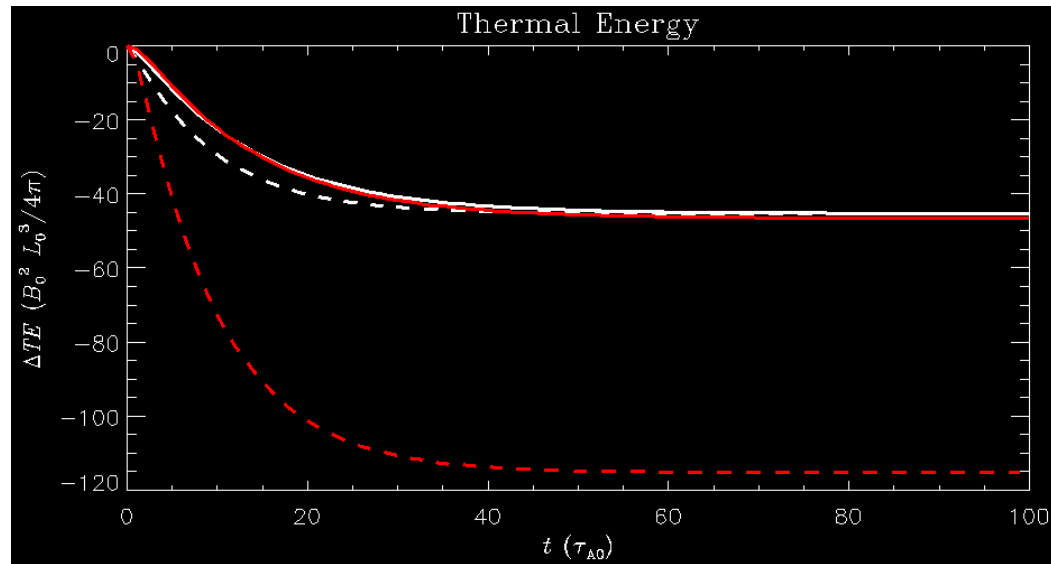
- **Perpendicular (solid)**
  - kth\_krl\_ip1\_it1/
  - Not terrible, but not good
- **Parallel (dashed)**
  - kth\_krh\_ip1\_it1/
  - Significant excess in the calculated heat flux
- **These discrepancies are what has made me look into this problem in the first place, since they mess up energy accounting**



For itemp=1, ipres=1, larger error for  $\kappa_{||}$  large

# Single Temperature Equation: ipres=0, itemp=1

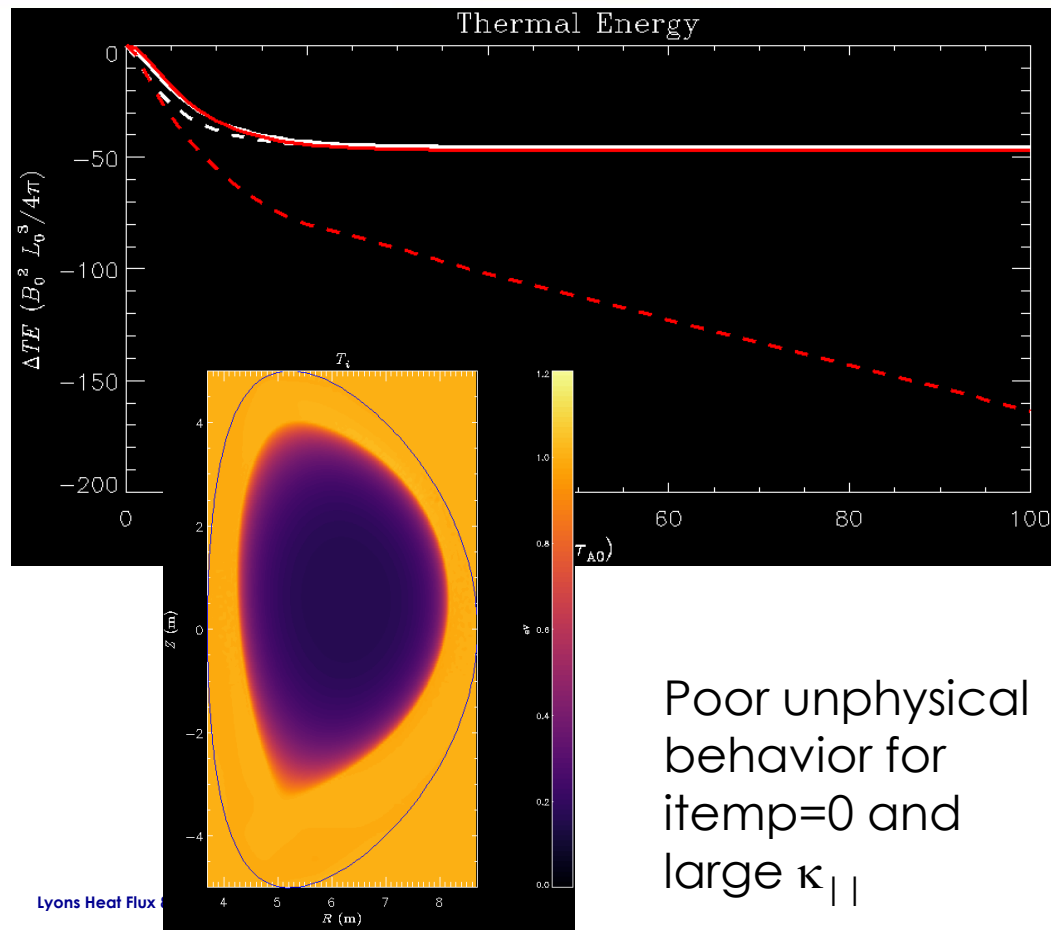
- **Perpendicular (solid)**
  - kth\_krl\_ip0\_it1/
  - Same as ipres=1
- **Parallel (dashed)**
  - kth\_krh\_ip0\_it1/
  - Roughly double the heat flux even though thermal decay is only very slightly different
  - Parallel heat flux appears to be double counted in flux\_heat()
  - Maybe because it accounts for the ion and electron temperature gradients?



For itemp=1  
ipres=0 gives larger  
error than ipres=1

# Single Pressure Equation: $i_{pres}=0$ , $i_{temp}=0$

- **Very strange behavior**
  - Temperature becomes lower in core than edge
  - Something wrong with heat flux implementation for open field lines?
- **Perpendicular (solid)**
  - $k_{th\_krl\_ip0\_it0/}$
  - No qualitative impact from temperature hole
- **Parallel (dashed)**
  - $k_{th\_krh\_ip0\_it0/}$
  - Shows persistent heat flux even after thermal energy is dissipated

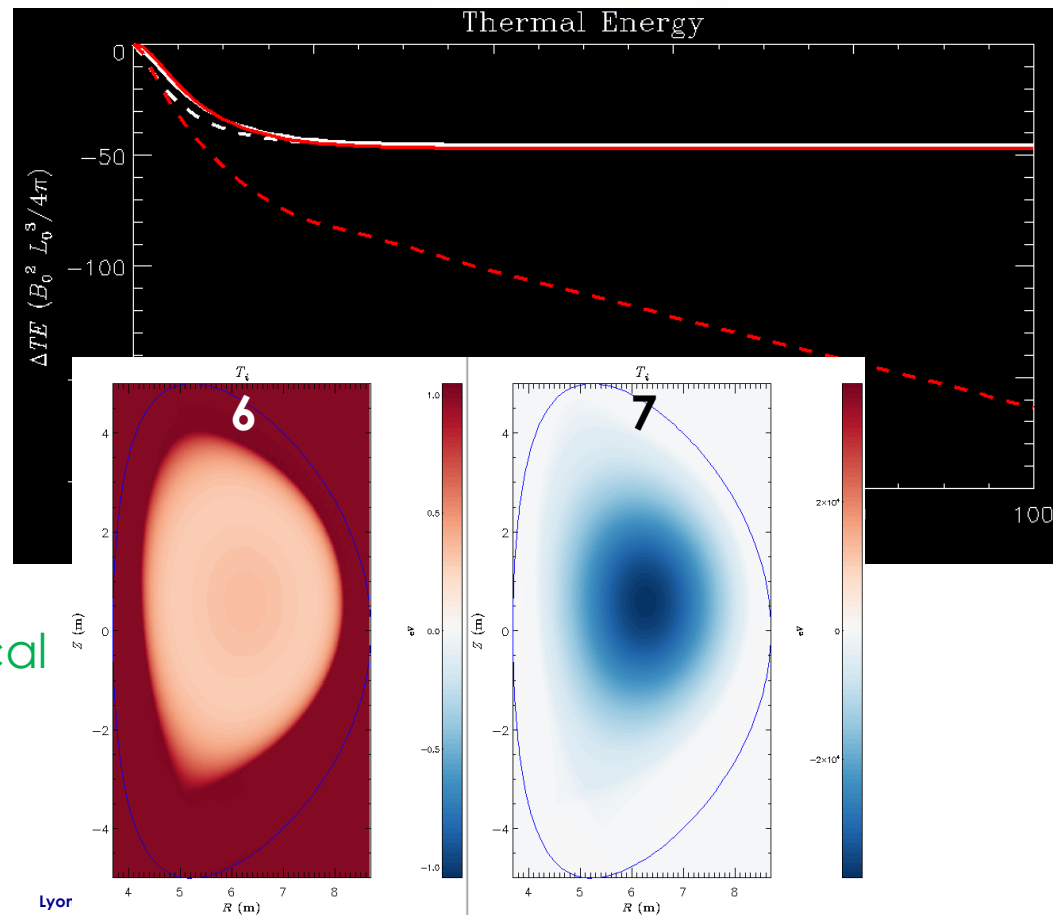


Poor unphysical behavior for  $i_{temp}=0$  and large  $\kappa_{||}$

# Two Pressure Equations: ipres=1, itemp=0

- **Even stranger behavior**
  - Ion temperature because huge and negative between outputs 6 and 7
  - No idea what to make of this
- **Traces are basically identical to ipres=0**
- **Perpendicular (solid):**  
kth\_kr1\_ip1\_it0/
- **Parallel (dashed):**  
kth\_krh\_ip1\_it0/

Poor unphysical behavior for itemp=0 and large  $\kappa_{||}$





# Conclusions

- Something is off with the `flux_heat()` diagnostic
- Nature of discrepancy is dependent on the `ipres` and `itemp`
- There may be something wrong about how the heat flux itself, not just the diagnostic, is implemented in the code for `itemp=0`

Do not use itemp=0 when kappar.ne.0

ltemp = 0

$$\begin{aligned}\frac{\partial p}{\partial t} &= \dots + \nabla \cdot \kappa_{\parallel} \frac{\mathbf{B}\mathbf{B}}{B^2} \cdot \nabla \left( \frac{p}{n} \right) \\ &= \dots + \nabla \cdot \kappa_{\parallel} \frac{\mathbf{B}\mathbf{B}}{B^2} \cdot \left( \frac{1}{n} \nabla p - \frac{p}{n^2} \nabla n \right)\end{aligned}$$

↗  
This term is a problem since p and n are solved separately. Will cause very large errors and possibly numerical instability

ltemp = 1

$$n \frac{\partial T}{\partial t} = \dots + \nabla \cdot \kappa_{\parallel} \frac{\mathbf{B}\mathbf{B}}{B^2} \cdot \nabla (T) \quad \text{Well behaved, diagonally dominant}$$

function flux\_heat

$$[\kappa_e \hat{n} \cdot \nabla T_e + \kappa_i \hat{n} \cdot \nabla T_i] \quad \text{HF\_perp}$$

$$-\frac{1}{B^2} [\hat{n} \cdot \nabla \psi \times \nabla \phi] \left[ \nabla \psi \times \nabla \phi \cdot (\kappa_{\parallel} \nabla T_e + \kappa_{\parallel} \nabla T_i) \right] \quad \text{HF\_par}$$

	itemp=1 ipres=1		temp=1 ipres=0	
	A		B	
kappar	HF_perp	HF_par	HF_perp	HF_par
1.	46.1	0.46	46.3	0.46
10	42.5	4.2	44.5	4.42
100	26	23	33.0	30.2
1000	06	48.5	10.6	84.2
10000	00.25	58	00.77	114.

For both A @ B: ion par HF = electron par HF

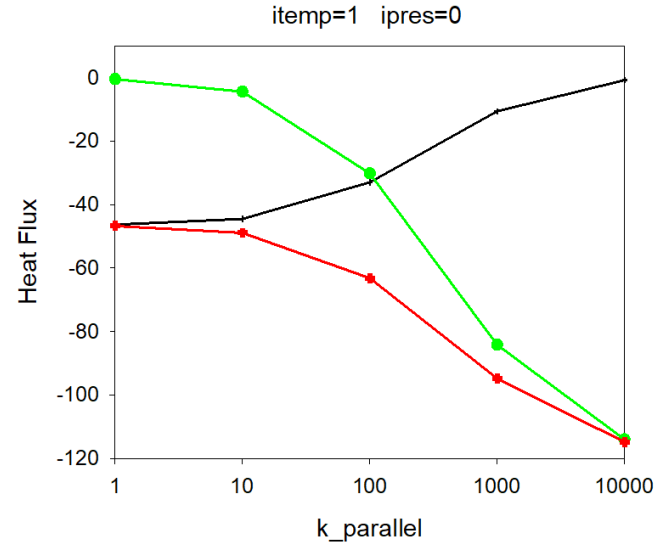
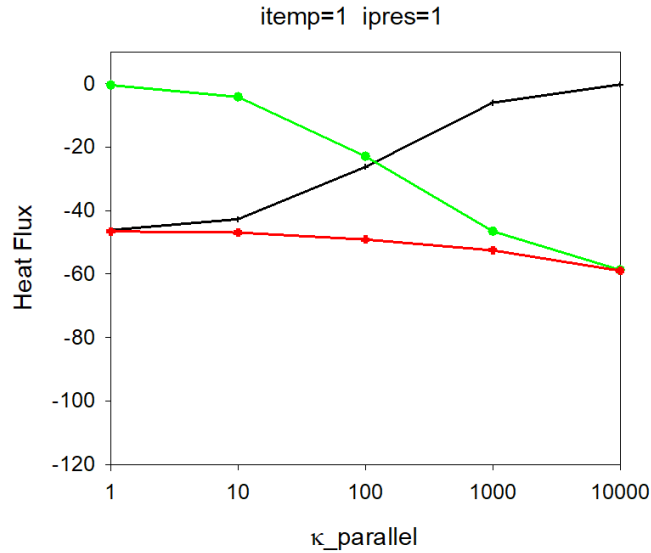
For both ions and elects : B\_par\_HF = 2 x A\_par\_HF

Te = Ti and are almost the same for both cases A and B

Graphs look much better if we multiply the HF\_par in case B by ½

Problem seems to be with the parallel heat flux with ipres=0

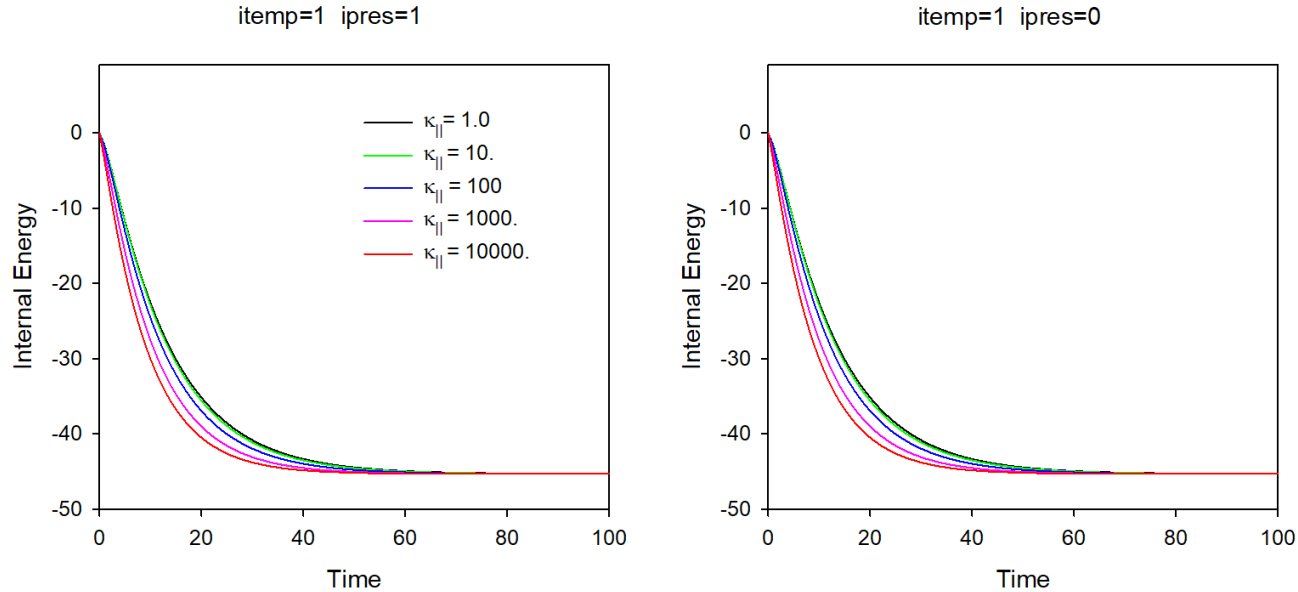
## Parallel heat flux is much larger for ipres=0



Note:

- Red line (total heat lost) should be horizontal if energy conserved for all  $\kappa_{\parallel}$
- Perpendicular heat loss almost the same for ipres=1 and ipres=0
- Parallel heat loss at large  $\kappa_{\parallel}$  slightly too large for ipres=1, much too large for ipres=0

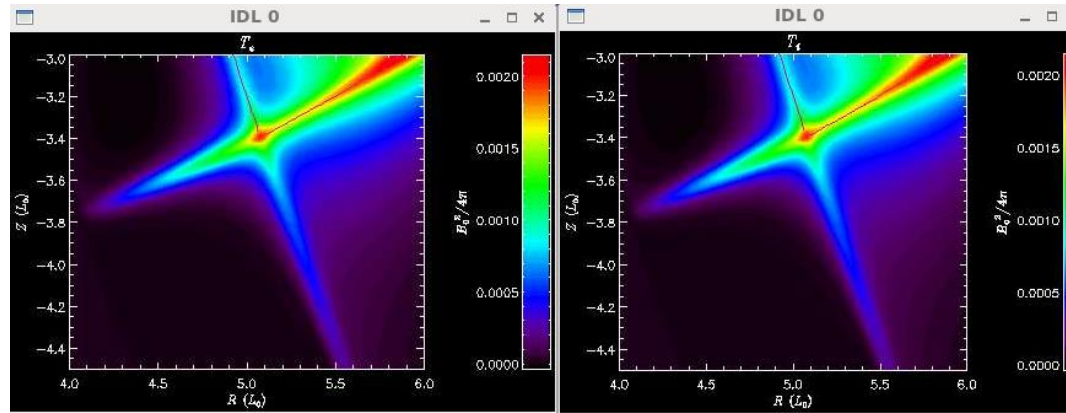
## Decay of internal energy similar for ipres=1,0



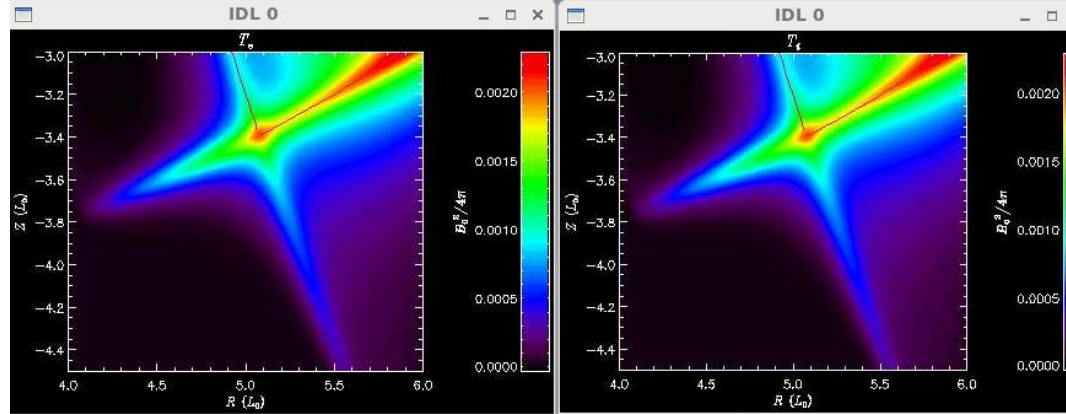
Only weakly dependent on  $\kappa_{||}$  for both cases, and very similar

dTe,i/dy profiles look nearly the same in SOL for itemp=1, ipres=0,1

ipres=1



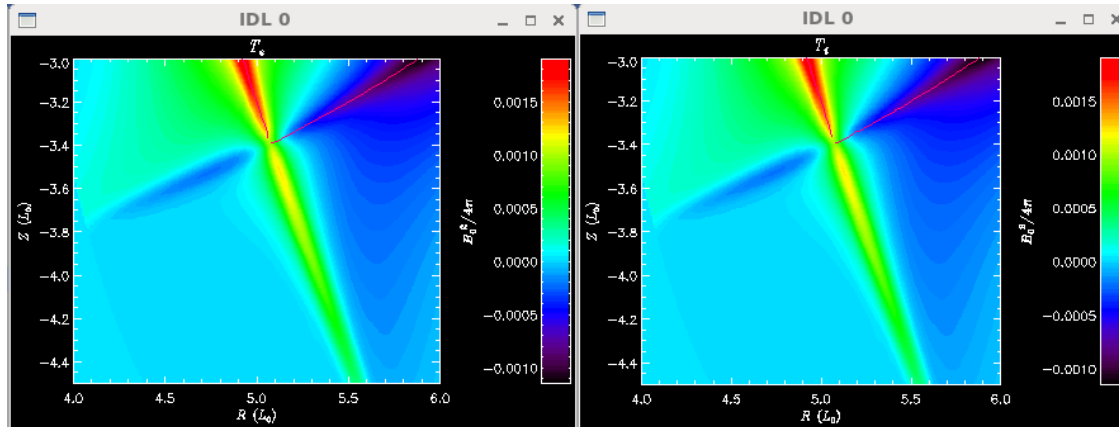
ipres=0



Time slice 2  
 $\kappa_{||} = 10^4$

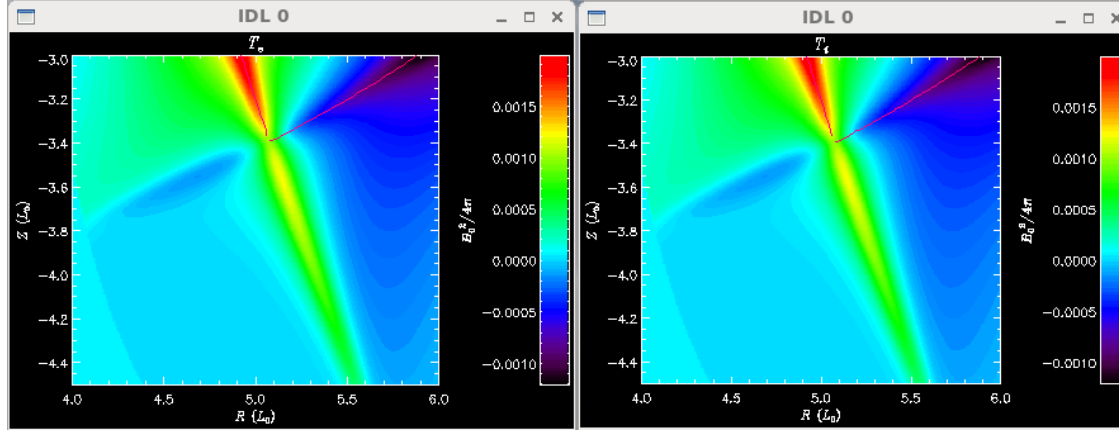
dTe,i/dx profiles look nearly the same in SOL for itemp=1, ipres=0,1

ipres=1

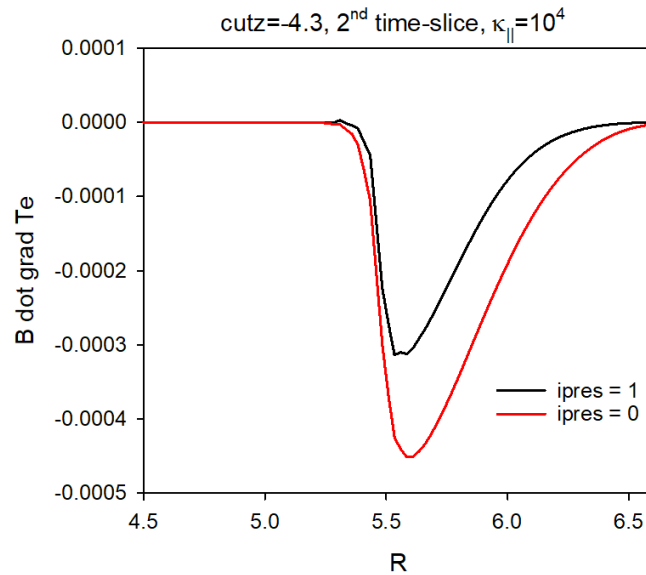
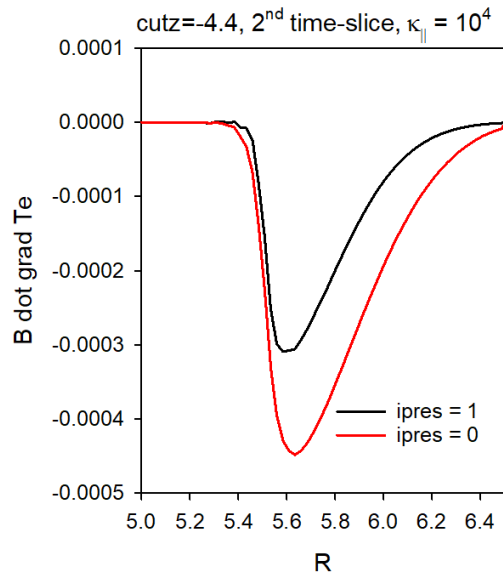


Time slice 2  
 $\kappa_{||} = 10^4$

ipres=0



$\mathbf{B} \cdot \nabla T_e$  along  $z=-4.4, 4.3$  much larger for  $\text{ipres}=0$





## Perpendicular temperature equation for ipres=0

$$n_e \dot{T}_e = \nabla \cdot \kappa \nabla T_e \quad (1)$$

$$n_i \dot{T}_i = \nabla \cdot \kappa (\text{kappai\_fac}) \nabla T_i$$

$$T_i = \frac{(1-\text{pefac})}{\text{pefac}} T_e$$

$$n_i \frac{(1-\text{pefac})}{\text{pefac}} \dot{T}_e = \nabla \cdot \kappa (\text{kappai\_fac}) \frac{(1-\text{pefac})}{\text{pefac}} \nabla T_e \quad (2)$$

Equation for  $T_e$  that conserves energy is obtained by adding (1)+(2)

$$\left[ n_e + n_i \frac{(1-\text{pefac})}{\text{pefac}} \right] \dot{T}_e = \nabla \cdot \left[ 1 + (\text{kappai\_fac}) \frac{(1-\text{pefac})}{\text{pefac}} \right] \kappa \nabla T_e$$

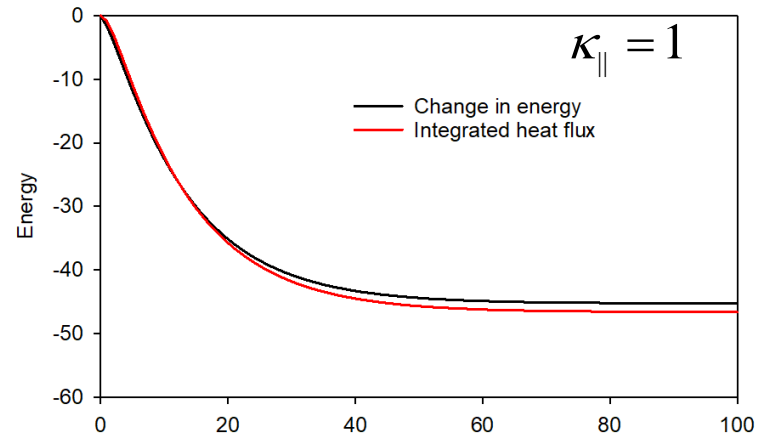
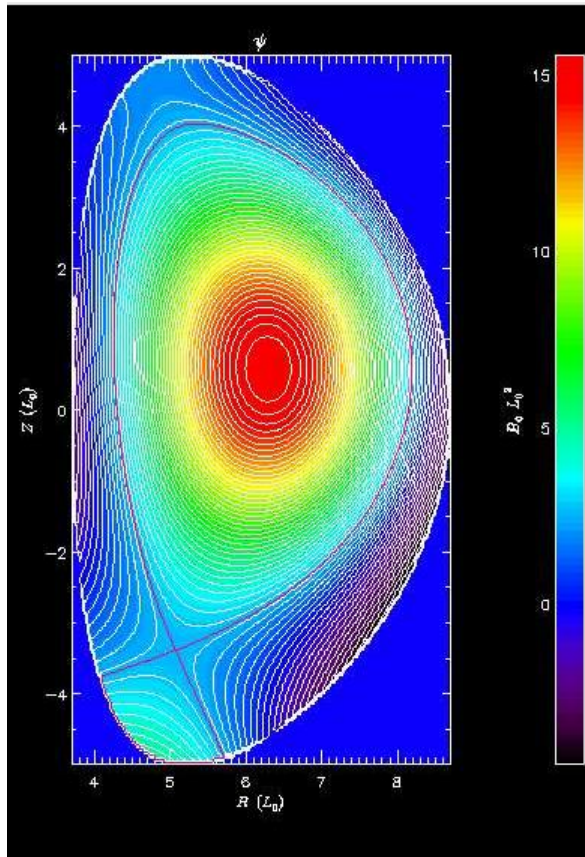
**Same needs to be done for parallel heat conduction for ipres=0!**

for ipres=0, replace (in temperature\_lin)

$$\text{kappar} \rightarrow \left[ 1 + \text{kappai\_fac} \frac{(1-\text{pefac})}{\text{pefac}} \right] \text{kappar}$$

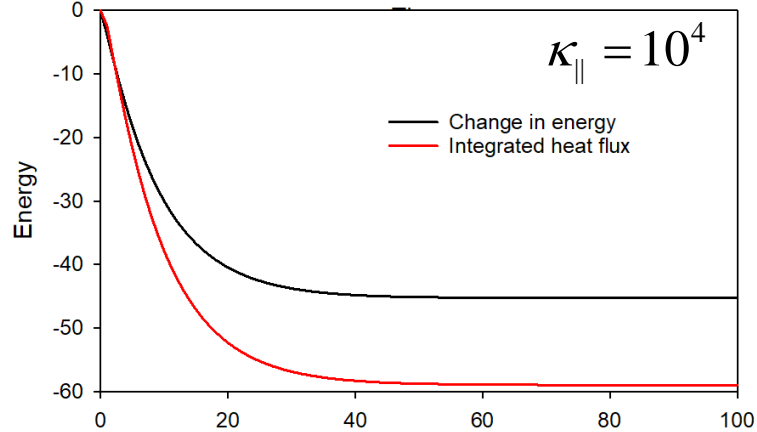
← This needs to be changed

## Parallel heat flux term requires high resolution



Now, exactly the same results for  $\text{ipres}=1$  and  $\text{ipres}=0$  for  $\text{itemp}=1$ .

Much better energy conservation at small values of  $\kappa_{||}$



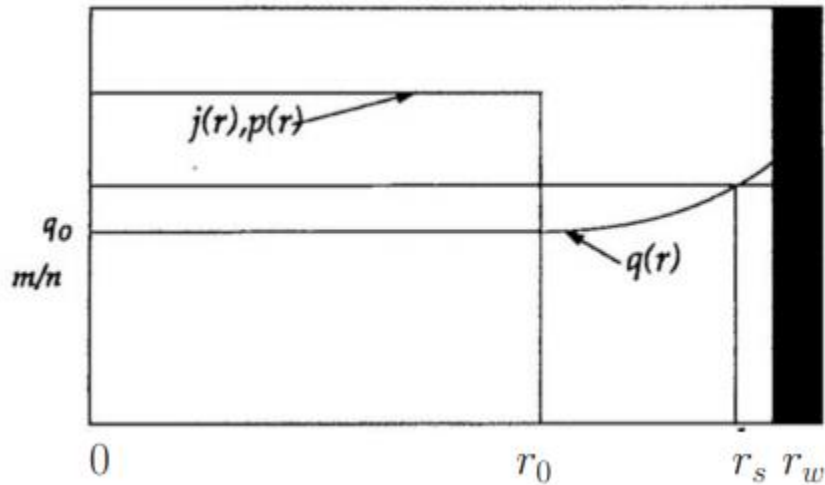
Large  $\kappa_{||}$  energy conservation may improve if mesh is refined

# Problem Solved !

1. Do not use itemp=0 if kappar .ne. 0
2. To make itemp=1, ipres=1 and itemp=1, ipres=0 identical  
for ipres=0, replace (in temperature\_lin)  
$$\text{kappar} \rightarrow \left[ 1 + \text{kappai\_fac} \frac{(1 - \text{pefac})}{\text{pefac}} \right] \text{kappar}$$
3. To get more precise energy balance for large kappar, may need to increase spacial resolution in SOL:
4. I have not yet committed the change as it makes one of the regression tests fail (RMP\_nonlin)

# Resistive Wall Mode in Periodic Cylinder

H. Strauss, 8/9/21



Plasma resistivity?  
What is  $f$ ?

RWM

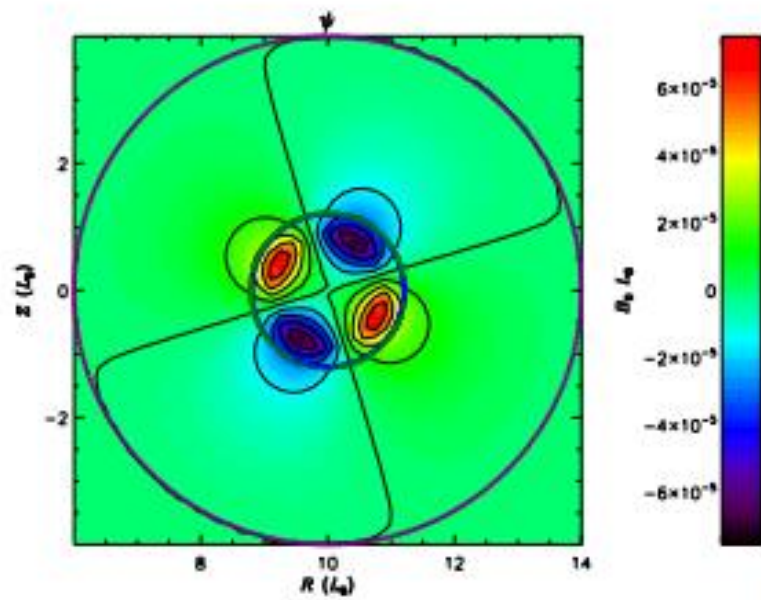
$$\gamma\tau_{wall} = -m \frac{1 - (m - nq)}{1 - (m - nq) - (r_0/r_w)^{2m}}$$

RWTM

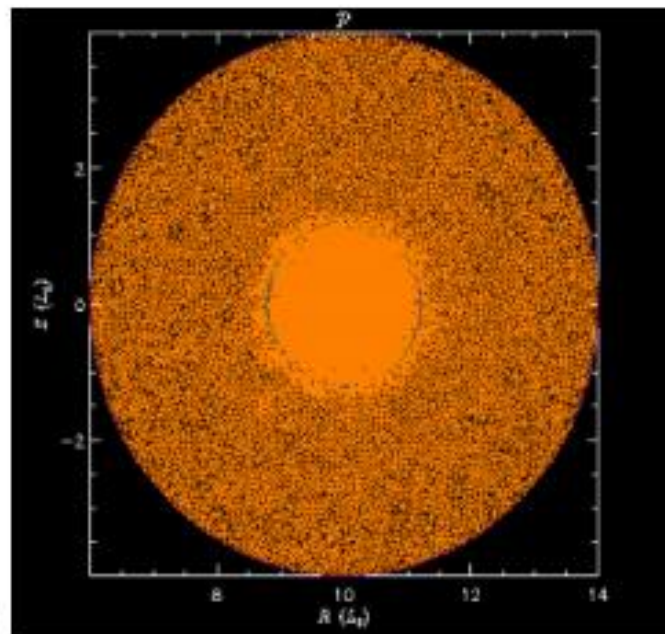
$$\gamma\tau_A = \frac{c_0}{S^{1/3} S_{wall}^{4/9}}$$

$$c_0 = 2.46 \left( \frac{q' r_s}{q} \right)^{2/9} f^{4/9} = 2.46 f^{4/9}$$

$$f = \frac{(r_s/r_w)^{2m}}{[1 - (r_s/r_w)^{2m}]^2}$$



(a)



(b)

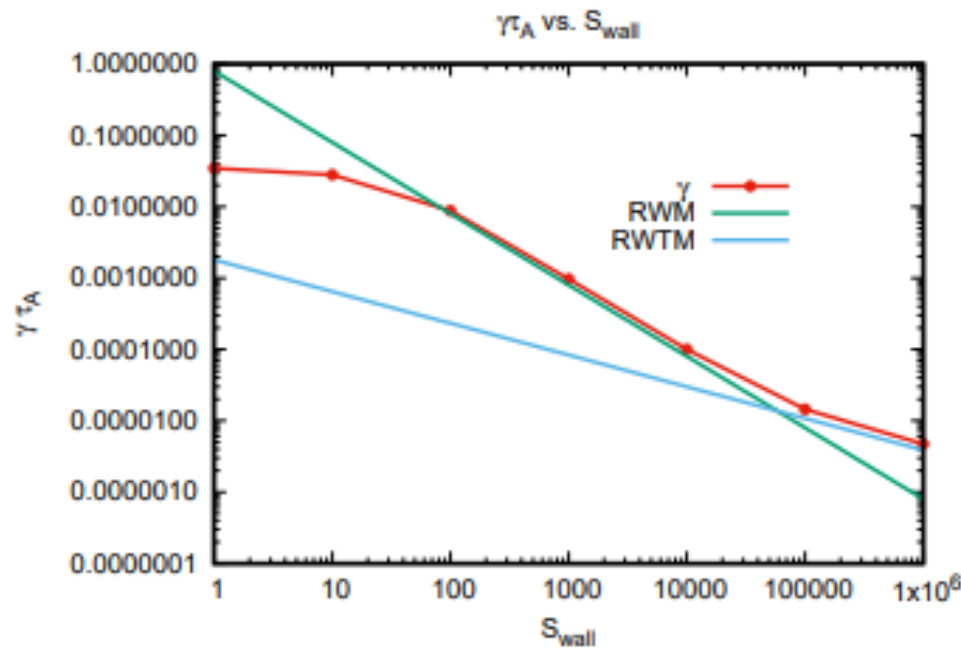
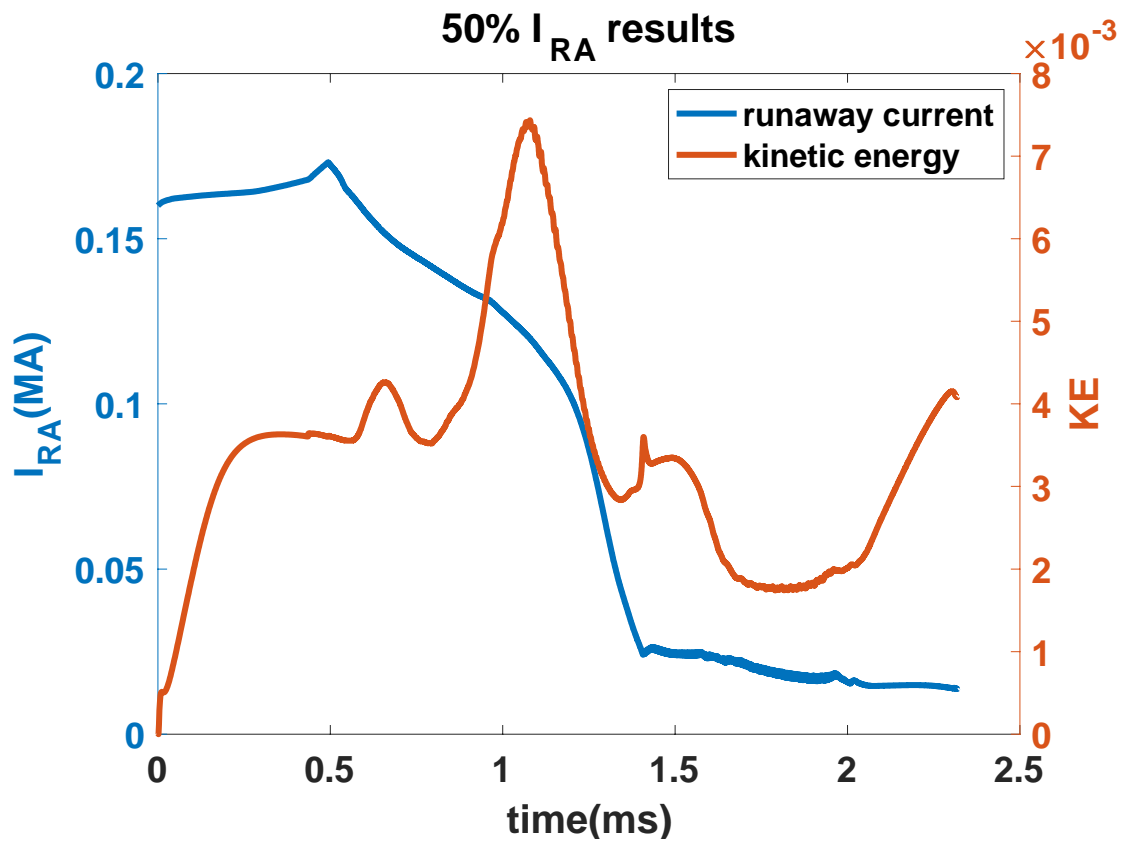


Fig.2(a) shows linear  $\psi$  for the case  $S_{wall} = 10^4$ . Fig.2(b) shows the adapted mesh used in all the cases. The mesh has a thin wall,  $\delta = 0.02$ . Fig.3 shows growth rate  $\gamma$  as a function of  $S_{wall}$ . For  $S_{wall} \geq 10^5$ , the most unstable mode appears to be a RWTM. The straight line fits are to  $S_{wall}^{-1}$  for a RWM and  $S_{wall}^{-4/9}$  for a RWTM.

Dependence on  $\eta$  ?  
Does mode structure change?

Sawteeth with 50% runaway  
current

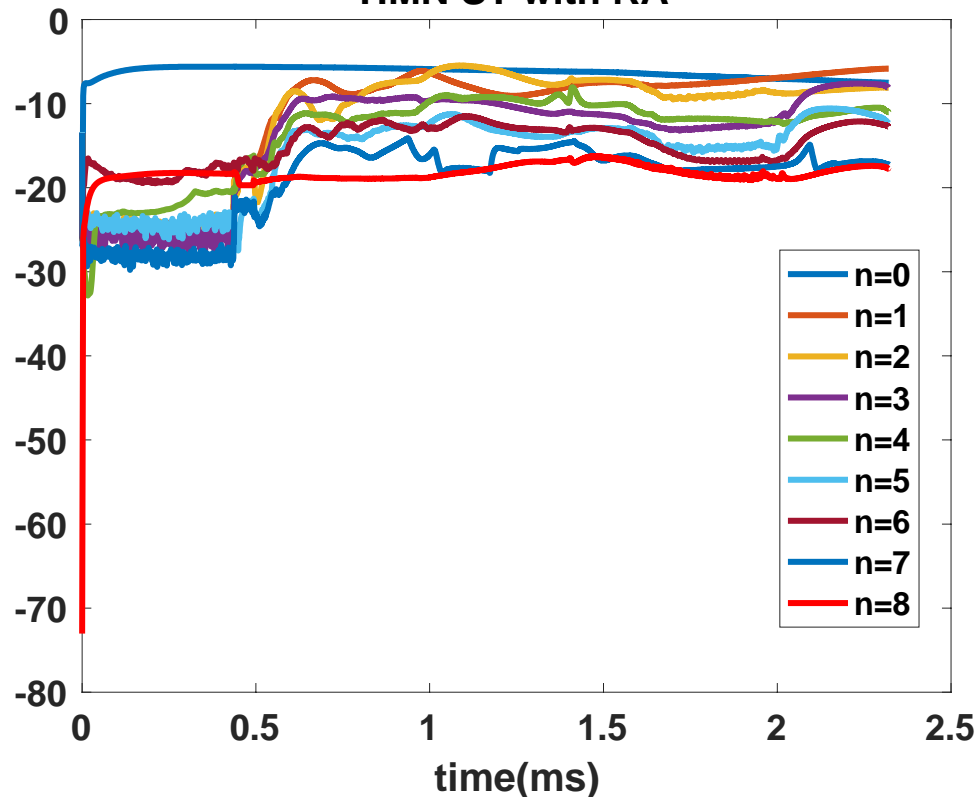
08092021



- The runaway electron current dropped to nearly 0 at about  $t=1.1$ ms by the MHD instability during ST.

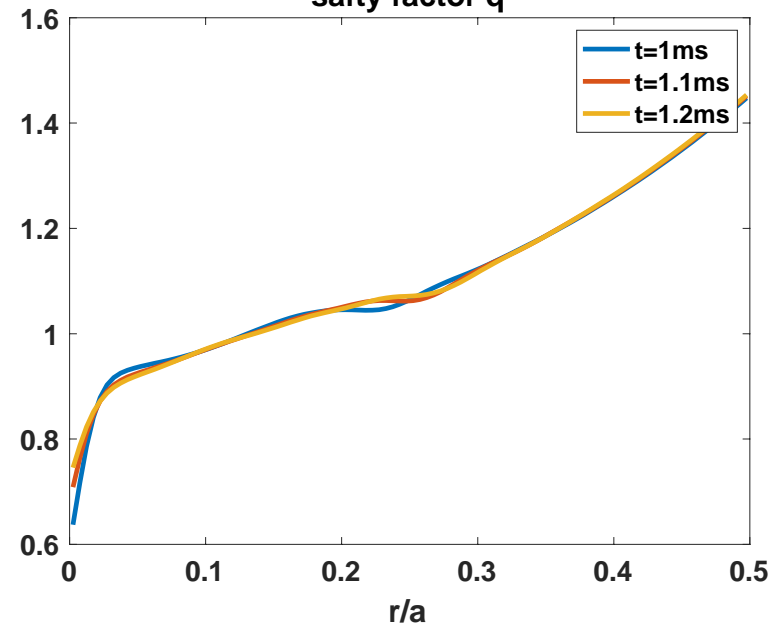


# HMN ST with RA

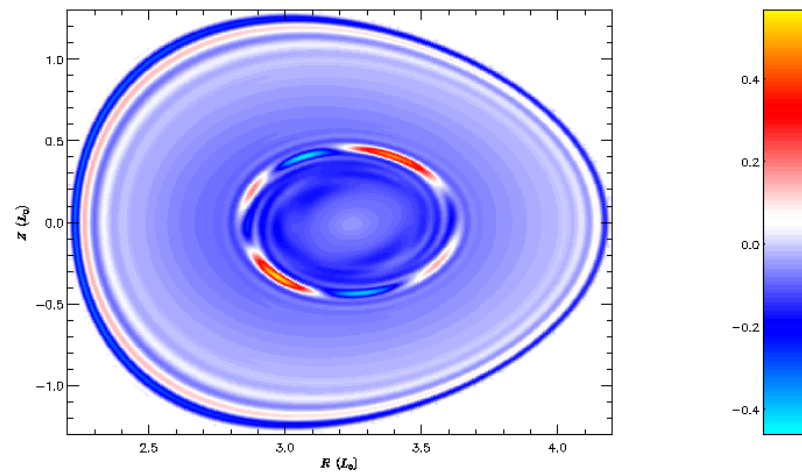


The highest is  $n=2$  at  $t=1.1\text{ms}$

# saftey factor q

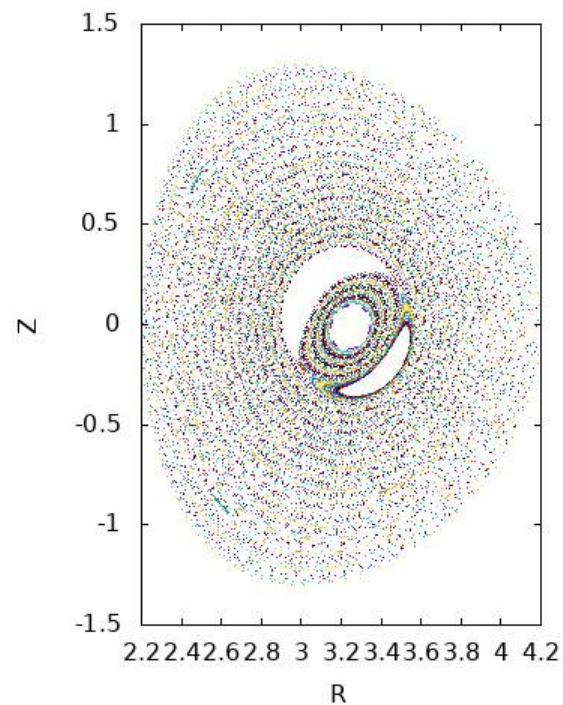


Runaway current density profile at  $t=1.1\text{ms}$



Maybe  $n=2$   $m=3$  mode

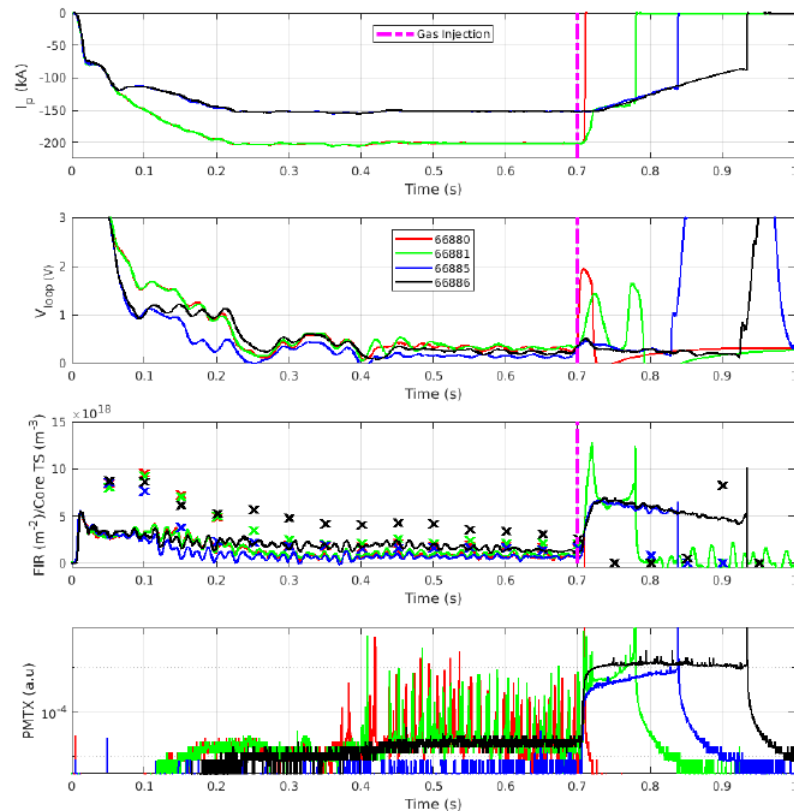
Poincare plot at  $t=1.1\text{ms}$



# Expulsion of RE Seed - Sawteeth

From Umar Sheikh  
(7/22/21)

- Sawteeth expel Res
  - Seed population lost - HXR spikes
  - No RE beam after MGI
- Lower  $I_p$  removed Sawteeth
  - HXR spikes gone
- Stable RE beams produced
- Sharp drop in  $I_p$  indicates larger Ohmic current contribution
  - Sharp drop in  $I_p$  if fewer REs (green)
  - Significant RE current before MGI



That's All I have

Anything Else ?