# The M3D-C1 User's Guide

# M3D-C1 Team RPI SCOREC

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# 1 Obtaining M3D-C1

#### 1.1 License Agreement

M3D-C1 is a code developed with funding from the US Department of Energy, and is intended for open scientific research. If you intend to run M3D-C1, please complete the following steps:

- 1. Sign the license agreement https://m3dc1.pppl.gov/M3D-C1\_License.pdf and return to nferraro@pppl.gov.
- 2. Request access to the code using the form https://pppl.tiny.us/code-release-form. This step will involve review by PPPL to clear any potential export control issues.

Please be aware that permission to run M3D-C1 does not carry an implicit agreement to provide technical support for compiling, running, modifying, or interpreting output of, M3D-C1.

#### 1.2 Preinstalled M3D-C1 Executables

For the general user, we recommend using precompiled executables and associated modules for release versions, where available. Installations are presently available on a number of systems, including:

```
GA Iris:
    module use /fusion/projects/codes/m3dc1/modules
    module load m3dc1/1.14
NERSC Perlmutter:
    module use /global/cfs/projectdirs/mp288/C1/modules/perlmutter
    module load m3dc1/1.14-cpu
PPPL Cluster:
    module use /p/m3dc1/modules
    module load m3dc1/1.14
PU Stellar:
    module use /projects/M3DC1/modules
    module load m3dc1/1.14
```

These modules will modify the user's environment variables appropriately to access the M3D-C1 executables, python libraries, and IDL routines.

#### 1.3 Accessing the M3D-C1 Source Code

If you choose to build the code yourself, either to use an unreleased version or to do code development yourself, you will need access to the M3D-C1 code repository. The M3D-C1 source code is located in the Github repository: PrincetonUniversity/M3DC1. To get access to this repository, please complete the license form and software access request form as described in section 1.1.

#### 1.4 Makefiles and Dependencies

Some of the build scripts depend on the following environment variables to be set, to specify the location of the M3D-C1 source code and the directory in which to install any compiled executables:

- M3DC1\_DIR should be set to the directory containing the M3D-C1 source code. For example: setenv M3DC1\_DIR \$HOME/src/M3DC1
- M3DC1\_INSTALL\_DIR should be set to the directory in which M3D-C1 will be installed. For example: setenv M3DC1\_INSTALL\_DIR \$HOME/M3DC1

It is recommended to set these values in your login script.

Makefiles for a number of systems are included in the repository, with filenames \$M3DC1\_DIR/unstructured/\*.mk. For most of these systems, environment modules are also included in the repository. These modules will load the appropriate software modules for building on that particular system, and can be loaded using:

NERSC Perlmutter: module use \$M3DC1\_DIR/modules/perlmutter module load m3dc1/devel

PPPL Cluster:

module use \$M3DC1\_DIR/modules/pppl
module load m3dc1/devel-centos7

PU Stellar:

module use \$M3DC1\_DIR/modules/stellar
module load m3dc1/devel

It is recommended to place the appropriate module use statement in your login script. The makefile should have the name  $\{M3DC1\_ARCH\}.mk$ . If  $M3DC1\_ARCH$  is not defined, it will default to HOST, stripped of any trailing numbers identifying a node index on multinode systems (*e.g.* if HOST==``sunfire06'' then  $M3DC1\_ARCH$  will default to "sunfire").

If you are building M3D-C1 on a system for which no makefile or development module is provided, please refer to the existing makefiles and modules as examples. In general, M3D-C1 requires the following dependencies:

- Compilers for C, C++, and Fortran;
- MPI

- HDF5 compiled with support for Fortran and MPI
- netcdf
- GSL
- FFTW
- PETSc compiled with support for Fortran, complex-valued functions, MUMPS and/or SuperLU\_dist
- PUMI meshing libraries

#### 1.5 Building

Once the appropriate makefile has been defined, the M3D-C1 executables can be built by entering <code>\$M3DC1\_DIR/unstructured</code> and running

make all

This will run the following commands:

make	OPT=1	Builds the 2D version
make	OPT=1 COM=1	Builds the complex version
make	OPT=1 3D=1 MAX_PTS=60	Builds the 3D version
make	OPT=1 3D=1 MAX_PTS=60 ST=1	Builds the stellar version
make	a2cc	Builds a utility for extracting coil currents from a-eqdsk files
make	bin	Puts executables into the _\$M3DC1_ARCH subdirectory

#### **1.6** Regression Tests

The M3D-C1 source code includes a suite of regression tests that should be run before commiting any new code to the repository. To run these tests:

#### cd \$M3DC1\_DIR/unstructured/regtest ./run <arch> <test>

where **<arch>** and **<test>** are optional arguments specifying the specific batch script to run, and the specific regression test to run, respectively. By default, **<arch>=\$M3DC1\_ARCH**.

If <test> is not specified, then all the regression tests will be run, using the batch scripts \$M3DC1\_DIR/unstructured/regtest/\*/base/batchjob.<arch>.

The run script will create new directories in which to run these regression tests, named \$M3DC1\_DIR/unstructured/regtest/\*/\${M3DC1\_VERSION}\_<arch>/.

If <test> is specified, then only \$M3DC1\_ARCH/unstructured/regtest/<test>/base/batchjob.<arch> will be run (again, in a new directory named as described above). If <test> is specified, then <arch> must also be specified.

To check the results of the regression tests

./check <arch>

# 2 GITHUB

Retrieve the current version of M3D-C1 from the GIT repository. For the first time, to check out the sources, do:

Initial access is with the *clone* command. This copies the source code from the master file into a working directory on your machine. You only do this once on each computer you work on.

 $module \ load \ git$ 

git clone https://github.com/PrincetonUniversity/M3DC1

Subsequent GIT commands used to commit:

- *add/commit/push*: You *add* files to a list of files to update, *commit* the chanes to your branch, and then *push* the changes to the master branch.
- git commit -m "message describing changes" (adding -a commits all changes)
- *diff* lists the changes you made from the last commit, even if you haven'tpushed your commits to github. To see how your files differ from what's on github, you can do:

git fetch origin master git diff origin/master

- *status* compares your branch with the master branch
- *pull origin master* updates your local branch to the current master branch
- $\bullet\ stash$  takes uncommitted changes, saves them for later use, and reverts files in working directory

stash list stash drop stash apply stash pop (apply and drop)

- stash pop removes changes from your stash and reapplies them to working copy
- stash apply keeps changes in stash, but reapplies them to working copy
- reset -hard discards any changes to local branch since last commit
- *branch* tells you what branch you are in
- log (–one line) (-after 2023-01-31) lists all the commits for the checked-out branch after that date
- *checkout hashtag* replaces your version with the version that has hastag "hashtag" (need only first 7 characters)

#### 2.1 Branches in GIT

To make a new branch called fp-phase2:

- git checkout master switch to the master branch
- git pull make sure the master branch is up to date
- git checkout -b fp\_phase 2 The =b creates a new branch. This will be identical tomaster to start.
- git push –set-upstream origin fp\_phase2 push this new branch to the remote so others can access it

#### 2.2 Committing changes

For example, to commit changes to newpar.f90

- git pull Always do thi before you start committing
- git add newpar.f90 This stages the current changes in newpar.f90 for commit. You could then make more change before committing, but you would have to add again to get those into the commit. Or, you could add -a to the following to commit all changes.
- git commit -m "Changed newpar.f90" Commit changes to your local branch.
- *git push* Push commits to the remote repo. –set-upstream only needs to be done the first time.

#### 2.3 Merging branches

To merge changes on master into fp\_phase2

- git checkout master switch to the master branch
- *git pull* get the latest commits on the master branch
- git checkout fp\_phase2 switch back to fp\_phase2
- *git pull* to get the latest commits to fp\_phase2
- git merge master merge any new commits into fp\_phase2. This makes a commit. You may need to resolve conflicts.
- git push Push the merged commit to remote repo

Inverting fb\_phase2 and master here would merge the development branch into master locally, then the push would send the merge to the remote master

# 3 Mesh Generation and Management

Given a mesh, all the mesh support needed to run  $M3D-C^1$  is provided by PUMI (Parallel Unstructured Mesh Infrastructure) developed at RPI Scientific Computation Research Center (SCOREC). The PUMI is used to manage the mesh information as it is processed within the  $M3D-C^1$  code. PUMI is a freely available open source software. For more information on how to get and install, visit http://www.scorec.rpi.edu/pumi.

The generation of M3D-C<sup>1</sup> meshes for Tokamak geometries using the M3D-C<sup>1</sup> mesh generation program involves both PUMI and Simmetrix mesh generation software. Simmetrix is commercial software which provides a set of tools and libraries for engineering simulation. As Simmetrix runs only with a valid license, any M3D-C<sup>1</sup> user who wants to install his/her own version of the M3D-C<sup>1</sup> mesh generation program should contact Simmetrix to purchase its license. For more information on Simmetrix, visit http://simmetrix.com. M3D-1<sup>1</sup> users with accounts on PPPL Portal or Princeton Stellar are allowed to run M3D-C<sup>1</sup> mesh mesh generation programs on those systems.

Note that Simmetrix software is not required to run a  $M3D-C^1$  simulation. Running a  $M3D-C^1$  simulation using a mesh generated by a different mesh generator requires the development of a tool that creates PUMI readable model and mesh files with  $M3D-C^1$  related control information required as the  $M3D-C^1$  input. For more information, contact RPI SCOREC (shephard@rpi.edu).

With respect to the mesh support, there are a number of files involved with housing the geometry and mesh information. First, the model file extensions are the following

- .smd: Simmetrix-readable binary format model file The model generated with Simmetrix is saved in this format.
- .dmg: PUMI-readable binary format model file The model generated from PUMI mesh
- .txt: M3D-C<sup>1</sup>-readable ascii format model file The model is generated from mesh generation tool (See Section 3.1)

Second, the mesh file extensions are the following.

- .sms: Simmetrix-readable binary format mesh file
  - The mesh generated with Simmetrix is saved in this format.
  - If a mesh is serial (1-part), the mesh file doesn't have a number before the extension
  - If a mesh is distributed (P-part, P>1), the mesh file has a number before the extension to represent the global part ID.
- .smb: PUMI-readable binary format mesh file
  - This format is used in M3D-C<sup>1</sup> to import/export a mesh
  - No matter if a mesh is serial (1-part) or distributed (P-part, P>1), the mesh file has a number before the extension to represent the global part ID.

• .vtu/pvtu: binary format mesh file for visualization with Paraview. For more information, visit http://paraview.org.

An overview of the Model/Mesh requirements for the  $M3D-C^1$  mesh generation process are as follows:

- The model and mesh shall be generated as described in Section 3.1.
- The mesh file must be PUMI-readable .smb file. Note that a mesh file name contains a number before the extension (.smb) to denote a global part ID.
- The model and mesh file must be present in the work directory
- The name of model and mesh file must be specified in Clinput file in the work directory
  - mesh\_model = model\_file
  - mesh\_filename = mesh\_file.smb (NOTE: do not specify a number before the file extension)
- In a 2D run with P processes, there should be P mesh files with part ID from 0 to P-1
- In a 3D run with P×N processes where 2D mesh is distributed to P parts,
  - there should be P mesh files with part ID from 0 to P-1
  - in Clinput file, specify nplanes to N (e.g. nplanes=8), where nplanes describes how many 2D mesh copies to be loaded
  - the M3D- $C^1$  code should be compiled with options "3D=1, MAX\_PTS=60".

The rest of this section is organized as follows: Section 3.1 describes a mesh generation program m3dc1\_meshgen. Section 3.2 describes a mesh generation program m3dc1\_mfmgen. Section 3.3 describes a mesh generation program polar\_meshgen. Section 3.5 presents a mesh partitioning program "split\_smb" and "collapse" which changes the number of parts of the mesh. For how to visualize a mesh with Paraview, see Appendix B.

#### 3.1 m3dc1\_meshgen

m3dc1\_meshgen requires an ascii input file of arbitrary name that contains the following parameters.

- modelType: 0, 1, 2, 3, or 4
  - Type 0: a parameterized vacuum region defined by five doubles for analytic expression. For five doubles  $X_0 X_1 X_2 Z_0 Z_1$ , vacuum boundary is defined by

$$X = X_0 + X_1 \cos(\theta + X_2 * \sin(\theta)) \tag{1}$$

$$Z = Z_0 + Z_1 \sin(\theta) \tag{2}$$

- Type 1: a vacuum region defined by piece-wise linear points
- Type 2: a vacuum region defined by piece-wise polynomials
- Type 3: spline-fitted 3-region model (plasma, wall and vacuum)
- Type 4: spline-fitted 3-region model (plasma, wall, and vacuum) with inner & outer boundary points to set resistive wall
- reorder: if 1, reorder PUMI mesh based on adjacency (default: 0) and generate vtk folders for mesh visualization. The mesh before and after reodering is saved in original-mesh.vtk and reordered-mesh.vtk, respectively. Note that the element order of Simmetrix mesh is not affected.
- inFile: (modelType 0) not required (modelType 1 and 2) geometry file describing the vacuum (modelType 3 and 4) geometry file describing the inner plasma wall
- bdryFile: (modelType 0-3) not required (modelType 4) geometry file describing the outer plasma wall
- outFile: output file name to save model and mesh
- meshSize: relative mesh size for each region (default 0.05) for modelType 3, set three doubles for plasma, resistive, vacuum, respectively
- useVacuumParams: for modelType 0 or 3, if 1, use parameterized vacuum wall (default 0)
- vacuumParams: five doubles to describe parameterized vacuum wall. Required if useVacuumParams=1.
- adjustVacuumParams: for modelType 0 or 3, if 1, multiply coordinates and parametric values of nodes on vacuum wall by vacuumFactor. Valid only if useVacuumParams=1 (default 0)
- vacuumFactor: for modelType 0 or 3, an optional double value used to multiply coordinates and parametric values of nodes on vacuum wall when adjustVacuumParams=1. Valid only if adjustVacuumParams=1 (default 2×PI)
- numVacuumPts: optional # interpolation points on parameterized vacuum wall. Valid only if useVacuumParams=1 (default 20)
- meshGradationRate: for modelType 3 or 4, optional mesh gradation rate (default: 0.3). This value should be greater than or equal to 0.3. Otherwise the mesh will be fine everywhere.
- resistive-width: for modelType 3, the width of resistive wall. If resistive-width=0, only plasma region is created (default 0.02)
- plasma-offsetX: for modelType 3, the offset in x direction to the left (default 0.0)
- plasma-offsetY: for modelType 3, the offset in y direction to the bottom (default 0.0)
- vacuum-width: for modelType 3 or 4, the width of vacuum region (default 2.5)
- vacuum-height: for modelType 3 or 4, the height of vacuum region (default 4.0)

Locate input parameter file and all files listed as bdryFile (if applicable) in the work folder and do m3dc1\_meshgen input\_param\_file, then the following output will be generated.

- The output model in three formats
  - $-M3D-C^1$ -readable .txt
  - Simmetrix-readable file .smd and
  - PUMI-readable .dmg
    - ▷ For modelType 0-2, the model is saved in outFile.\*
    - ▷ For modelType 3 with resistive width R, vacuum-width W and vacuum-height H, the model is saved in outFile-R-W-H.\*.
    - ▷ For modelType 4 with vacuum-width W and vacuum-height H, the model is saved in outFile-W-H.\*.
- The output mesh in three formats
  - Simmetrix-readable.sms
  - $-M3D C^1/\text{PUMI}$  readable .smb
  - Paraview
    - $\triangleright$  For modelType 0-2 with # mesh faces F,
      - if F > 1000, the mesh is saved in outFile-(F/1000).\*
      - if F < 1000, the mesh is saved in outFile-F.\*
    - $\triangleright\,$  For model Type 3 with # mesh faces F, resistive width R, vacuum width W, vacuum-height H,
      - if F > 1000, the mesh is saved in outFile-R-W-H-(F/1000).\*
      - if F < 1000, the mesh is saved in outFile-R-W-H-F.\*
    - $\triangleright$  For modelType 4 with # mesh faces F, vacuum-width W and vacuum-height H,
      - if F > 1000, the mesh is saved in outFile-W-H-(F/1000).\*
      - if F < 1000, the mesh is saved in outFile-W-H-F.\*

#### 3.1.1 Type 0 (parameterized vacuum)

The figure 1 illustrates a mesh generated by the following input file.

modelType 0
outFile analytic
meshSize 0.04
useVacuumParams 1
vacuumParams 1.65908 0.46 0.2 -0.02504 0.8
numVacuumPts 20
adjustVacuumParams 0
vacuumFactor 6.28319

The figure 2 illustrates a mesh generated with the same vacuum region parameters and a higher mesh size value.



Figure 1: A mesh with vacuum region defined by five parameters for analytic expression



Figure 2: A mesh generated with vacuum region parameters and mesh size  $0.1\,$ 



Figure 3: A mesh with vacuum region defined by piece-wise polynomials

#### 3.1.2 Type 2 (piece-wise polynomial vacuum)

The figure 3 illustrates a mesh generated by the following input file. The vacuum region's geometry information is defined by piece-wise polynomials and stored in the file in-poly.

modelType 2
inFile in-poly
outFile poly

The vacuum region's geometry information is defined by piece-wise polynomials and stored in the file in-poly and the example file can be found in

/p/tsc/m3dc1/lib/SCORECLib/rhel7/intel2019u3-openmpi4.0.3/16.0-220226/bin.

#### 3.1.3 Type 3 (three-regions with inner wall points)

With the model type 3, a geometric model consists of three model faces where each represents plasma region, resistive region and vacuum region, respectively. An ascii file name which describes inner plasma wall boundary has to be provied with the parameter inFile.

The figure 4 illustrates a mesh generated by the following input file. In the figure, geometric model faces are different colored.



Figure 4: A mesh with spline-fitted 3-region model

```
modelType 3
inFile in-circle
outFile circle
meshSize 0.1 0.5 0.1
useVacuumParams 1
adjustVacuumParams 1
vacuumParams 5.0 1.5 0.0 0.0 1.5
numVacuumPts 20
meshGradationRate 0.4
resistive-width 0.4
```

The example files circle-input and in-circle can be found in

/p/tsc/m3dc1/lib/SCORECLib/rhel7/intel2019u3-openmpi4.0.3/16.0-220226/bin.

#### 3.1.4 Type 4 (three-regions with inner & outer wall points)

With the model type 4, a geometric model consists of three model faces where each represents plasma region, resistive region and vacuum region, respectively. The parameter inFile denotes a file name that contains inner plasma wall boundary. The parameter bdryFile denotes a file name that contains resistive wall boundary.



Figure 5: A mesh with inner & outer plama wall boundary points

The figure 5 illustrates a mesh generated by the following input file. In the figure, geometric model faces are different colored.

modelType 4
inFile inner\_bdry.pts
bdryFile outer\_bdry.pts
outFile iter
meshSize 0.7 0.5 0.7
useVacuumParams 1
adjustVacuumParams 1
vacuumParams 8.25 8.0 0.2 0.0 12.5
meshGradationRate 1

The example files bdry-input, inner\_bdry.pts and outer\_bdry.pts can be found in

/p/tsc/m3dc1/lib/SCORECLib/rhel7/intel2019u3-openmpi4.0.3/16.0-220226/bin.

#### $3.2 m3dc1_mfmgen$

m3dc1\_mfmgen requires an ascii input file of arbitrary name that contains the following parameters. The parameters can be in any order.

- numBdry: the number of boundary files defined by peice-wise linear points given for the construction of the loops (default: 0)
  - Each boundary file corresponds to a loop in PUMI
  - For numBdry=N, N lines of bdryFile should be provided,  $N \ge 0$
- bdryFile: For each boundary file, the user has to provide its file name followed by the unique loop ID and desired mesh size on the loop.
  - Each boundary file corresponds to a loop in PUMI
  - The unique ID can be an arbitrary integer defined by the user
  - The unique ID is used with input parameter "faceBdry" to specify the boundaries (loops) of model face
  - For more than one boundary files (numBdry>1), the boundary files can be in any order
- useVacuum: A parameter to control the vacuum boundary
  - The first number sets the mode of vacuum boundary and can be 0, 1 or 2. If 0, no vacuum boundary will be created. If 1, vacuum boundary will be created without user defined parameters. If 2, a parameterized vacuum boundary will be created by using the parameters defined in the parameter "vacuumParams".
  - The second number is the desired unique loop ID for the vacuum loop
  - The third number defines the mesh size on the vacuum boundary
- vacuumParams: if "useVacuum = 2", the user has to provide five doubles to define parameterized vacuum wall
- numVacuumPts: if "useVacuum = 2", # interpolation points on parameterized vacuum wall (default=20)
- thickWall: three integers and one double to control finite thickness wall
  - The first number can either be 0 or 1. If it is set to 0, no finite thickness wall will be created. If it is set to 1, a finite thickness wall will be created
  - The second number is the loop ID that will be offset for given thickness
  - The third number is the desired unique loop ID for the new loop created from offsetting for the finite thickness wall
  - The last number is the desired wall thickness
- layeredMesh: two integers to create an extruded layeded mesh on the finite thickness wall
  - The first integer is 0, no layered mesh will be created. If 1 (default), an extruded mesh with desired number of mesh layers will be created
  - The second integer defines the number of mesh layers
- numFace: the number of geometric model faces in PUMI (default 1)
  - Each geometric face corresponds to regions (e.g. plasma, resistive, vacuum) in M3DC1
  - For numFace=N, N lines of faceBdry should be provided, N > 0

- faceBdry: For each model face, the user has to provide the number of loops, loop ID(s), and desired mesh size
  - The first number gives the total number of loops bounding the face
  - the first number is followed by the loops IDs of the bounding loops. If number of loops = n, there should be n loop ID
  - The last number is the desired mesh size on the geometric face
- meshGradationRate: Global mesh gradation rate for the meshing. This parameter is optional and if not specified a default mesh gradation rate = 0.3 is used. This value should be greater than or equal to 0.3. Otherwise the mesh will be fine everywhere.
- outFile: output file name to save model and mesh

Locate input parameter file and all files listed as bdryFile (if applicable) in the work folder and do m3dc1\_mfmgen input\_param\_file. The output files are the same as those of m3dc1\_meshgen.

#### 3.2.1 Mesh with parameterized vacuum wall

This section presents a mesh created with a parameterized vacuum wall. This is equivalent to Type 0 mesh of m3dc1\_meshgen.

numBdry 0 useVacuum 1 1 0.1 numFace 1 faceBdry 1 1 0.09 outFile analytic-0.09

The figure 6 presents the mesh generated by the input file above with two different mesh sizes.

#### 3.2.2 Mesh with single boundary file

```
numBdry 1
bdryFile loop1.dat 3 0.1
numFace 1
faceBdry 1 3 0.2
outFile input1
```

The figure 7 presents the mesh generated by the input file above.



Figure 6: Mesh with a parameterized vacuum region and different mesh size for model face (left)  $0.2~({\rm right})~0.09$ 



Figure 7: Mesh with single boundary file and no vacuum wall



Figure 8: Mesh with two boundary files and a parameterized vacuum region

#### 3.2.3 Mesh with two boundary files and a parameterized vacuum wall

This section presents a mesh created with two boundary files and a parameterized vacuum wall. This is equivalent to Type 4 mesh of m3dc1\_meshgen.

numBdry 2 bddyFile loop1.pts 1 0.5 bdryFile loop2.pts 2 0.5 useVacuum 2 3 0.5 vacuumParams 8.25 8.0 0.2 0.0 12.5 numFace 3 faceBdry 1 1 0.7 faceBdry 2 1 2 0.5 faceBdry 2 2 3 0.7 meshGradationRate 1 outFile iter

The figure 8 presents the mesh generated by the input file above. As you can see, the mesh in Figure 8 and 5 are almost identical.





#### 3.2.4 Mesh with three boundary files

```
numBdry 3
bdryFile loop1.dat 3 0.1
bdryFile loop2.dat 10 0.05
bdryFile loop3.dat 11 0.09
numFace 3
faceBdry 1 3 0.2
faceBdry 2 3 10 0.1
faceBdry 2 10 11 0.09
outFile input3
```

The figure 9 presents the mesh generated by the input file above.

#### 3.2.5 Mesh with seven boundary files and a vacuum wall

numBdry 7
bdryFile loop1.dat 3 0.2
bdryFile loop2.dat 10 0.3
bdryFile loop3.dat 11 0.4



Figure 10: Mesh with seven boundary files and a parameterized vacuum wall

bdryFile loop4.dat 21 0.4 bdryFile loop5.dat 25 0.4 bdryFile loop6.dat 17 0.2 bdryFile loop7.dat 19 0.1 useVacuum 1 9 0.1 vacuumParams 1.8 1.5 0.4 0.0 2.5 numFace 8 faceBdry 1 3 0.2 faceBdry 1 10 0.3 faceBdry 1 11 0.1 faceBdry 2 21 25 0.2 faceBdry 2 25 17 0.09 faceBdry 2 17 19 0.1 faceBdry 2 19 9 0.1 faceBdry 4 3 10 11 21 0.11 outFile input7

The figure 10 presents the mesh generated by the input file above.

#### 3.2.6 Mesh with finite thickness wall and layers

Add text here

#### 3.3 polar\_meshgen

polar\_meshgen requires an ascii file of arbitrary name that contains input parameters as the following:

- inFile: input file name containing equilibrium generation by jsolver
- outFile: output file name to save model and mesh
- meshSize: relative mesh size for each region (default 0.05)
- reorder: if 1, reorder PUMI mesh based on adjacency (default: 0) and generate vtk folders for mesh visualization. The mesh before and after reodering is saved in original-mesh.vtk and reordered-mesh.vtk, respectively. Note that the element order of Simmetrix mesh is not affected.

The following presents an example input file "polar\_input".



Figure 11: Add Caption Here



Figure 12: Add Caption Here

inFile POLAR outFile polar meshSize 0.04

To run polar\_meshgen, place polar\_input and POLAR in your work folder and do "polar\_meshgen polar\_input". The program will read POLAR and generate various model and mesh files starting with "polar". For instance, polar-2K0.smb, pol-2K.sms, pol-2K.vtk, polar.dmg, polar.smd, polar.txt. If the resulting mesh is too fine, increase the value of meshSize. If the resulting mesh is too coarse, decrease the value of meshSize is not specified in the input file, the default value is 0.05.

The program read\_jsolver generates equilibrium and stores in the file POLAR. Given the input file POLAR, m3dc1\_meshgen generates the following files:

- model.dmg: PUMI-readable model file
- model.txt: M3DC1-readable model file
- mesh0.smb: PUMI/M3DC1-readable mesh file
- mesh.vtk: Paraview data files
- norm\_curv: ascii file containing nodes' normal/curvature information

### 3.4 Mesh Control with SimModeler

- 1. From the menu "File  $\rightarrow$  Open Model", load a model file (.smd) generated by m3dc1\_meshgen
- 2. In the upper panel, in the views section, click on Front to view the model, then go to Meshing tab
- 3. Select outer region, click + in Mesh Attributes and select Mesh Size→relative. Enter a value (typically 0.1)
- 4. Select wall region, click + in Mesh Attributes and select Mesh Size→relative. Enter a value (typically 0.02)
- 5. Select inner region, click + in Mesh Attributes and select Mesh Size→relative. Enter a value (typically 0.04). Here, one can already generate the mesh by clicking on Generate Mesh and verify if the mesh sizes are suitable

- 6. Select both inner and wall regions (holding shift key), click + in Mesh Attributes and select Mesh Size→relative. Enter a function, e.g. 0.01×abs(\$y+1.5)^2+0.004 to specify an anisotropic mesh density on top of previous settings There are many available parameters for fine-tuning the mesh density. For example, Mesh Curvature Refinement with parameter packs more elements near the edges of the resistive wall.
- 7. Generate Mesh and Show Mesh to view result in new windows
- 8. If the result is satisfactory, from the menu File→Save Mesh, give it a meaningful name with the extension .sms. The original model file .smd has been automatically saved by the program with your mesh modifications.
- 9. Close simmodeler then it will release a license. Until you quit Simmodeler, no one cannot run neither m3dc1\_meshgen nor simmodeler.
- 10. Copy the .txt, .smd and .sms files to the simulation directory and run the following splitting routine to obtain PUMI-readable .smb mesh files.

/p/tsc/C1/m3dc1-sunfire.r6-1.5/bin/part\_mesh.sh model\_file.smd mesh\_file.sms X, where X is the number of parts you need in the .smb mesh.

11. Modify the Clinput file accordingly

mesh\_filename = 'part.smb'
mesh\_model = 'filename.txt'

#### 3.5 Mesh Partitioning

#### 3.5.1 Splitting

The program  $split\_smb$  increases the number of parts in a mesh from P to N (P<N). In each machine, the program  $split\_smb$  is available in  $SCOREC\_UTIL\_DIR$  provided in hostname.mk file.

In order to split P-part mesh to N parts (N>P), run "mpirun -np N ./split\_smb input-mesh(.smb) output-mesh(.smb) X"

- the file extension of input-mesh should be .smb
- the file extension of output-mesh should be .smb
- N is the number of parts in the output mesh
- For a P-part input mesh, X must be N/P
- For both input and output mesh, do not specify a number before the file extension
- split\_smb will insert a number in the output mesh file. The number represents a global part ID.

Make sure that the output mesh doesn't have any empty part. Otherwise, the program crashes with the following error message:
 APF warning: 1 empty parts split\_smb: .../mds/mds.c:614: check\_ent: Assertion 'e >= 0' failed

Examples on portal:

- 1. To split a serial (1-part) mesh to 6 parts, run
  "mpirun -np 6 ./split\_smb struct-curveDomain.smb part.smb 6"
  - Input mesh: struct-curveDomain0.smb
  - Output mesh: part0.smb, part1.smb, part2.smb, part3.smb, part4.smb, part5.smb
- 2. To split a 2-part mesh to 6 parts, run "mpirun -np 6 ./split\_smb struct-curveDomain.smb part.smb 3"
  - Input mesh: struct-curveDomain0.smb, struct-curveDomain1.smb
  - Output mesh: part0.smb, part1.smb, part2.smb, part3.smb, part4.smb, part5.smb

See readme.split\_smb for detailed instructions and trouble shooting tips.

#### 3.5.2 Mesh Merging

The program collapse decreases the number of parts in a mesh from N to P (P<N). In each machine, the program collapse is available in  $COREC_UTIL_DIR$  provided in hostname.mk file.

In order to merge N-part.smb mesh to P parts (P>0), run "mpirun -np N ./collapse input-mesh(.smb) output-mesh(.smb) X"

- the file extension of input-mesh should be .smb
- the file extension of output-mesh should be .smb
- N is the number of parts in the input mesh
- For a P-part output mesh, X must be N/P
- For both input and output mesh, do not specify a number before the file extension
- collapse will insert a number in the output mesh file. The number represents a global part ID.

Example on portal:

In order to merge 4-part mesh into a serial (1-part) mesh, run "mpirun -np 4 ./collapse part.smb serial.smb 4"

• Input mesh: part0.smb, part1.smb, part2.smb, part3.smb

• Output mesh: serial0.smb

See readme.collapse for detailed instructions and trouble shooting tips.

## 3.6 Miscellaneous

#### 3.6.1 Verification

The program check\_smb investigates an input mesh and prints any invalid aspects of the mesh. It prints out the mesh size (the number of global, local, and owned entities per dimension) at the end.

In order to run, do "mpirun -np N ./check\_smb input-mesh(.smb)".

4 Mesh Adaptation by Error Estimator

# 5 Output

## 5.1 HDF5 Output

By default, M3D-C1 will output data to a HDF5 file named C1.h5. The file is organized as follows:

- state variables
- scalars/
- equilibrium/
- time\_###/

#### 5.1.1 State variables

#### 5.1.2 The scalars/ group

The scalars/ group contains one-dimensional time series data, output at every MHD timestep. Some of the time series that are output include the following:

Scalar	Units	Description
time	$t_0$	The physical time at each MHD timestep
dt	$t_0$	The physical time per MHD timestep
$loop\_voltage$	$V_0$	The loop voltage applied at the domain boundary
$toroidal\_current$	$I_0$	The total toroidal current in the MHD region
$particle_number$	$n_0 L_0^3$	The total number of main ions in the MHD region
$electron\_number$	$n_0 L_0^{3}$	The total number of electrons in the MHD region
power_injected	$p_0 L_0^{3}/t_0$	The total power from the heat source in the MHD region

For simulations that make use of the KPRAD module, the following time series are also output:

Scalar	$\mathbf{Units}$	Description
radiation	$p_0 L_0^3 / t_0$	Total power loss due to all radiation sources
$line_rad$	$p_0 L_0^3 / t_0$	Power loss due to line radiation
brem_rad	$p_0 L_0^3/t_0$	Power loss due to bremsstrahlung
ion_loss	$p_0 L_0^3/t_0$	Power loss due to ionization
$\operatorname{reck}_{\operatorname{rad}}$	$p_0 L_0^3 / t_0$	
recp_rad	$p_0 L_0^3 / t_0$	
$kprad_n$	$n_0 L_0^3$	Total number of impurity nuclei (both ionized and neutral)
kprad_n0	$n_0 L_0^{3}$	Total number of neutral impurity atoms
$kprad\_dt$	$t_0$	Time step of KPRAD subcycle



Figure 13: The geometry of a mesh element in the toroidal plane<sup>[2]</sup>

#### 5.1.3 The equilibrium/ and time\_###/ groups

The equilibrium/ and time\_###/ groups contain the mesh and field data upon initialization (equilibrium/ and time\_000/) and at each field output time slice (time\_###). The number of MHD timesteps per field output time slice is determined by the Clinput parameter ntimepr.

In calculations with eqsubtract=1, the equilibrium/ group contains the equilibrium part of the fields whereas time\_###/ contain the perturbed part of the fields.

The data in equilibrium/ and time\_### are written to separate files named equilibrium.h5 and time\_###.h5. These files are linked to C1.h5 using the linking feature of HDF5 so that the data can be accessed as if it were a regular group in C1.h5.

Variable	$\mathbf{Units}$	Description
time	$t_0$	Physical time of time slice
nspace	1	2 for 2D / complex; 3 for 3D simulations
$\operatorname{ntimestep}$	1	MHD timestep associated with this times lice
version	1	Version number for output data
$\mathrm{mesh}/$		Mesh data group
$\mathrm{mesh/nelms}$	1	Number of mesh elements
$\operatorname{mesh}/\operatorname{period}$	1	Toroidal period of mesh
$mesh/ifull\_torus$	1	1 if full-torus simulation; 0 if partial-torus simulation
$\mathrm{mesh/nperiods}$	1	If partial torus, number of periods per full torus
$\mathrm{mesh/phi}$	1	1D array containing the toroidal angles of the toroidal planes
mesh/elements		2D array containing mesh data
mesh/adjacency		2D array containing adjacency data
fields/		Group containing field data

mesh/elements is a 2D array consisting of nelms rows (one per mesh element), with 8 (for 2D simulations) or 10 (for 3D simulations) columns:

$$a, b, c, \theta, x_1, y_1$$
, ibound, izone,  $d, \varphi_1$ 

The first five columns (a through  $y_1$ ) describe the geometry of the element in the toroidal plane, as illustrated in figure 13. d and  $\varphi_1$  represent the toroidal extent of the element and the toroidal coordinate of the first bounding plane (in radians if itor=1, in units of  $L_0$  otherwise), and are only present for 3D simulations.

**Fields**/ is a group containing the field data. Each field is a 2D array consisting of **nelms** rows (one per element) and 20 (in 2D) or 80 (in 3D) columns, representing the coefficients of each shape function within the given mesh element.

The field  $\phi$  at the coordinate system  $(\xi, \eta, \zeta)$  local to the mesh element is given by

$$\phi(\xi,\eta,\zeta) = \sum_{j=0}^{J} \sum_{i=1}^{20} a_{i+20j} \xi^{m_i} \eta^{n_i} \zeta^j$$

where

$$m = \{0, 1, 0, 2, 1, 0, 3, 2, 1, 0, 4, 3, 2, 1, 0, 5, 3, 2, 1, 0\}$$
$$n = \{0, 0, 1, 0, 1, 2, 0, 1, 2, 3, 0, 1, 2, 3, 4, 0, 2, 3, 4, 5\}$$

and where J = 3 for 3D simulations and J = 0 for 2D / complex simulations.

The fields that are written include the following. (Note that in 2D complex simulations, these field names contain the real part of the field data; for each field an additional field with the suffix "\_i" is written to contain the imaginary part.)

Field	Units $(itor=1)$	Units $(itor=0)$	Description
psi	$B_0 L_0^2$	$B_0L_0$	$\psi$
Ι	$B_0L_0$	$B_0$	F
f	$B_0L_0$	$B_0 L_0^2$	f
$_{\rm phi}$	$v_0$	$v_0L_0$	U
V	$v_0/L_0$	$v_0$	$\omega$
chi	$v_0 L_0^3$	$v_0L_0$	$\chi$
Р	$p_0$	$p_0$	p
Pe	$p_0$	$p_0$	$p_e$
$\operatorname{ti}$	$T_0$	$T_0$	$T_i$
te	$T_0$	$T_0$	$T_e$
den	$n_0$	$n_0$	$n_i$
ne	$n_0$	$n_0$	$n_e$

# 6 PETSc Option

#### 6.1 2D

The petsc option to run 2D modes is default to superlu\_dist. You can add the following line on your "srun" command line to change to to mumps:

-pc\_factor\_mat\_solver\_type mumps

#### 6.2 3D

When running M3D- $C^1$  in the 3D nonlinear mode, you need to include PETSc Options file. There is a number "8" in the file below. It must be equal to the number of toroidal planes. It should be changed whenever you change the number of toroidal planes in the C1input file. The recommended *options\_bjacobi* file is as follows:

```
-pc_type bjacobi
-pc_bjacobi_blocks 8
    (for 8 toroidal planes should be equal to nplanes in C1input)
-sub_pc_type lu
-sub_pc_factor_mat_solver_package superlu_dist
    (can exchange mumps for superlu_dist)
-mat_superlu_dist_rowperm NOROWPERM (only needed for superlu_dist)
-mat_mumps_icntl_14 50 (only needed for mumps)
    (50 means 50% of memory increase when needed.
     Users can make it 100 or more if encountering a runtime memory issue.)
-sub_ksp_type preonly
-ksp_type fgmres
-ksp_gmres_restart 220
-ksp_rtol 1.e-9
-ksp_max_it 10000
-on_error_abort
-hard_pc_type bjacobi
-hard_pc_bjacobi_blocks 8
    (for 8 toroidal planesshould be equal to nplanes in Clinput)
-hard_sub_pc_type lu
-hard_sub_pc_factor_mat_solver_type superlu_dist
    (can change mumps for superlu_dist)
-mat_superlu_dist_rowperm NOROWPERM
    (only needed for superlu_dist)
-mat_mumps_icntl_14 50
    (only needed for mumps.)
    (50 means 50% of memory increase when needed.
```
```
Users can make it 100 or more if encountering a runtime memory issue.)

-hard_sub_ksp_type preonly

-hard_ksp_type lgmres

-hard_ksp_lgmres_argument 4

-hard_ksp_gmres_restart 220

-hard_ksp_rtol 1.e-9

-hard_ksp_max_it 10000
```

#### 6.3 More

More examples are in regtest/pellet/base/ directory, such as

```
options_bjacobi.type_superludist
options_bjacobi.type_mumps
```

The following are additional optional arguments:

```
-ksp_converged_reason
-ksp_view
```

-help

-options\_table -options\_left

-trdump -malloc\_log

for diagnosing purpose.

## 7 Running Jobs

Users can find almost all of the needed example batch scripts and input files to run a job on available computing facilities from

unstructured/regtest/\*/base/

directories.

#### 7.1 Running 2D or Linear Jobs

In 2D, the run can be either linear or nonlinear, depending on the C1input parameter *linear*:

linear = 0 (non-linear run: must compile with the option RL=1)

linear = 1 (linear run: must compile with the option COM=1)

In both cases, set

nplanes = 1

For the linear case, use

ntor = nn

set the toroidal mode number.

To run your job on a scratch directory, copy the following files over:

```
executable (m3dc1_2d for non-linear run or m3dc1_2d_complex for linear run)
C1input
AnalyticModel (or MultiEdgeAnalyticModel)
struct-dmg.sms
(and geqdsk if needed)
```

To run non-linear job

mpirun np 8 ./m3dc1\_2d

To run linear job

mpirun np 24 ./m3dc1\_2d\_complex -pc\_factor\_mat\_solver\_package mumps

#### 7.2 Running 3D Nonlinear Jobs

For the 3D nonlinear run, set *linear* = 0 and set *nplanes* equal to the number of toroidal planes in C1input file. The number of bjacobi blocks in the PETSc options file must also be equal to *nplanes*. The total number of processors to request must be the product of nplanes and M (the number of processors per plane, which equals the number of mesh partitions per plane).

Files required to be present to the local directory are:

executable (m3dc1 or m3dc1\_st)
C1input,
partnn.smb (one for each poloidal plane partition)
options\_bjacobi
m3dc1.xml (if using ADIOS)
geqdsk (if needed)

To run linear job

mpirun np 16 ./ m3dc1 ipetsc options\_file options\_bjacobi (nplane=2, M=8)

See the previou section for the format of the PETSc option file options\_bjacobi. In this example job, there are M = 8 mesh partition files:

part0.smb part1.smb part2.smb part3.smb part4.smb part5.smb part6.smb part7.smb

#### 7.3 Graphics Files

The graphics files are of two types. There is a single file called: C1.h5 that contains all the timedependent scalar information. This must be saved and be present in the directory of a job so that it can be added to.

In addition to this file, each plot cycle will produce a file: time\_nnn.h5, where nnn is the plot cycle number. The equilibrium is written into a file called equilibrium.h5. These must be stored in the same directory as the C1.h5 file.

#### 7.4 Restarting Jobs

By default hdf5 files are written in every time step. Therefore jobs can be restarted from the hdf5 plot file, the same one that is used by the idl routines to make plots.

By default, the hdf5 files are written in single precision. If idouble\_out is set to 1 in C1input file, hdf5 files are written in double precision.

#### 7.4.1 Reading restart files for 2D real, 2D complex, or 3D real runs

To start a normal simulation with the hdf5 files, set the C1input parameter irestart to 1.

However, the files C1.h5 and the final time\_nnn.h5 file must be present in the working directory. You may also restart from an intermediate time by setting irestart\_slice=nn where nn is the nnth plot file. If this is not set, it will restart from the final plot file.

#### 7.4.2 Reading real restart files to initialize 2D complex calculation

- Run 2D linear=0
- Copy 2D C1.h5 and the final time\_nnnn.h5 to the working directory
- Run 2D (complex) linear=1. In the initial restart, the time and cycle number will start from t=0 and N=0 for the complex run

#### 7.4.3 Running 3D real simulation from 2D real restart files

To start a 3D simulation with 2D restart files, do the following:

- Run 2D
- Copy 2D C1.h5 and the final time\_nnn.h5 to the working directory
- In 3D work folder, set the C1input parameter irestart= 1 Regardless of the time step when the restart files were written, the 3D simulation starts with time step 1.

#### 7.5 Monitoring Jobs

You can monitor the progress of your running job in several ways:

A. C1ke file. Each time step, one line will be added to the ASCII C1ke file in the run directory that you can open with a text editor. The first 4 fields are:

cycle time kinetic\_energy growth\_rate

B. C1.h5 file: You can monitor a time dependent run by using the idl utility described below. Especially useful is the

plot\_scalar,ke

command and also

### plot\_scalar,ke,/growth

C. You can use a text editor to monitor the log file slurm-nnnn.out file (where nnnn is the job number assigned by SLURM)

### 7.6 Exporting Node/Vector/Matrix for Standalone Study

### 7.7 Archiving Data at PPPL

## 8 Boundary Conditions

In all cases, f = 0 on the boundary, and therefore also  $\mathbf{\hat{t}} \cdot \nabla f = 0$ . Some other boundary conditions that may be specified are as follows:

No normal flow (inonormalflow=1) Holds  $\hat{\mathbf{n}} \cdot \mathbf{u}$  constant.

No poloidal flow (inoslip\_pol=1) Holds  $\hat{\mathbf{t}} \cdot \mathbf{u}$  constant.

No toroidal flow (inoslip\_tor=1) Holds  $\hat{\varphi} \cdot \mathbf{u}$  constant.

No normal current (Winocurrent\_pol inocurrent\_norm=1) Holds  $\hat{\mathbf{n}} \cdot \mathbf{J}$  constant.

No poloidal current (inocurrent\_pol=1) Holds  $\hat{\mathbf{t}} \cdot \mathbf{J}$  constant.

No toroidal current (inocurrent\_tor=1) Holds  $\hat{\varphi} \cdot \mathbf{J}$  constant.

$$\hat{\mathbf{n}} \cdot \mathbf{u} = -R\hat{\mathbf{t}} \cdot \nabla U + \frac{1}{R^2} \hat{\mathbf{n}} \cdot \nabla \chi$$
(3)

$$\hat{\mathbf{t}} \cdot \mathbf{u} = R\hat{\mathbf{n}} \cdot \nabla U + \frac{1}{R^2} \hat{\mathbf{t}} \cdot \nabla \chi$$
(4)

$$\hat{\varphi} \cdot \mathbf{u} = R\omega \tag{5}$$

$$\hat{\mathbf{n}} \cdot \mathbf{B} = -\frac{1}{R} \hat{\mathbf{t}} \cdot \nabla \psi - \frac{1}{R^2} \hat{\mathbf{n}} \cdot \nabla f_{\varphi}$$
(6)

$$\hat{\mathbf{t}} \cdot \mathbf{B} = \frac{1}{R} \hat{\mathbf{n}} \cdot \nabla \psi \tag{7}$$

$$\hat{\varphi} \cdot \mathbf{B} = \frac{F}{R} \tag{8}$$

$$\hat{\mathbf{n}} \cdot \mathbf{J} = -\frac{1}{R} \hat{\mathbf{t}} \cdot \nabla F + \frac{1}{R^2} \hat{\mathbf{n}} \cdot \nabla \psi_{\varphi}$$
(9)

$$\hat{\mathbf{t}} \cdot \mathbf{J} = \frac{1}{R} \hat{\mathbf{n}} \cdot \nabla (F + f_{\varphi\varphi}) + \frac{1}{R^2} \hat{\mathbf{t}} \cdot \nabla \psi_{\varphi}$$
(10)

$$\hat{\varphi} \cdot \mathbf{J} = -\frac{1}{R} \Delta^* \psi \tag{11}$$

In the above definitions,  $\hat{\mathbf{n}}$  is the unit vector normal to the boundary surface, and  $\hat{\mathbf{t}} = \hat{\varphi} \times \hat{\mathbf{n}}$ .

#### 9 Discretization

#### 9.1 Finite Elements

Each field is represented as a linear combination of N basis functions  $\nu_i$  on the computational domain

$$U = \sum_{i=1}^{N} \nu_i U_i.$$

The finite element used in M3D- $C^1$  is the reduced quintic element [2], in which the basis functions are fifth order polynomials. At each time step, the projection of the equations onto the basis functions are computed and solved. For example, the equation

$$\frac{\partial U}{\partial t} = F(U)$$

becomes the system of projection equations

$$\int dV \,\nu_i \frac{\partial U}{\partial t} = \int dV \,\nu_i F(U).$$

These projections equations are known collectively as the *weak form* of the equation. Solving the equation in this manner is known as the *Galerkin method*. Hereafter the index i will be dropped from  $\nu_i$ .

Once the equations are cast in the weak form, integrations by parts may be carried out in order to reduce the order of the differential operators acting on the physical fields. For example,

$$\int dV \,\nu \nabla^2 U = \int dV \,\nabla \cdot (\nu \nabla U) - \nabla \nu \cdot \nabla U$$
$$= \oint d\mathbf{A} \cdot \nabla U \nu - \int dV \,\langle \nu, \nabla U \rangle$$
$$= -\int dV \,\langle \nu, U \rangle .$$

It is found that using integrations by parts to re-cast the equations into a form in which a roughly equal number of derivatives acts on the trial function as on the physical fields improves the numerical stability of methods for solving the equations. Thus, in the above example, the form  $-\langle \nu, U \rangle$  is preferable to  $\nu \nabla^2 U$ .

#### 9.1.1 Weak form of Physical Equations

#### Integration Identities

Rather than performing integrations by parts directly on each term in equations (??), it is simpler to begin directly from the vector form, equations (??) and use the following identities when applying the operations to extract the scalar equations:

$$-\int dV R^2 \nu \nabla \varphi \cdot \nabla \times \mathbf{A} = -\int dV \mathbf{A} \cdot \left[\nabla (R^2 \nu) \times \nabla \varphi\right]$$
$$\int dV \nu \nabla \cdot \mathbf{A} = -\int dV \nabla \nu \cdot \mathbf{A}.$$

(Note that the torodal operator,  $R^2 \nu \nabla \varphi$ , is not a differential operator and therefore the integration by parts cannot be performed *a priori*.)

Similary, useful identities for the operators that will act on the stress tensor  $\Pi$  are:

$$R^{2}\nu\nabla\varphi\cdot\nabla\times(\nabla\cdot\Pi) = R^{2}\partial_{Z}\nu\nabla\varphi\cdot\Pi\cdot\nabla\varphi-\nabla\nu\cdot\Pi\cdot\nabla Z \qquad (12a)$$
$$+r\nabla\varphi\cdot\left[\nabla\nabla(\nu r)\dot{\times}\Pi\right]+\nabla\cdot\mathbf{A}_{1}$$

$$-R^{2}\nu\nabla\varphi\cdot(\nabla\cdot\Pi) = R^{2}\nabla\nu\cdot\Pi\cdot\nabla\varphi+\nabla\cdot\mathbf{A}_{2}$$
(12b)

$$-\nu\nabla\cdot(\nabla\cdot\Pi) = -\nabla\nabla\nu:\Pi + \nabla\cdot\mathbf{A}_3 \tag{12c}$$

where

$$\begin{split} \mathbf{A}_1 &= -R^2 \nu \nabla \varphi \times (\nabla \cdot \Pi) - r \Pi \cdot [\nabla \varphi \times \nabla (r\nu)] + \nu \Pi \cdot \nabla Z \\ \mathbf{A}_2 &= -R^2 \nu \Pi \cdot \nabla \varphi \\ \mathbf{A}_3 &= \nabla \nu \cdot \Pi - \nu \nabla \cdot \Pi. \end{split}$$

(These identities hold for any symmetric tensor  $\Pi$ .) The total divergences vanish upon integration.

#### 9.1.2 Physical Equations after Integrations by Parts

$$\int dV N_n = \int dV \left[ N_{nU} + N_{n\chi} + N_{nD} \right]$$
(13a)

$$\int dV \left[ U_{Un} + U_{\chi n} \right] = \int dV \left[ U_{UUn} + U_{VVn} + U_{U\chi n} + U_{\chi\chi n} + U_{\psi\psi} + U_{FF} + U_{U\mu} + U_{\chi\mu} + U_g + U_{UD} + U_{\chi D} + U_{\Pi_{\parallel}} + U_{\Pi_{\star}} \right]$$
(13b)

$$\int dV V_{Vn} = \int dV \left[ V_{VUn} + V_{V\chi n} + V_{\psi F} + V_{V\mu} + V_{VD} + V_{\Pi_{\parallel}} + V_{\Pi_{\star}} \right]$$
(13c)

$$\int dV [X_{Un} + X_{\chi n}] = \int dV [X_{UUn} + X_{VVn} + X_{U\chi n} + X_{\chi\chi n}$$

$$+ X_p + X_{\psi\psi} + X_{FF} + X_{U\mu} + X_{\chi\mu} + X_g$$

$$+ X_{UD} + X_{\chi D} + X_{\Pi_{\parallel}} + X_{\Pi_{\star}}]$$
(13d)

$$\int dV \Psi_{\psi} = \int dV \left[ \Psi_{\psi U} + \Psi_{\psi \chi} + \Psi_{\psi Fn} + \Psi_{\psi \eta} \right]$$
(13e)

$$\int dV F_F = \int dV \left[ F_{FU} + F_{\psi V} + F_{F\chi} + F_{\psi n} + F_{Fn} + F_{p_e n} + F_{F\eta} \right]$$
(13f)

$$\int dV P_p = \int dV \left[ P_{pU} + P_{p\chi} + P_{p_e Fn} + P_{\eta\psi} + P_{\eta F} + P_{\kappa} + P_{\kappa} + P_{\kappa_{\parallel}} + P_{\kappa_{\times}} \right]$$
(13g)

The terms in the above equations are categorized and defined in the following sections. Each term has been integrated by parts to arrive at the simplest expression having for which the order of differentiation on the trial function is roughly equal to that on the physical fields.

#### Basic Terms

The terms in this section are the basic terms in the two-fluid equations, which do not depend on any specific choice of closure.

$$N_{n}(\nu, \dot{n}) = \nu \dot{n}$$

$$N_{nU}(\nu, n, U) = \nu [U, n]$$

$$N_{n\chi}(\nu, n, \chi) = n \langle \nu, \chi \rangle$$

$$N_{nD}(\nu, n, D) = -D \langle \nu, n \rangle$$
(14)

$$\begin{aligned} U_{Un}(\nu, \dot{U}, n) &= -\frac{1}{R^2} n \left\langle R^2 \nu, \dot{U} \right\rangle \\ U_{\chi n}(\nu, \dot{\chi}, n) &= -R^2 \nu \left[ n, \dot{\chi} \right] \\ U_{UUn}(\nu, U, U, n) &= \frac{1}{R^2} n \Delta^* U \left[ R^2 \nu, U \right] + \frac{1}{2R^2} \left\langle U, U \right\rangle \left[ R^2 \nu, n \right] \\ U_{VVn}(\nu, V, V, n) &= \frac{1}{2R^2} \left[ \nu, R^2 \right] VVn \\ U_{U\chi n}(\nu, U, \chi, n) &= \frac{1}{R^2} n \Delta^* U \left\langle R^2 \nu, \chi \right\rangle - \left[ U, \chi \right] \left[ R^2 \nu, n \right] \\ U_{\chi\chi n}(\nu, \chi, \chi, n) &= \frac{1}{2} \left\langle \chi, \chi \right\rangle \left[ R^2 \nu, n \right] \\ U_{\psi\psi}(\nu, \psi, \psi) &= -\frac{1}{R^2} \left[ R^2 \nu, \psi \right] \Delta^* \psi \\ U_{FF}(\nu, F, F) &= -R^2 \nu F \left[ F, \frac{1}{R^2} \right] \\ U_{UD}(\nu, U, D) &= \frac{1}{R^2} \left\langle R^2 \nu, U \right\rangle D \\ U_{\chi D}(\nu, \chi, D) &= - \left[ R^2 \nu, \chi \right] D \end{aligned}$$
(15)

$$V_{Vn}(\nu, V, n) = \nu n \dot{V}$$

$$V_{VUn}(\nu, V, U, n) = \nu n [U, V]$$

$$V_{V\chi n}(\nu, V, \chi, n) = -\nu n \langle \chi, V \rangle$$

$$V_{\psi F}(\nu, \psi, F) = \nu [F, \psi]$$

$$V_{VD}(\nu, V, D) = -\nu VD$$
(16)

$$\begin{aligned} X_{Un}(\nu, \dot{U}, n) &= \nu \left[ n, \dot{U} \right] \\ X_{\chi n}(\nu, \dot{\chi}, n) &= -n \langle \nu, \dot{\chi} \rangle \\ X_p(\nu, p) &= \langle \nu, p \rangle \\ X_{UUn}(\nu, U, U, n) &= -\frac{1}{R^2} n \Delta^* U \langle \nu, U \rangle + \frac{1}{2} n \left\langle \nu, \frac{\langle U, U \rangle}{R^2} \right\rangle \\ X_{VVn}(\nu, V, V, n) &= \frac{1}{2} n V V \left\langle \frac{1}{R^2}, \nu \right\rangle \\ X_{U\chi n}(\nu, U, \chi, n) &= (n \nabla^2 \nu + \langle n, \nu \rangle) [U, \chi] + n \Delta^* U [\nu, \chi] \\ X_{\chi \chi n}(\nu, \chi, \chi, n) &= \frac{1}{2} n \langle \nu, \langle \chi, \chi \rangle \rangle \\ X_{\psi \psi}(\nu, \psi, \psi) &= \frac{1}{R^2} \Delta^* \psi \langle \nu, \psi \rangle \\ X_{FF}(\nu, F, F) &= \frac{1}{R^2} F \langle \nu, F \rangle \\ X_{UD}(\nu, U, D) &= [\nu, U] D \\ X_{\chi D}(\nu, \chi, D) &= \langle \nu, \chi \rangle D \end{aligned}$$

$$(17)$$

$$\begin{aligned}
\Psi_{\psi}(\nu,\dot{\psi}) &= \nu\dot{\psi} \\
\Psi_{\psi U}(\nu,\psi,U) &= \nu [U,\psi] \\
\Psi_{\psi \chi}(\nu,\psi,\chi) &= -\nu \langle \chi,\psi \rangle \\
\Psi_{\psi Fn}(\nu,\psi,F,n) &= d_i \nu \frac{1}{n} [\psi,F] \\
\Psi_{\psi \eta}(\nu,\psi,\eta) &= -\frac{1}{R^2} \langle \psi, R^2 \nu \eta \rangle
\end{aligned} \tag{18}$$

$$F_{F}(\nu, \dot{F}) = \nu \dot{F}$$

$$F_{FU}(\nu, F, U) = R^{2}\nu \left[U, \frac{F}{R^{2}}\right]$$

$$F_{\psi V}(\nu, \psi, V) = R^{2}\nu \left[\frac{V}{R^{2}}, \psi\right]$$

$$F_{F\chi}(\nu, F, \chi) = \frac{F}{R^{2}} \langle R^{2}\nu, \chi \rangle$$

$$F_{\psi n}(\nu, \psi, \psi, n) = d_{i} \frac{\Delta^{*}\psi}{R^{2}n} \left[\psi, R^{2}\nu\right]$$

$$F_{Fn}(\nu, F, F, n) = d_{i}R^{2}\nu F \left[\frac{1}{R^{2}n}, F\right]$$

$$F_{pen}(\nu, p_{e}, n) = d_{i}R^{2}\nu \left[\frac{1}{n}, p_{e}\right]$$

$$F_{F\eta}(\nu, F, \eta) = -\frac{1}{R^{2}}\eta \langle R^{2}\nu, F \rangle$$
(19)

$$P_{p}(\nu, \dot{p}) = \nu \dot{p}$$

$$P_{pU}(\nu, p, U) = \nu [U, p]$$

$$P_{p\chi}(\nu, p, \chi) = \Gamma p \langle \nu, \chi \rangle + (\Gamma - 1)\nu \langle p, \chi \rangle$$

$$P_{p_{e},F,n}(\nu, p_{e}, F, n) = d_{i} \left(\frac{1}{n}\nu [p_{e}, F] + \Gamma \nu p_{e} \left[\frac{1}{n}, F\right]\right)$$

$$P_{\eta,\psi}(\nu, \eta, \psi, \psi) = (\Gamma - 1)\nu \frac{(\Delta^{*}\psi)^{2}}{R^{2}}$$

$$P_{\eta,F}(\nu, \eta, F, F) = (\Gamma - 1)\nu \frac{F^{2}}{R^{2}}$$
(20)

#### Gravity

These terms are obtained assuming a gravitational force of the form given by equation (??).

$$U_g(\nu, n) = g_r \nu [n, R] - g_Z r \nu \langle n, R \rangle$$
  

$$X_g(\nu, n) = \frac{n}{R^2} (g_r \langle \nu, R \rangle + g_Z r [\nu, R])$$
(21)

#### Heat Flux Terms

These terms are obtained assuming a heat flux density of the form described in section ??.

$$P_{\kappa}(\nu,\kappa,T) = -(\Gamma-1)\kappa \langle \nu,T \rangle$$

$$P_{\kappa_{\parallel}}(\nu,\kappa_{\parallel},T,\psi,\psi,B^{-2}) = -(\Gamma-1)\kappa_{\parallel}\frac{1}{B^{2}}[\psi,\nu][\psi,T]$$

$$P_{\kappa_{\times}}(\nu,\kappa_{\times},T,F,B^{-2}) = (\Gamma-1)\kappa_{\times}\frac{F}{B}[\nu,T]$$
(22)

$$T = p/n$$
  
$$B^{2} = \frac{1}{R^{2}} \left[ \langle \psi, \psi \rangle + F^{2} \right]$$

Isotropic Viscosity

These terms result from isotropic viscosity of the form given by equation (??).

$$U_{U\mu}(\nu, U, \mu) = \frac{1}{R^2} \left[ \left( \left\langle \mu, R^2 \nu \right\rangle + \mu \Delta^*(R^2 \nu) \right) \Delta^* U + \nabla^2 \mu \left\langle R^2 \nu, U \right\rangle + \Delta^*(R^2 \nu) \left\langle \mu, U \right\rangle \right] U_{\chi\mu}(\nu, \chi, \mu) = -\nabla^2 (R^2 \nu) \left[ \mu, \chi \right] - \Delta^* \mu \left[ R^2 \nu, \chi \right] - \frac{1}{R^2} \Delta^* (R^2 \chi) \left[ R^2 \nu, \mu \right] V_{V\mu}(\nu, V, \mu) = \left[ \left\langle \nu, \mu \right\rangle + \frac{1}{R^2} \mu \Delta^* (R^2 \nu) \right] V X_{U\mu}(\nu, U, \mu) = \nabla^2 \nu \left[ \mu, U \right] + \nabla^2 \mu \left[ \nu, U \right] + \Delta^* U \left[ \nu, \mu \right] X_{\chi\mu}(\nu, \chi, \mu, \mu_c) = \nabla^2 \nu \left\langle \mu, \chi \right\rangle + \nabla^2 \mu \left\langle \nu, \chi \right\rangle + 2\mu_c \nabla^2 \nu \nabla^2 \chi$$
(23)

#### Parallel Viscosity

These terms are obtained assuming a parallel viscosity of the form given in equation (??). These equations were obtained using equations (12). For compactness, derivatives are written as subscripts in the following expressions (*i.e.*  $\nu_Z = \partial_Z \nu$ ).

$$U_{\Pi_{\parallel}U}(\nu, U) = \mu_{\parallel}D_{U}$$
  

$$U_{\Pi_{\parallel}V}(\nu, V) = \mu_{\parallel}D_{U}$$
  

$$U_{\Pi_{\parallel}\chi}(\nu, \chi) = \mu_{\parallel}D_{U}$$
(24)

$$V_{\Pi_{\parallel}U}(\nu, U) = \mu_{\parallel}D_V$$

$$V_{\Pi_{\parallel}V}(\nu, V) = \mu_{\parallel}D_V$$

$$V_{\Pi_{\parallel}\chi}(\nu, \chi) = \mu_{\parallel}D_V$$
(25)

$$X_{\Pi_{\parallel}U}(\nu, U) = \mu_{\parallel}D_X$$
  

$$X_{\Pi_{\parallel}V}(\nu, V) = \mu_{\parallel}D_X$$
  

$$X_{\Pi_{\parallel}\chi}(\nu, \chi) = \mu_{\parallel}D_X$$
(26)

$$D_U = \frac{3}{B^2} \left\{ -\frac{1}{2} R^2 \left[ \nu, \frac{\langle \psi, \psi \rangle}{R^2} \right] + \langle \psi, [\nu, \psi] \rangle - \frac{1}{R^2} F^2 \nu_Z \right. \\ \left. - \frac{2}{R^2} \left[ \nu_Z (\psi_Z^2 - \psi_R^2) + 2\nu_r \psi_r \psi_Z \right] \right\} \\ D_V = -3 \frac{F}{B^2} \left[ \nu, \psi \right] \\ D_X = -\nabla^2 \nu \left( 1 - \frac{3}{R^2} \frac{\langle \psi, \psi \rangle}{B^2} \right) \\ \left. + \frac{3}{R^2 B^2} \left( \frac{1}{2} R^2 \left\langle \nu, \frac{\langle \psi, \psi \rangle}{R^2} \right\rangle - \langle \psi, \langle \nu, \psi \rangle \rangle + \frac{1}{R} F^2 \nu_r \right) \right\}$$

$$\begin{split} \mu_{\parallel_{U}} &= \eta_{0} \frac{p_{i} \tau_{i}}{R^{2} B^{2}} \left( -\frac{1}{2} R^{2} \left[ U, \frac{\langle \psi, \psi \rangle}{R^{2}} \right] + \langle \psi, [U, \psi] \rangle - \frac{1}{R^{2}} F^{2} U_{Z} \right) \\ \mu_{\parallel_{V}} &= -\eta_{0} p_{i} \tau_{i} \frac{F}{B^{2}} \left[ \psi, \frac{V}{R^{2}} \right] \\ \mu_{\parallel_{\chi}} &= \eta_{0} \frac{p_{i} \tau_{i}}{R^{2} B^{2}} \left( \frac{1}{2} R^{2} \left\langle \chi, \frac{\langle \psi, \psi \rangle}{R^{2}} \right\rangle - \langle \psi, \langle \chi, \psi \rangle \rangle + \frac{1}{R} F^{2} \chi_{r} \\ &+ \nabla^{2} \chi \left\langle \psi, \psi \right\rangle ) \end{split}$$

Gyroviscosity

These terms are obtained using equations (12) assuming a gyroviscosity of the form given by equation (??).

$$\begin{split} U_{\Pi_{\times}U}(\nu,U) &= -\frac{p_{i}F}{2R^{3}B^{2}} \\ &\times \left\{ \begin{array}{l} \left(1 + \frac{3}{2R^{2}}\frac{\langle\psi,\psi\rangle}{B^{2}}\right) \begin{bmatrix} ([R^{3}\nu_{Z}]_{Z} - [R^{3}\nu_{r}]_{r}) \left(\begin{bmatrix} U_{R}\\R \end{bmatrix}_{Z} + \begin{bmatrix} U_{Z}\\R \end{bmatrix}_{r} \right) \\ -([R^{3}\nu_{r}]_{Z} + [R^{3}\nu_{Z}]_{r}) \left(\begin{bmatrix} U_{Z}\\R \end{bmatrix}_{Z} - \begin{bmatrix} U_{R}\\R \end{bmatrix}_{r} \right) \end{bmatrix} \\ &+ \frac{9}{2rB^{2}} \\ &\times \begin{bmatrix} (\psi_{Z}^{2} - \psi_{R}^{2}) \left(r\nu_{Z} \left[ \left(\frac{U_{Z}}{R}\right)_{Z} - \left(\frac{U_{R}}{R}\right)_{r} \right] - \frac{1}{R^{3}}U_{Z} \left[ (R^{3}\nu_{Z})_{Z} - (R^{3}\nu_{r})_{r} \right] \right) \\ &+ 2\psi_{r}\psi_{Z} \left(r\nu_{Z} \left[ \left(\frac{U_{R}}{R}\right)_{Z} + \left(\frac{U_{Z}}{R}\right)_{r} \right] - \frac{1}{R^{3}}U_{Z} \left[ (R^{3}\nu_{r})_{Z} + (R^{3}\nu_{Z})_{r} \right] \right) \end{bmatrix} \end{split}$$

$$\begin{split} U_{\Pi_{\times}V}(\nu,V) &= -\frac{p_i}{B^2} \\ \times \left\{ \begin{array}{l} \frac{1}{4R^2} \left( 1 - \frac{3F^2}{B^2R^2} \right) \\ &\times \left( \left\langle \frac{V}{R^2}, R^4 \left[ \psi, \nu \right] \right\rangle - \left\langle \psi, R^4 \left[ \nu, \frac{V}{R^2} \right] \right\rangle + \frac{1}{R^2} \left[ \nu, R^6 \left\langle \frac{V}{R^2}, \psi \right\rangle \right] \right) \\ &- \frac{3}{4B^2} \left[ \psi, \frac{V}{R^2} \right] \\ &\times \left( 2 \left\langle \psi, \langle \psi, \nu \rangle \right\rangle - R^2 \left\langle \nu, \frac{\langle \psi, \psi \rangle}{R^2} \right\rangle - \Delta^* \nu \left\langle \psi, \psi \right\rangle + 6 \psi_Z \left[ \nu, \psi \right] \right) \\ &+ \frac{9F^2}{2B^2R^2} \nu_Z \left\langle \psi, \frac{V}{R^2} \right\rangle \end{split}$$

$$U_{\Pi_{\times\chi}}(\nu,\chi) = -\frac{p_i F}{2R^3 B^2} \\ \times \begin{cases} \left[ (\chi_{RR} - \chi_{ZZ}) \left( [R^3 \nu_r]_r - [R^3 \nu_Z]_Z \right) + 2\chi_{RZ} \left( [R^3 \nu_r]_Z + [R^3 \nu_Z]_r \right) \right] \\ + \frac{3}{R^2 B^2} \begin{bmatrix} \left( \Delta^* \chi [\psi_Z^2 - \psi_R^2] - \chi_{ZZ} \psi_Z^2 + \chi_{RR} \psi_R^2 \right) \left( [R^3 \nu_r]_r - [R^3 \nu_Z]_Z \right) \\ + 2\chi_{RZ} \left( \psi_Z^2 [R^3 \nu_r]_Z + \psi_R^2 [R^3 \nu_Z]_r \right) \\ - 2\psi_r \psi_Z \left( \left[ \chi_{ZZ} - \frac{1}{R} \chi_r \right] [R^3 \nu_r]_Z + \left[ \chi_{RR} - \frac{1}{R} \chi_r \right] [R^3 \nu_Z]_r \right) \end{bmatrix} \end{cases}$$

$$\begin{split} V_{\Pi_{\times}U}(\nu,U) &= \frac{p_i}{4rB^2} \\ \times \left\{ \begin{array}{l} \left(1 - \frac{3}{R^2}\frac{F^2}{B^2}\right) \left(\langle\psi, R\left[U,\nu\right]\rangle + \langle\nu, R\left[U,\psi\right]\rangle - \frac{1}{R^3}\left[U, R^4\left\langle\nu,\psi\right\rangle\right] + U_r\left[\nu,\psi\right] + \frac{2}{R}\psi_Z\left\langle\nu,U\right\rangle\right) \\ &+ \frac{3}{RB^2}\left[\psi,\nu\right] \left(2\left\langle\psi,\left\langle U,\psi\right\rangle\right\rangle - \Delta^*U\left\langle\psi,\psi\right\rangle - \frac{1}{R^2}\left\langle U, R^2\left\langle\psi,\psi\right\rangle\right\rangle + \left[\psi, R^2\right]\left[\psi,U\right]\right) \\ &- \frac{18}{R^2}\frac{F^2}{B^2}\left[U, R\right]\left\langle\nu,\psi\right\rangle \end{split} \right\}$$

$$V_{\Pi_{\times}V}(\nu,V) = \frac{p_i F R^2}{4B^2} \left(1 - \frac{3}{R^2} \frac{\langle \psi, \psi \rangle - F^2}{B^2}\right) \left[\nu, \frac{V}{R^2}\right]$$

$$\begin{split} V_{\Pi_{\times\chi}}(\nu,\chi) &= \frac{p_i}{B^2} \\ &\times \left\{ \begin{array}{l} \left( \frac{1}{R^2} \left\langle \chi, R^2 \left\langle \nu, \psi \right\rangle \right\rangle - \left\langle \nu, \left\langle \chi, \psi \right\rangle \right\rangle - \left\langle \psi, \left\langle \nu, \chi \right\rangle \right\rangle \right) \\ &+ \frac{3}{2rB^2} \left[ \psi, \nu \right] \left( \left\langle \psi, R \left[ \chi, \psi \right] \right\rangle - \frac{1}{2} r \left[ \chi, \left\langle \psi, \psi \right\rangle \right] \right) \\ &+ \frac{3}{4R^2} \frac{F^2}{B^2} \left( \left\langle \psi, \left\langle \chi, \nu \right\rangle \right\rangle + \left\langle \nu, \left\langle \chi, \psi \right\rangle \right\rangle - \left\langle \chi, \left\langle \nu, \psi \right\rangle \right\rangle - 2\Delta^* \chi \left\langle \nu, \psi \right\rangle \right) \end{array} \right\} \end{split}$$

$$\begin{split} X_{\Pi_{\times}U}(\nu,U) &= \frac{p_{i}F}{2R^{2}B^{2}} \\ &\times \left\{ \begin{array}{l} \langle \langle \nu,U \rangle \rangle - R^{2}\left[ [\nu,U] \right] + \frac{1}{R} \left[ U_{r}(\nu_{ZZ} - \nu_{RR}) - 2U_{Z}\nu_{RZ} - \frac{1}{R}U_{r}\nu_{r} \right] \\ &+ \frac{3}{RB^{2}} \left[ \begin{array}{l} \left( \left[ \frac{U_{Z}}{R} \right]_{Z} - \left[ \frac{U_{R}}{R} \right]_{r} \right) \left( \nu_{ZZ}\psi_{R}^{2} - \nu_{RR}\psi_{Z}^{2} + \frac{1}{R}\nu_{r}[\psi_{Z}^{2} - \psi_{R}^{2}] \right) \\ &+ 2\nu_{RZ} \left( \left[ \frac{U_{R}}{R} \right]_{Z}\psi_{R}^{2} + \left[ \frac{U_{Z}}{R} \right]_{r}\psi_{Z}^{2} - \frac{1}{R^{2}}U_{Z}[\psi_{Z}^{2} - \psi_{R}^{2}] \right) \\ &- 2\psi_{r}\psi_{Z} \left( \begin{array}{l} \left[ \frac{U_{R}}{R} \right]_{Z}\nu_{RR} + \left[ \frac{U_{Z}}{R} \right]_{r}\nu_{ZZ} - \frac{1}{R^{2}}U_{Z}[\nu_{ZZ} - \nu_{RR}] \\ &- \frac{1}{R}\nu_{r} \left[ \left( \frac{U_{R}}{R} \right)_{Z} + \left( \frac{U_{Z}}{R} \right)_{r} \right] \end{array} \right) \end{array} \right] \end{split}$$

$$\begin{split} X_{\Pi_{\times}V}(\nu,V) &= \frac{p_i}{4B^2} \\ &\times \left\{ \begin{array}{l} \left(1 - \frac{3}{R^2} \frac{F^2}{B^2}\right) \left(\frac{1}{R^2} \left\langle \nu, R^2 \left\langle \frac{V}{R^2}, \psi \right\rangle \right\rangle - \left\langle \psi, \left\langle \frac{V}{R^2}, \nu \right\rangle \right\rangle - \left\langle \frac{V}{R^2}, \left\langle \psi, \nu \right\rangle \right\rangle \right) \\ &+ \frac{6}{B^2} \left[\psi, \frac{V}{R^2}\right] \left(\frac{1}{R} \left\langle \psi, R\left[\nu, \psi\right] \right\rangle - \frac{1}{2} \left[\nu, \left\langle \psi, \psi \right\rangle \right]\right) \\ &- 6 \frac{F^2}{B^2} \left\langle \psi, \frac{V}{R^2} \right\rangle \left[ \left(\frac{\nu_Z}{R^2}\right)_Z + \left(\frac{\nu_R}{R^2}\right)_r \right] \end{split} \right\} \end{split}$$

$$\begin{split} X_{\Pi_{\times\chi}}(\nu,\chi) &= -\frac{p_i F}{B^2} \\ &\times \left\{ \begin{array}{l} \left(1 + \frac{3}{2R^2} \frac{\langle \psi,\psi \rangle}{B^2}\right) [\langle \nu,\chi \rangle] \\ &+ \frac{3}{2B^2} \left[ \begin{array}{l} \left(-\frac{1}{2} \left[\chi,\langle\psi,\psi\rangle\right] + \frac{1}{R} \left\langle\psi,R\left[\chi,\psi\right]\right\rangle\right) \left(\left[\frac{\nu_R}{R^2}\right]_r + \left[\frac{\nu_Z}{R^2}\right]_Z\right) \\ &- \left(-\frac{1}{2} \left[\nu,\langle\psi,\psi\rangle\right] + \frac{1}{R} \left\langle\psi,R\left[\nu,\psi\right]\right\rangle\right) \left(\left[\frac{\chi_R}{R^2}\right]_r + \left[\frac{\chi_Z}{R^2}\right]_Z\right) \end{array} \right] \end{array} \right\} \end{split}$$

#### 9.1.3 Spatial Integration

The integrals required to calculate the weak-form equations of the Galerkin method are computed numerically using a 79-point Gaussian quadrature. That is, the value of each field is calculated at 79 points for each triangular element, and a weighted sum of these values is computed to approximate the integral.

$$\int dA \ f(x) \simeq \sum_{i=1}^{79} w_i f(x_i),$$

where the integrand f(x) is restricted to a single element. The sampling points and weights appropriate for a equilateral triangle are taken from ref. [1].

The coordinates of the sampling points are given in the "natural coordinates"  $(\alpha, \beta, \gamma)$  in ref. [1]. These coordinates may be converted to cartesian coordinates (r, Z) for an equilateral triangle e having vertices

$$\left\{ \left(-\frac{\sqrt{3}}{2}, -\frac{1}{2}\right), \left(\frac{\sqrt{3}}{2}, -\frac{1}{2}\right), (0, 1) \right\}$$

using the linear transformation

$$\phi_{n \to e}(\alpha, \beta, \gamma) = \left(\frac{\sqrt{3}}{2}(\beta - \gamma), \frac{1}{2}(3\alpha - 1)\right).$$

The weights must be multiplied by the Jacobian of this transformation,

$$\mathcal{J}_{\phi_{n\to e}} = \frac{3\sqrt{3}}{4}.$$

To find the coordinates of the sampling points for a general triangle g having vertices  $\{(-b, 0), (a, 0), (0, c)\}$ , as in ref. [2], one may use the linear transformation

$$\phi_{e \to g}(r, Z) = \left(\frac{a+b}{\sqrt{3}}x + \frac{a-b}{3}(1-y), \frac{c}{3}(2y+1)\right)$$
  
$$\mathcal{J}_{\phi_{e \to g}} = \frac{2c}{3\sqrt{3}}(a+b).$$

The transformation from natural coordinates to cartesian coordinates for a triangle having vertices  $\{(-b, 0), (a, 0), (0, c)\}$  is therefore

$$\phi_{n \to g}(r, Z) = \left(\frac{1}{2}(a+b)(\beta-\gamma) + \frac{1}{2}(a-b)(1-\alpha), c\alpha\right)$$
$$\mathcal{J}_{\phi_{n \to g}} = \frac{1}{2}(a+b)c.$$

The 79-point quadrature gives the exact results for integrands which are polynomials of degree 20 (or less). In the case of quintic finite elements, this means the integration is exact for terms involving products of three fields or fewer, not including the degree-five trial function  $\nu$ . In cylindrical geometry, the presence factors of 1/R will cause the quadrature not to be exact, as 1/R is not in the form of a polynomial. The weights  $w_i$  must also be multiplied by  $R_i$  in cylindrical coordinates to account for the Jacobian of the transformation from cartesian to cylindrical coordinates.

#### 9.2 Time Step

#### 9.2.1 Implicit Time Advance

For the implicit time advance, equations (13) are evaluated at the  $\theta$ -advanced time (e.g.  $F(\psi) \rightarrow F(\psi + \theta \delta t \dot{\psi} + \cdots)$ ), linearized (*i.e.*  $\mathcal{O}(\delta t^2)$  and higher are dropped), and then discretized temporally according to the chosen time integration method (*i.e.*  $\dot{\psi} \rightarrow (\psi^{(n+1)} - \psi^{(n)})/\delta t$ ).

$$\begin{pmatrix} S_{11}^{v} & R_{11}^{v} & S_{12}^{v} & R_{12}^{v} & S_{13}^{v} & 0 & R_{14}^{v} & R_{13}^{v} \\ R_{11}^{B} & S_{11}^{B} & R_{12}^{B} & S_{12}^{B} & R_{13}^{B} & S_{13}^{B} & 0 & 0 \\ S_{21}^{v} & R_{21}^{v} & S_{22}^{v} & R_{22}^{v} & S_{23}^{v} & 0 & R_{24}^{v} & R_{23}^{v} \\ R_{21}^{B} & S_{21}^{B} & R_{22}^{v} & S_{22}^{B} & R_{23}^{B} & S_{23}^{B} & 0 & 0 \\ S_{31}^{v} & R_{31}^{v} & S_{32}^{v} & R_{32}^{v} & S_{32}^{v} & R_{33}^{u} & S_{33}^{v} & 0 \\ R_{31}^{n} & 0 & R_{32}^{n} & 0 & R_{33}^{n} & 0 & S^{n} & 0 \\ R_{31}^{n} & 0 & R_{32}^{p} & 0 & R_{33}^{n} & 0 & S^{n} & 0 \\ R_{31}^{p} & 0 & R_{32}^{p} & 0 & R_{33}^{p} & 0 & 0 & S^{p} \end{pmatrix} \begin{pmatrix} U \\ \psi \\ R_{11}^{v} \\ Q_{11}^{p} & Q_{12}^{v} & Q_{12}^{v} & Q_{13}^{v} & 0 & Q_{14}^{v} & Q_{13}^{v} \\ Q_{11}^{P} & D_{12}^{p} & Q_{12}^{p} & Q_{13}^{p} & D_{13}^{P} & 0 & 0 \\ D_{21}^{v} & Q_{21}^{v} & D_{22}^{v} & Q_{22}^{v} & D_{23}^{v} & 0 & Q_{24}^{v} & Q_{23}^{v} \\ Q_{21}^{B} & D_{21}^{B} & Q_{22}^{B} & D_{22}^{B} & Q_{23}^{B} & D_{23}^{B} & 0 & 0 \\ D_{31}^{v} & Q_{31}^{v} & D_{32}^{v} & Q_{32}^{v} & D_{33}^{v} & 0 & Q_{34}^{v} & Q_{33}^{v} \\ Q_{31}^{B} & D_{31}^{v} & Q_{32}^{0} & 0 & Q_{33}^{v} & 0 & 0 \\ Q_{31}^{v} & 0 & Q_{32}^{v} & 0 & Q_{33}^{v} & 0 & 0 & D^{p} \end{pmatrix} \begin{pmatrix} U \\ \psi \\ V \\ F \\ \chi \\ p_{e} \\ n \\ p \end{pmatrix}$$
(27)

#### 9.2.2 Split Time Step Method

Time is advanced using a split time-step method in which the velocity field is advanced first, then the density and total pressure fields are advanced separately, and finally the magnetic field and electron pressure are advanced together. Though the velocity and magnetic field are advanced separately, the Alfvén and magnetosonic waves are treated implicitly by using equations (??–??) to calculate analytically the advanced-time values of the pressure and magnetic field for use in the velocity time step.

$$\begin{pmatrix} S_{11}^{v} & S_{12}^{v} & S_{13}^{v} \\ S_{21}^{v} & S_{22}^{v} & S_{23}^{v} \\ S_{31}^{v} & S_{32}^{v} & S_{33}^{v} \end{pmatrix} \begin{pmatrix} U \\ V \\ \chi \end{pmatrix}^{(n+1)}$$

$$= \begin{pmatrix} D_{11}^{v} & D_{12}^{v} & D_{13}^{v} \\ D_{21}^{v} & D_{22}^{v} & D_{23}^{v} \\ D_{31}^{v} & D_{32}^{v} & D_{33}^{v} \end{pmatrix} \begin{pmatrix} U \\ V \\ \chi \end{pmatrix}^{(n)} + \begin{pmatrix} Q_{11}^{v} & Q_{12}^{v} & Q_{13}^{v} \\ Q_{21}^{v} & Q_{22}^{v} & Q_{23}^{v} \\ Q_{31}^{v} & Q_{32}^{v} & Q_{33}^{v} \end{pmatrix} \begin{pmatrix} \psi \\ F \\ p \end{pmatrix}^{(n)}$$

$$+ \begin{pmatrix} Q_{1}^{v} \\ Q_{2}^{v} \\ Q_{31}^{v} & Q_{32}^{v} & Q_{33}^{v} \end{pmatrix} \begin{pmatrix} \psi \\ F \\ p \end{pmatrix}^{(n)}$$

$$S^{n}n^{(n+1)} = D^{n}n^{(n)} + \begin{pmatrix} R_{1}^{n} & R_{2}^{n} & R_{3}^{n} \end{pmatrix} \begin{pmatrix} \mathbf{U} \\ \mathbf{V} \\ \chi \end{pmatrix}^{(n+1)} + \begin{pmatrix} Q_{1}^{n} & Q_{2}^{n} & Q_{3}^{n} \end{pmatrix} \begin{pmatrix} \mathbf{U} \\ \mathbf{V} \\ \chi \end{pmatrix}^{(n)}$$
(29)

$$S^{p}p^{(n+1)} = D^{p}p^{(n)} + \begin{pmatrix} R_{1}^{p} & R_{2}^{p} & R_{3}^{p} \end{pmatrix} \begin{pmatrix} \boldsymbol{U} \\ \boldsymbol{V} \\ \boldsymbol{\chi} \end{pmatrix}^{(n+1)} + \begin{pmatrix} Q_{1}^{p} & Q_{2}^{p} & Q_{3}^{p} \end{pmatrix} \begin{pmatrix} \boldsymbol{U} \\ \boldsymbol{V} \\ \boldsymbol{\chi} \end{pmatrix}^{(n)}$$
(30)

$$\begin{pmatrix}
S_{11}^{B} & S_{12}^{B} & S_{13}^{B} \\
S_{21}^{B} & S_{22}^{B} & S_{23}^{B} \\
S_{31}^{B} & S_{32}^{B} & S_{33}^{B}
\end{pmatrix} \begin{pmatrix}
\psi \\
F \\
p_{e}
\end{pmatrix}^{(n+1)}$$

$$= \begin{pmatrix}
D_{11}^{B} & D_{12}^{B} & D_{13}^{B} \\
D_{21}^{B} & D_{22}^{B} & D_{23}^{B} \\
D_{31}^{B} & D_{32}^{B} & D_{33}^{B}
\end{pmatrix} \begin{pmatrix}
\psi \\
F \\
p_{e}
\end{pmatrix}^{(n)} + \begin{pmatrix}
R_{11}^{B} & R_{12}^{B} & R_{13}^{B} \\
R_{21}^{B} & R_{22}^{B} & R_{23}^{B} \\
R_{31}^{B} & R_{32}^{B} & R_{33}^{B}
\end{pmatrix} \begin{pmatrix}
U \\
V \\
\chi
\end{pmatrix}^{(n+1)}$$

$$+ \begin{pmatrix}
Q_{11}^{B} & Q_{12}^{B} & Q_{13}^{B} \\
Q_{21}^{B} & Q_{22}^{B} & Q_{23}^{B} \\
Q_{31}^{B} & Q_{32}^{B} & Q_{33}^{B}
\end{pmatrix} \begin{pmatrix}
U \\
V \\
\chi
\end{pmatrix}^{(n)} + \begin{pmatrix}
O_{1}^{B} \\
O_{2}^{B} \\
O_{3}^{B}
\end{pmatrix}$$
(31)

#### Linear Calculations

Linear calculations may be performed by calculating each matrix once, and recalculating the matrixvector products each time step with the updated vectors. This method is very efficient because the S matrices need only be inverted once, and in all subsequent time steps the only matrix operations carried out are addition and matrix-vector multiplication. Non-linear simulations require all the matrices to be recalculated each time step, and the S matrices must be inverted each time step.

#### Implementation of Electron Pressure Equation

Because it is the electron pressure which appears in the generalized Ohm's law, equation (??), if the electron pressure equation is retained, it is solved with the magnetic field as the third row in equation (31) so as to keep the fast magnetosonic wave implicit. In this case, the full pressure is evolved independently in equation (30). If the electron pressure equation is not included, the third row in equation (31) is the total pressure equation, and the electron pressure is assumed to remain always at a specific fraction of the total pressure, which is determined by the initial conditions.

#### 9.2.3 Crank-Nicholson

The Crank-Nicholson time step is defined by the following discretization:

$$\frac{\partial U}{\partial t} \rightarrow \frac{U^{(n+1)} - U^{(n)}}{\delta t} U \rightarrow \theta U^{(n+1)} + (1-\theta)U^{(n)}.$$

By Taylor expanding about U,

$$U^{(n+1)} = U + \theta \,\delta t \,\dot{U} + \frac{1}{2} \theta^2 \delta t^2 \ddot{U} + \frac{1}{6} \theta^3 \delta t^3 \ddot{U} + \cdots$$
$$U^{(n)} = U + (\theta - 1) \delta t \,\dot{U} + \frac{1}{2} (\theta - 1)^2 \delta t^2 \ddot{U} + \frac{1}{6} (\theta - 1)^3 \delta t^3 \ddot{U} + \cdots$$

the trunctation error of the time-derivative operator can be calculated directly:

$$\Delta_{CN}(\delta t, \theta) = \frac{U^{(n+1)} - U^{(n)}}{\delta t} - \dot{U}$$
  
=  $\left(\theta - \frac{1}{2}\right) \delta t \ddot{U} + \frac{1}{2} \left(\theta^2 - \theta + \frac{1}{3}\right) \delta t^2 \ddot{U} + \cdots$ 

When  $\theta = 1/2$ , the time differencing is "time-centered" because the two points involved in the time differencing are equidistant in the logical time coordinate from the point at which the field itself is evaluated. In this case, the leading-order truncation error is  $\mathcal{O}(\delta t^2)$ :

$$\Delta_{CN}(\delta t, \theta = 1/2) = \frac{1}{24} \delta t^2 \ddot{U} + \cdots .$$
(32)

#### 9.2.4 BDF2

The BDF2 time step is defined by the following discretization:

$$\begin{array}{rcl} \frac{\partial U}{\partial t} & \rightarrow & \frac{3U^{(n+1)} - 4U^{(n)} + U^{(n-1)}}{2\,\delta t} \\ U & \rightarrow & U^{(n+1)}. \end{array}$$

Taylor expanding about U:

$$U^{(n+1)} = U$$
  

$$U^{(n)} = U - \delta t \, \dot{U} + \frac{1}{2} \delta t^2 \ddot{U} - \frac{1}{6} \delta t^3 \ddot{U} + \cdots$$
  

$$U^{(n-1)} = U - 2 \, \delta t \, \dot{U} + 2 \, \delta t^2 \ddot{U} - \frac{4}{3} \delta t^3 \ddot{U} + \cdots$$

and the truncation error is

$$\Delta_{BDF2}(\delta t) = \frac{3U^{(n+1)} - 4U^{(n)} + U^{(n-1)}}{2\,\delta t} - \dot{U} = -\frac{1}{3}\delta t^2 \ddot{U} + \cdots .$$
(33)

# 10 Input Parameters

## 10.1 Model Options

Option	Default	Description
numvar	3	MHD model. 1: 2-field; 2: 4-Field; 3: 6-Field.
linear	0	1: linear (perturbation terms only, no matrix recalculation)
eqsubtract	0	1: remove equilibrium terms from equations
icsubtract	0	set to 1 if PF coils are in the domain. These are defined in
		the files "coils.dat" and "current.dat"
extsubtract	0	1: subtract fields from non-axisymmetric coils
idens	1	1: include density equation
ipres	0	1: include electron pressure equation
ipressplit	0	1: seperate pressure solve from the magnetic field solves when
		isplitstep=1. (ipressplit must be 0 for isplitstep= $0$ )
itemp	0	1: Advance temperatures rather than pressures (for ipress-
		plit=1 only)
gyro	0	1: include Braginskii gyroviscous term. (note: needs db to
		be nonzero also)
igauge	0	0: loop voltage applied to boundary psi only
inertia	1	1: Include $\mathbf{u} \cdot \nabla \mathbf{u}$ terms
itwofluid	1	1: Include $\mathbf{J} \times \mathbf{B}$ and $\nabla p_e$ terms in Ohm's law (electron form).
		2: ion form (not recommended) 3: parallel pressure gradient
		in Ohm's law only (not recommended)
ibootstrap	0	1: include bootstrap current
ibootstrap_model	0	1: $J_BS = alpha F jp, psi; B$
bootstrap_alpha	0	alpha parameter in bootstrap current model
imp_bf	0	1: include implicit equation for f (recommended for 3D and
		2D complex)
nosig	0	1: drop sigma terms from momentum equation
itor	0	0: cartesian; 1: cylindrical
istatic	0	1: Do not advance velocity
iestatic	0	1: Do not advance magnetic fields
chiiner	1.	factor to multiply the chi equation inertial terms
$ieq_bdotgradt$	1.	1: include equilibrium parallel T gradient
no_vdg_T	0	1: do notinclude V dot grad T in Temp equation (debug)
iwall_is_limiter	1	1: wall acts as limiter
kinetic	0	1: Use kinetic PIC for hot pressure, 2: Incompressible CGL,
		3. Full CGL

Option	Default	Description
iadiabat	1	1: Corrects several problems with itemp=1 option
irunaway	0	1: include runaway electron model
cre	0	runaway speed
imp_temp	0	0: compute temperatures for isplitstep=0, itemp=0
iohmic_heating	1	1: Include Ohmic heating term in heating
irad_heating	1	1: Include radiation heat sink
gravr	0	gravitational acceleration in R-direction
gravz	0	gravitational acceleration in Z-direction

# 10.2 Initial Conditions Options

Option	Default	Description
itaylor	0	Pre-defined initial conditions.
		for itor=1 (toroidal geometry) $($
		0: Tilting cylinder
		1: Calls Grad-Shafranov solver
		2: magneto-rotational equilibrium
		3: rotational instability
		40: Fixed boundary stellarator
		41: Free boundary stellarator
		for itor=0 (slab geometry)
		0: Tilting cylinder
		1: Taylor Reconnection
		2: Force-Free equilibrium (Taylor
		state)
		3: GEM Reconnection
		4: Wave Propagation
		5: Gravitational Instability
		6: Strauss equilibrium
		7: circular field init
		8,9: biharmonic
		10,11,12,13:: analytic RWM test
		problem
		14: 3D wave test
		15: 3D diffusion test
		16: FRS cylindrical equilibrium
		17: ftz init
		18: eigen init
		19: ASDEX profiles similar to YU's
		20: kstar profiles with multiple $q=1$
		surfaces
		21,22: fixed $q(r)$ and $p(r)$ profiles
		23: Startsev equilibrium with
		$J = (1/R_0q_0)(1-r^2)^{-1}$
		27: cylindrical test problem
		29: basicJ profiles
iupstream	0	1: adds diffusion term to convection-like upstream differenc-
_		ing
magus	5.e-2	magnitude of the upstream diffusion term
iflip	0	1: Flip coordinate system handedness
iflip_b	0	1: Flip sign of toroidal field
iflip_j	0	1: Flip sign of toroidal current
iflip_v	0	1: Flip sign of toroidal velocity
iflip_z	0	1: Flip equilibrium across $z=0$ plane

Option	Default	Description
icsym	0	Symmetry of random perturbations
		0: No symmetry
		1: Odd up-down symmetry (in $U$ )
		2: Even up-down symmetry (in $U$ )
bzero	1	$B_{arphi}$ at rzero
bx0	0	Initial field in x-direction for some test problems
vzero	0	Initial toroidal velocity for some test problems
phizero	0	Initial velocity stream function for some test problems
v0_cyl	0	Central toroidal velocity for some test problems
v1_cyl	0	$VZ=v0_cyl + v1_cyl^*psin^{**}beta$
idevice	0	define coils for a particular device
		-1: reads coil.dat file
		0: generic dipole configuration
		1: CDX-U
		2: NSTX
		3: ITER
		4: DIII-D
iwave	0	defines what wave to initialize in wave propagation test
eps	0.01	Size of random perturbation
maxn	200	Maximum wavenumber of initial random noise
verzero	0	magnitude of initial vertical velocity
irmp	0	1: apply nonaxisymmetric fields
		throughout plasma. reads
		$rmp\_coil.dat$ for $(R,Z)$ of window
		$pane coils. reads rmp\_current.dat$
		for $(+-)$ currents in kA and phases
		in degrees. toroidal mode number of
		current specified by ntor. requires
		$\texttt{type\_ext\_field} = 0.$
		2: apply nonaxisymmetric fields only
	1	at boundaries.
type_ext_field	-1	External field type 0: RMP or error field for tokamak ge-
		ometry. 1: For free boundary stellara-
rmp attan	0	tor only: FIELDLINES or MGRID.
	0	$\frac{1}{1}$ irmp-2
irond ovt field	0	1 mp-2
heta	0	narameter used in some model equilibrium initializations
In	0	length scale parameter used in some model equilibrium
111 alongation	1	alongation used in Solovov equilibrium
erougarion	T	ciongation used in polovev equilibrium

Option	Default	Description
isample_ext_field	1	factor to down sample external field data toroidally
isample_ext_field_pol	1	factor do down sample external field data poloidally
scale_ext_field	1	factor to scale external field
shift_ext_field	0	toroidal shift (in deg) of external fields
ibasicj_solvep	0	0: uniform p, solve for F; 1: uniformF, solve for p
basicj_nu	1	exponent in basicj equilibrium
basicj_j0	1	On-axis current density in basicj equilibrium
basicj_voff	1	Radial extent of flat toroidal rotation in basicj equilibrium
basicj_vdelt	1	Width of velocity drop-off, as fraction of ln, in basicj equilib-
		rium
basicj_dexp	1	parameter for basicj equilibrium
basicj_dvac	1	parameter for basicj equilibrium
basicj_q0	0	parameter for basicj equilibrium
basicj_qa	0	parameter for basicj equilibrium
pf_shift	0	(array) horizontal shift of PF coil
pf_shift_angle	0	(array) direction of PF shift in degrees
pf_tilt	0	(array) Angle of PF from vertical in degrees
pf_tilt_angle	0	(array) Axis of rotation for PF tilt in degrees
tf_shift	0	horizontal shift of TF coils
tf_shift_angle	0	direction of TF shift in degrees
tf_tilt	0	angle of TF from vertical in degrees
$tf_tilt_angle$	0	axis of rotation for TF tilt in degrees

Option	Default	Description
inumgs	0	1: Use numerical def. of p and g from profile-p and profile-g
		files
igs	80	Max number of Picard iterations
eta_gs	1000.	factor for smoothing nonaxisymmetries in psi in 3D GS solve
igs_pp_ffp_rescale	0	1: rescale p' and FF' to match p and F
nv1equ	0	1:use numvar $=1$ equilibrium for numvar .GT. 1
tcuro	1	(scaled) plasma current in GS equilibrium
xmag	1	<i>R</i> -coordinate of magnetic axis
zmag	0	Z-coordinate of magnetic axis
xmag0	0	if nonzero, target magnetic axis $R$ for feedback
zmag0	0	if nonzero, target magnetic axis $Z$ for feecback
xlim	0	R-coordinate of limiter
zlim	0	Z-coordinate of limiter
xlim2	0	R-coordinate of limiter $#2$
zlim2	0	Z-coordinate of limiter #2
rzero	1	nominal major radius of device for itor=1
libetap	1.2	approximate value of $l_i/2 + \beta_P$ for free-boundary equ
p0	0.01	Pressure at magnetic axis
pi0	0.005	Ion pressure at magnetic axis
p1	0	$p'(\Psi)$ at magnetic axis
p2	0	$p''(\Psi)$ at magnetic axis
pedge	0	Pressure in vacuum region
tedge	0	temperature in vacuum region (if .GT. 0). Only used in $GS$
		solve. Boundary value of electron temp is $twall = pedge \times$
		$pefac/den\_edge$
tiedge	0	ion temperature in vacuum region
expn	0	Fraction of pressure gradient due to density gradient: $n =$
		$p^{expn}$ .
q0	1	Safety factor at magnetic axis
djdpsi	0	$J'_{\varphi}(\Psi)$ at magnetic axis
th_gs	0.8	implicitness of GS Picard iterations
tol_gs	$10^{-8}$	convergence criteria for GS iteration
pscale	1.	factor multiplying pressure profile
bscale	1.0	Factor multipying toroidal field
bpscale	1.0	Factor multiplying F' (keeping F0 constant)
vscale	1.0	Factor multiplying toroidal rotation profile
iread_bscale	0	1: read profile_bscale for factor to scale F
${\tt iread\_pscale}$	0	1: read profile_pscale for factor to scale $p$ and $p'$

# 10.3 Grad-Shafranov Solver Options

Option	Default	Description
batemanscale	1	Bateman scale the TF, keeping curent profile fixed
irot	0	1: include toroidal rotation in equilibrium calculatin
iscale_rot_by_p	1	see below
alpha0	0	$\alpha_0$ in analytic rotation profile
alpha1	0	$\alpha_1$ in analytic rotation profile
alpha2	0	$\alpha_2$ in analytic rotation profile
alpha3	0	$\alpha_3$ in analytic rotation profile

For iread\_omega=0, the function  $\alpha(\psi)$  is parameterized by:

 $\tilde{\alpha} = \alpha_0 + \alpha_1 s + \alpha_2 s^2 + \alpha_3 s^3$ 

For iscale\_rot\_by\_p = 0:  $\alpha = \tilde{\alpha} \times n(\psi)/p(\psi)$ . For iscale\_rot\_by\_p = 1:  $\alpha = \tilde{\alpha}$ . For iscale\_rot\_by\_p = 2:  $\alpha = \left[\alpha_0 + \alpha_1 e^{-[(\psi - \alpha_2)/\alpha_3]^2}\right] \times n(\psi)/p(\psi)$ In all cases, the angular velocity is then determined by:

$$\omega = \left[\frac{2\alpha p(\psi)}{R_0^2 n(\psi)}\right]^{\frac{1}{2}}$$

idenfunc	0	0: $n = \text{den0} \times (p/p0)^{\text{expn}} + \text{denedge}$
		1: $n = \operatorname{den}0 \times \frac{1}{2} \times \left[1 + \operatorname{tanh}\left(\frac{\psi - (\psi_B + n_O(\psi_B - \psi_M))}{\Delta \times (\psi_B - \psi_M)}\right)\right]$
		2: $n = \text{den}0 + \frac{1}{2} \left( \text{den}_{-}\text{edge} - \text{den}0 \right) \times \left[ 1 + \tanh\left(\frac{\tilde{\psi} - n_O}{\Delta}\right) \right]$
		3: if $\tilde{\psi}$ .LT. $n_O$ and $(\psi - \psi_M) \times$
		$[d\psi/dx(x-x_{MA})+d\psi/dz(z-z_{MA})]$ .GT. 0, then $n =$
		den0. Else, $n = \text{den}_{-}\text{edge}$ .
		$(\psi_B = \text{psibound}, \psi_M = \text{psimin}, \tilde{\psi} = (\psi - \psi_M)/(\psi_B - \psi_M))$
den_edge	0.0	edge density. If 0, set to $den0^*(pedge/p0)^{**}expn$
den0	1.0	(scaled) central density
denoff	1.0	$n_O$ : offset for idenfunc=1,2,3
dendelt	0.1	$\Delta$ : width of transition region for idenfunc=1,2
divertors	0	Number of divertors $(0-2)$
divcur	0.1	Divertor $\operatorname{current}(s)$ , as fraction of teuro
xdiv	0	r-coordinate of divertor current(s)
zdiv	0	z-coordinate of divertor current. If $divertors = 2$ , the sec-
		ond divertor has $z = -zdiv$ .
xnull	0	Guess for $r$ -coordinate of x-point
znull	0	Guess for $z$ -coordinate of x-point
mod_null_rs	0	if 1: you can reset xnull and znull from C1input
xnull0	0	Target R-Coordinate of x-point for feedback
znullO	0	Target Z-Coordinate of x-point for feedback
xnull2	0	Guess for $r$ -coordinate of inactive x-point
znull2	0	Guess for $z$ -coordinate of inactive x-point
mod_null_rs2	0	if1: you can reset xnull2 and znull2 from C1input

Option	Default	Description
gs_pf_psi_width	0	width of psi smoothing into provate flux region
$gs\_vertical\_feedback$	0	proportional feedback of each coil to (zmag-zmag0) (array)
gs_vertical_feedback_i	0	integral feedback of each coil to (zmag-zmag0) (array)
gs_vertical_feedback_x	0	proportional feedback of each coil to (znull-znull0) (array)
gs_vertical_feedback_x_i	0	integral feedback of each coil to (znull-znull0) (array)
gs_radial_feedback	0	proportional feedback of each coil to (xmag-xmag0) (array)
gs_radial_feedback_i	0	integral feedback of each coil to (xmag-xmag0) (array)
gs_radial_feedback_x	0	proportional feedback of each coil to (xnull-xnull0) (array)
gs_radial_feedback_x_i	0	integral feedback of each coil to (xnull-xnull0) (array)
igs_extend_p	0	1: extend p past pls=1 using ne and Te profiles
igs_feedfac	1	proportionality factor for external field feedback
igs_forcefree_lcfs	-1	1: ensure that GS solution is force free at LCFS
igs_start_xpoint_search	0	number of GS iterations before searching for x-point
sigma0	0	width of Gaussian for initial current distribution for GS iter-
		ation
igs_extend_diagmag	1	1: extend diamagnetic rotation past $psi=1$
adapt_qs	0	Safety factor values to pack around (array)
adapt_zlow	0	Z-coordinate below which SOL adaption is coarse
adapt_zup	0	Z-coordinate above which SOL adaptation is coarse

Option	Default	Description
ivisfunc	0	select viscosity function
		0: visc = amu
		1: visc = amu + $\frac{1}{2}$ amu_edge × $\left[1. + \tanh\left[\frac{\psi - (\psi_l + \nu_0(\psi_l - \psi_0))}{\nu_\Delta}\right]\right]$
		2: visc = amu + $\frac{1}{2}$ amu_edge × $\left[1. + \tanh\left[\frac{\psi - \nu_0}{\nu_\Delta}\right]\right]$
		or, if amuoff2 .ne. 0 and amudelt2.ne.0)
		$\operatorname{visc} = \operatorname{amu} + \frac{1}{4}\operatorname{amu}_{\operatorname{edge}} \times \left[2. + \operatorname{tanh}\left[\frac{\tilde{\psi} - \nu_0}{\nu_{\Delta}}\right] + \operatorname{tanh}\left[\frac{\tilde{\psi} - \nu_{02}}{\nu_{\Delta 2}}\right]\right]$
		3: visc = amu or amu_edge depending on criteria in de-
		fine_fields
amu	0	core viscosity for ivisfunc $=0,,3$
amu_edge	0	edge viscosity for ivisfunc $= 1,,3$
amuoff	0	$\nu_0$ in ivisfunc = 1,2
amuoff2		$\nu_{02}$ in ivisfunc = 1,2
amudelt	0	$\nu_{\Delta}$ in ivisfunc = 1,2
amudelt2	0	$\nu_{\Delta 2}$ in ivisfunc = 1,2
amuc	0	Compressional viscosity coefficient
amupar	0	Parallel viscosity coefficient
amue	0	bootstrap viscosity coefficient
iresfunc	0	select resistivity function
		0: $eta = etar + eta0/Te^{**}(3/2)$
		1: eta = etar + $\frac{1}{2}$ eta $0 \times \left[1. + \tanh\left[\frac{\psi - (\psi_l + \text{etaoff} \times (\psi_l - \psi_0))}{\text{etadelt} \times (\psi_l - \psi_0)}\right]\right]$
		2: $\operatorname{eta} = \operatorname{etar} + \frac{1}{2}\operatorname{eta} \times \left[1. + \tanh\left[\frac{\varphi - \operatorname{eta}}{\operatorname{eta}}\right]\right]$
		The following two options are applied in a way that they
		should not have negative valueseven if the idl plots
		indicate otherwise
		3: eta = etar for $\psi$ < etaoff othrwise eta0
		4: Spitzer resistivity with offset.
		Define $T_{wall} = \text{pedge}^* \text{pefac}/\text{den}_e \text{dge}$
		for $T_e > T_{wall} - T_e^{off}, \eta = (T_e - T_e^{off})^{-3/2}$
		for $T_e < T_{wall} - T_e^{off}$ , $\eta = (T_{wall} - T_e^{off})^{-3/2}$
		can be increased by inputing $\text{eta}_{\text{fac}} > 1$ .
		5: Simple neoclassical model:
		$\eta = \mathrm{eta0} \times (n_e/p_e)^{3/2}/(1 1.46(r/R)^{1/2})$
etar	0	see description of iresfunc
eta0	0	see description of iresfunc
etaoff	0	see description of iresfunc
etadelt	0	see description of iresfunc
eta_te_offset	0	$T_e^{off}$ for iresfunc=4
ikprad_te_offset	0	if 1, $T_e^{off}$ also used in kprad and ablation
eta_fac	0	for iresfunc=4, resistivity multiplied by eta_fac
eta_mod	0	if 1: remove d/dphi terms in resistivity

## 10.4 Transport Coefficients

Option	Default	Description
eta_max	0	maximum resistivity in plasma (defaults to etavac)
eta_min	0	minimum resistivity in plasma
ikappafunc	0	select thermal conductivity function
		0: $\kappa = \text{kappat} + \text{kappa0} \times (n^3/p)^{1/2}$
		1: $\kappa = \text{kappa0} \times \frac{1}{2} \left[ 1 + \tanh \left[ \frac{\psi - \left(\psi_l + \kappa^{0ff} \times (\psi_l - \psi_0)\right)}{\kappa_{\Delta} \times (\psi_l - \psi_0)} \right] \right]$
		2: $\kappa = \text{kappa0} \times \frac{1}{2} \left[ 1 + \tanh \frac{\tilde{\psi} - \kappa^{off}}{\kappa_{\Delta}} \right] \text{for} \tilde{\psi} < 1$
		2: $\kappa = \text{kappa0} \times \frac{1}{2} \left[ 1 + \tanh \frac{2 - \psi - \kappa^{0/2}}{\kappa_{\Delta}} \right] \text{ for } \psi > 1$
		3: $\kappa = \text{kappat} + \text{kappa0} \times 1/(pn)^{1/2}$
		4: $\kappa = \text{kappad} + \text{kappad} \times (1. + \text{kappadelt} \times  \nabla T_e ^2)$
		5: $\kappa = \text{kappat} + \text{kappa}/T_e$ limited by kappa_max
		10: read from profile_kappa file in $m^2$ /sec
		11: read from profile_kappa file in normalized units
	0	12: option to go with itaylor= $27$
kappa_max	0	it .NE. 0, max $\kappa$ for ikappatunc=5
kappai_fac	1	ion thermal conduction is kappai_fac* kappa
ikapscale	0	if 1: kappar gets scaled by kappa
ikappar_ni	0	1: include 1/n terms in parallel heat flux
kappaoff	0	$\kappa^{off}$ see ikappafunc
kappadelt	0	$\kappa_{\Delta}$ see ikappafunc
kappat	0	isotropic thermal conductivity
kappa0	0	see ikappafunc
ikapparfunc	0	select parallel thermal conductivity (PTC) function 0: PTC = kappar $= 2\pi T^{2}$
_	0	1: PTC = kappar/ $[(T_{crit}/T)^{5/2} + 1]$
kappar	0	Parallel thermal conductivity
tcrit	0	$T_{crit}$ for ikappartunc = 1
kapparı_fac	1	ion parallel thermal conductivity is kapparifac x electron value
kappax	0	coefficient of $B \times \nabla T$ temperature diffusion
kappah	0	if nonzero, kappa = kappah × $\tanh^2 \left[ (\tilde{\psi} - 1.)/2 \right]$
kappaf	1	Factor multiplying kappa when $\nabla p < \nabla p_{crit}$
kappag	0	Thermal diffusion proportional to pressure gradient
$gradp_crit$	0	$ abla p_{crit}$ for kappaf,kappag model
k_fac	1	Factor by which TF is multiplied in denominator of $kappa_par$
temin_q0	0	Min temperature used in equipartition for ipres=1
idenmfunc	0	selects from of particle diffusion (PD)
		0: PD = denm
		1: $PD = denm + denmt/Te$
		10: read from file profile_denm in $m^2/\text{sec}$
		11: read from file profile_denm in normalized units
denm	0	see idenmfunc
denmt	0	multiplier of $1/\text{Te}$ for idenmfunc=1
denmmin	0	minimum value of denm
denmmax	$1.\mathrm{E6}$	maximum value of denm

# 10.5 Hyper-Diffusivity

Option	Default	Description
imp_hyper	0	switch for evaluating hyper resistivity
		0: $\lambda_H \nabla^2 \mathbf{J}$ explicit for $\psi$ implicit for F
		1: $\lambda_H \nabla^2 \mathbf{J}$ implicit for $\psi$ implicit for F 2: $(\mathbf{B}/B^2) \nabla \bullet \lambda_H \nabla \sigma$ implicit for $\psi$ and F ( $\sigma = \mathbf{J} \bullet \mathbf{B}/B^2$ )
deex	1	scale length used in hyper
hyper	0	hyper coefficient for $\psi$ equation
hyperc	0	hyper coefficient for poloidal velocity
hyperi	0	hyper coefficient for toroidal field
hyperp	0	hyper coefficient for pressure
hyperv	0	hyper coefficient for toroidal flow
ihypdx	2	hyper terms multiplied by deex <sup>**</sup> ihypdx
ihypeta	1	swithc for multipliers of hyper terms
	1	magnetic field hyper multiplied by eta
	2	magnetic field hyper multiplied by p
ihypamu	1	1: velocity hyper coefficient multiplied by amu
ihypkappa	1	1: pressure hyper coefficient multiplied by kappa

## 10.6 Unit Normalizations

Option	Default	Description
n0_norm	$10^{14}$	Density normalization (in cgs)
b0_norm	$10^{4}$	Magnetic field normalization (in cgs)
10_norm	100	Length normalization (in cgs)

Option	Default	Description
isurface	1	include surface terms in Galerkin method
icurv	2	if $> 0$ , include curvature from mesh
nonrect	0	1: non-rectangular boundary
ifixedb	0	Set $\psi = 0$ on boundary
inonormalflow	1	1: No-normal-flow boundary
inoslip_pol	0	1: No-slip boundaries for poloidal velocity
$inoslip_tor$	1	1: No-slip boundaries for toroidal velocity
inostress_tor	0	1: No-normal-stress boundary for toroidal velocity
iconst_bz	1	1: Toroidal field held constant on boundary
iconst_bn	1	1: Hold normal field constant on boundary
iconst_n	0	1: Density held constant on boundary
iconst_p	1	1: Pressure held constant on boundary
iconst_t	0	1: Temperature held constant on boundary
inograd_p	0	1: No normal pressure gradient
$inograd_t$	0	1: No normal temperature gradient
inograd_n	0	1: No normal density gradient
com_bc	0	1: $ abla^2 \chi = 0$
vor_bc	0	1: $\Delta^* U = 0$
$inocurrent_pol$	0	1: no poloidal current on boundary
inocurrent_tor	0	1: no toroidal current on boundary
inocurrent_norm	0	1: no normal current on boundary
ifbound	-1	boundary condition on $f'$
		1: Dirichlet
		2: Neumann
iconstflux	0	1: conserves toroidal flux in nonlinear calculation
tebound	-1	boundary condition for electron temperature
tibound	-1	boundary condition for ion temperature
iper	0	1: Left/right boundaries periodic
jper	0	2: Top/bottom boundaries periodic

## 10.7 Boundary Conditions

# 10.8 Time Step

Option	$\mathbf{Default}$	Description
dt	0.1	Initial size of ime step
ntimemax	20	Total number of time steps
integrator	0	0: Crank-Nicholson (CN); 1: BDF2
imp_mod	0	0: $\theta$ -implicit
		1: Implicit leapfrof (isplitstep = 1 only)
thimp	0.5	Implicitness parameter
thimpsm	1	Implicitness of the smoother functions
isplitstep	1	0: Fully implicit time step; 1: split time step.
iteratephi	0	1: Calculate transport coefficients after field advance, then
		recalculate field advance.
idiff	0	1: solve for difference between n and $n+1$ in B,p
idifv	0	1: solve for difference between n and $n+1$ for V
$irecalc_eta$	0	1:recalculate transport coefficients after density solve
iconst_eta	0	1" don't evolve resistivity
$\texttt{itime\_independent}$	0	1:exclude d/dt terms
harned_mikic	0	coefficient of Harned-Mikic 2F stabilization term
isources	0	1: include source terms in velocity advance
nskip	1	number of time steps per matrix recalculation
pskip	1	number of timesteps the preconditioner is reused
iskippc	1	number of times preconditioner is reused
ddt	0	change in timestep per timestep
frequency	0	frequency in time-independent calculation
dtmin	4.0	minimum timestep for variable timestep calculation
dtmax	40.	maximum timestep for variable timestep calculation
dtkecrit	0	lower timestep if ekin es above this $(0.01 \text{ typical})$
dtfrac	.10	max fractional change of timestep in 1 cycle
max_repeat	3	max number time step is repeated for ksp_max iterations exceeded
ksp_max	10000	max number of ksp iterations before repeating time step
ksp_min	1200	increase dt if ksp ; ksp_min
ksp_warn	1600	decrease dt if ksp ¿ ksp_warn

### 10.9 Mesh

Option	Default	Description
nplanes	1	number of toroidal planes for 3D nonlinear
xzero	0	R-coordinate of bottom left corner of domain
zzero	0	z-coordinate of bottom left corder of domain
tiltangled	0	angle a rectangular mesh is tilted
mesh_model		model file name from which the mesh is generated
mesh_filename		mesh name of .smb files (without number)
imatassemble	0	1: use petsc matrix parallel assembly instead of scorec
$imulti\_region$	0	1: Mesh has multiple regions that include resistive wall and
		vacuum. Wall resistivity is "eta_wall", vacuum resistivity is
		"eta_vac"
$toroidal\_pack\_angle$	0	toroidal angle of maximum mesh packing
$toroidal\_pack\_factor$	1	ratio of longest to shortest toroidal element
iread_vmec	0	1: read geometry from VMEC file
vmec_filename	geometry.nc	name of vmec data file
nperiods	1	model $1/nperiods$ of a torus when ifull_torus=0
$iread_planes$	0	Read positions of toroidal planes from plane_positions
$bloat_distance$	0	factor to expand VMEC domain
$bloat_factor$	0	factor to expand VMEC domain
$ifull_torus$	0	0: one field period, 1: full torus
igeometry	0	0: default, identity
xcenter	0	center of logical mesh $(\mathbf{x})$
zcenter	0	center of logical mesh $(z)$
bound_type(i)	0	Boundary conditions to apply on mesh loop i. $0 = None, 1$
		= First wall, $2 =$ Domain boundary
<pre>zone_type(i)</pre>	0	Physics model of mesh zone i. $1 = $ plasma, $2 = $ conductor,
		3 = vacuum.

### 10.10 Solver

Option	Default	Description
solver_type	0	for PETSc only, 0: direct solve, 1: iterative solver. for Trilin-
		ios, iterative solver is used
solver_tol	1.e-9	solver tolerance
num_iter	100	maximum number of iterations

## 10.11 Mesh Adaptation (will be depricated soon)

Option	Default	Description
iadapt	0	0: no adaptation
		1:adapt mesh from the flux field in equilibrium

# 10.12 Numerical Options

Option	Default	Description
int_pts_main	25	Sampling points for integrations in main time step matrices
int_pts_aux	25	Sampling points for integrations in calculations of auxiliary variables
int_pts_diag	25	Sampling points for integrations in diagnostic calculations
int_pts_tor	5	Max number of toroidal integration points
ivform	0	0: $\mathbf{u} = \nabla U \times \nabla \varphi + V \nabla \varphi + \nabla \chi$
		1: $\mathbf{u} = R^2 \nabla U \times \nabla \varphi + R^2 \omega \nabla \varphi + R^{-2} \nabla \chi$
		Now depricated to ivform=1 only
jadv	0	1: Use toroidal current density equation instead of poloidal
		flux equation.
max_ke	1.0	Maximum value of kinetic energy before solution is rescaled
		in linear simulations. $(0 = \text{don't rescale})$
equilibrate	0	1: scale trial function so L2 norm $= 1$
regular	0	regularization constant in chi equation
iset_pe_floor	0	1: do not let pe drop below pe_floor
pe_floor	0	minimum value for pe when iset_pe_floor=1
iset_pi_floor	0	1: do not let pi drop below pe_floor
pi_floor	0	minimum value for pi when iset_pi_floor=1
iset_te_floor	0	1: do not let te drop below te_floor
te_floor	0	minimum value for te when iset_te_floor=1
iset_ti_floor	0	1: do not let ti drop below ti_floor
ti_floor	0	minimum value for ti when iset_ti_floor=1
iset_ne_floor	0	1: do not let ne drop below ne_floor
ne_floor	0	minimum value for ne when iset_ne_floor=1
iset_ni_floor	0	1: do not let ni drop below ni_floor
ni_floor	0	minimum value for ni when iset_ni_floor=1
iprecompute_metric	0	1: precompute metric temsor

Option	Default	Description
iread_eqdsk	0	1: Read EFIT g-file 'geqdsk'
		2: Read psi from geqdsk, but use analytic profiles for p and
		$\mathbf{F}$
		3: read profiles from geqdsk, but not the psi
iread_dskbal	0	1: Read BAL file 'dskbal'
iread_jsolver	0	1: Read Jsolver file 'fixed'
$iread_omega$	0	1: reads in rotation profile
iread_omega_ExB	0	1: read ExB rotation
iread_omega_e	0	1: read electron rotation
iread_ne	0	1: read in electron density profile
iread_te	0	1: read in temperature profile
iread_p	0	1: read pressure profile from profile_p
iread_neo	0	1: read velocity profiles from NEO output
ineo_subtract_diamag	0	1: subtract diamagnetic term from input vel when reading
		neo velocity
iread_heatsource	0	1: read heat source profile (psi normalized) scaled by ghs_rate
${\tt iread\_partilesource}$	0	1: read particle source profile (psi normalized) scaled with
		pellet_rate
iread_f	0	1: read R x BT from file
iread_j	0	1: read current density from a file

# 10.13 Input Options

Option	Default	Description
ntimepr	5	Number of timesteps per full field output
iprint	0	1: Print detailed info to stdout
ifout	-1	1: ourput f field
idouble_out	0	1: use double precision floating in putput hdf5 files
itemp_plot	0	1: output vdotgradt, deldotq_perp, deldotq_par, eta_jsq
ibdgp	0	optons for plotting partial terms for bdgp plot
		(1) $[\psi, \Phi], (2) (f', \Phi), (3) - R^{-2}F\Phi'$
iveldif	0	option for plotting partial terms for veldif plot
		(1) $ /psi, U $ , (2) $R^{-3}(\psi, chi) + R^{-2} \chi, f' $
		$(3)R^{-1}\Phi'(4) R^{-1}\Phi' + R^{-1}FU'$
		(5) $ \psi, U  + R(U, f') - R^{-1}FU'$ (6) $R^{-2} /chi, f' $
icalc_scalars	1	1 : calculate scalar diagnostics
ike_only	0	1 " only calculate ke scalar diagnostic
ike_harmonics	0	number of toroidal harmonics of kinetic energy to be calcu-
		lated for diagnostics
ibh_harmonics	0	number of toroidal harmonics of magnetic energy to be cal-
		culated for diagnostics
irestart	0	0: start from timestep 0
		1: normal restart (can restart 3D from 2D)
		3: start a 2D complex run from a 2D real restart
$irestart_slice$	-1	if set to an integer, restart from that time slice
itimer	0	1: output internal timeing data
iwrite_transport_coefs	1	1: output transport coefficients fields
iwrite_aux_vars	1	1:output auxiliary variable fields

# 10.14 Output Options

Option	Default	Description
xray_detector_enabled	0	1: enable xray detector
xray_r0	0	R coordinate of xray detector
xray_phi0	0	phi coordinate of xray detector
xray_z0	0	Z coordinate of xray detector
xray_theta	0	angle of xray detector chord (degrees)
xray_sigma	1	spread of xray detector chord (degrees)
imag_probes	0	number of magnetic probes
<pre>mag_probe_x(i)</pre>	0	R coordinate of magnetic probe i
<pre>mag_probe_phi(i)</pre>	0	phi coordinate of magnetic probe i
<pre>mag_probe_z(i)</pre>	0	Z coordinate of magnetic probe i
<pre>mag_probe_nx(i)</pre>	0	R component of normal vector of mag probe i
<pre>mag_probe_nphi(i)</pre>	0	phi component of normal vector of mag probe i
<pre>mag_probe_nz(i)</pre>	0	Z component of normal vector of mag probe i
iflux_loops	0	number of flux loops
flux_loop_x(i)	0	R coordinate of flux loop i
flux_loop_z(i)	0	Z coordinate of flux loop i
$ifixed_temax$	0	1: temax evaluated at (xmag0,0,zmag0)

## 10.15 Diagnostics

# 10.16 Sources/Sinks

Option	Default	Description	
ibeam	0	neutral beam source	
		1:include neutral beam particle, energy, and momentum source:	
		$S = \frac{nb_n}{4r\pi^2 nb_{dr}} \exp \left[ (r - nb_r)^2 + (z - nb_z)^2 \right] / 2nb_{dr}^2$	
		2: include only particle and energy source	
		3: include only energy source	
		4: include only momentum and energy source	
		5: include only momentum source	
beam_x	0	R coordinate of beam center (in m)	
beam_z	0	Z coordinate of beam center (in m)	
beam_v	1.e4	beam voltage (in volts)	
beam_rate	0	ions/second deposited by beam	
beam_dr	.1	dispersion of beam deposition	
beam_dv	100.	dispersion of beam voltage (in volts)	
beam_fracpar	1	cosine of beam angle relative to parallel (for momentum source)	
vloop	0	initial loop voltage, NOTE: to change vloop at restart time, must	
		have control_type = $-1$	
tcur	0	target (scaled) plasma current for current control: $\mu_0 I_P$	
tcuri	0	if tcuri .ne. tcurf, the target current is a function of time	
tcurf	0	$tcur = tcuri + (tcurf-tcuri) \times .5 \times (1. + tanh((t-tcur_t0)/tcur_tw))$	
$tcur_t0$	0	see tcurf	
tcur_tw	0	see tcurf	
control_type	-1.	0 old current control algorithm (not recommended)	
		1: standard PID control with the following control parameters	
$control_p$	0	proportional control coefficient	
control_i	0	integral control coefficient	
control_d	0	derivative control coefficient	
ipellet	0	density source if non-zero (3D part equals 1 for 2D runs). Double-	
		digit values have volume integrals normalized to 1. Make netative	
		for initial perturbation only.	
		Define: $G_{2D} = \frac{1}{2\pi R V_p^2} \exp\left[-\frac{(R-R_P)^2 + (Z-Z_P)^2}{2V_p^2}\right]$	
		1: $S = G_{2D} \times \frac{R}{\sqrt{2\pi}V_{*}} \exp\left[-\frac{RR_{P}(1-\cos(\phi-\phi_{P}))}{V_{*}^{2}}\right]^{2}$	
		2: $S = \text{den0} \times (\max(p, p_{edge})/p_0)^{expn}$ 2D and 3D	
		3: Gaussian source proportional to pressure	
		$S = p \times G_{2D} \times \frac{R}{\sqrt{2\pi}V_p} \exp\left[-\frac{RR_P(1-\cos(\phi-\phi_P))}{V_c^2}\right]$	
		4: Same distribution as ipellet=1 in $3D$	
		$S = \sqrt{2\pi}RV_n \times G_{2D} \times \frac{1}{2\pi} \exp\left[-\frac{RR_P(1-\cos(\phi-\phi_P))}{2\pi}\right]$	
		$V_{t} = 2D - 2\pi V_{p} V_{t} = 1 \begin{bmatrix} V_{t}^{2} \\ V_{t} \end{bmatrix}$	
Option	Default	Description	
-----------------	---------	--	--
		11: Same as $\#1$ but numerically normalized	
		12: Spherical, Cartesian Gaussian numerically normalized	
		$2D: S = RG_{2D}$	
		3D: $S = \exp \left[-\frac{(R\cos\phi - R_p\cos\phi_p)^2 + (R\sin\phi - R_p\sin\phi_p)^2 + (Z - Z_p)^2}{2V^2}\right]$	
		13: Axisymmetric, toroidal Gaussian, numerically normalized 2D and 3D: $S = G_{2D}$ 14: Toroidal distribution is a blend of a von Mises and Cauchy dis-	
		tribution	
		$S = G_{2D} \times \frac{R}{\sqrt{2\pi}V_p} \times$	
		$\left[ (1 - f_c) \exp\left[-\frac{RR_p \left(1 - \cos(\phi - \phi_p)\right)}{V_t^2}\right] + f_c \times \frac{\cosh\left(\frac{V_t}{\sqrt{RR_p}}\right) - \cos(\phi_p)}{\cosh\left(\frac{V_t}{\sqrt{RR_p}}\right) - \cos(\phi - \phi_p)} \right] \right]$	
		15: Toroidal von-Mises distribution with angular half-width	
		$S = G_{2D} \times \exp\left[\cos(\phi - \phi_n)/V_r^2\right]$	
		Pellet_var_tor radians for ipellet=15, distance otherwise	
ipellet_z	0	Atomic number of pellet (0 for main-ion species)	
ipellet_abl	0	Turn on pellet ablation. (Recommend double-digit ipellet for particle	
-		conservation)	
		1: Include ablation model [Parks NF94] calibrated on DIII-D (Li)	
		2: Include new ablation model [Parks, 2015] for small pellets (Li)	
		3: Parks model developed 6/20/2017 (for Neon)	
temin_abl	0	Minimum temperature at which ablation turns on	
$iread_pellet$	0	0: Single pellet defined by scalar parameters below	
		1: Read pellet.dat, with one row per pellet. The 13 space-delimited	
		columns are:	
		pellet_r, pellet_phi, pellet_z, pellet_rate, pellet_var, pellet_var_tor,	
		$pellet\_velr, \ pellet\_vel\_phi, \ pellet\_vel\_z, \ r\_p, \ cloud\_pel, \ pellet\_mix,$	
		cauchy_fraction	
pellet_rate	0	Particle number injection rate	
$pellet_var$	1	Variance of injection profile	
pellet_var_tor	0	toroidal spatial dispersion of pellet source $(V_t)$ . If zero, pellet_var_tor = pellet_var.	
pellet_R	0	R-coordinate of injection profile	
pellet_z	0	Z-coordinate of injection profile	
pellet_velr	0	initial radial velocity of the pellet	
$pellet_velphi$	0	initial toroidal velocity of pellet	
$pellet_velz$	0	initial vertical velocity of pellet	
r_p	1.e-3	initial pellet radius	
cloud_pel	1	Parameter used to change the width of the density source if ablating.	
		In this case, $pellet_var = cloud_pel * r_p$	
pellet_mix	0	Molar fraction of diatomic main ion molecures in pellet (e.g., $D_2$ )	

Option	Default Description		
irestart_pellet	0	1: will read the following from C1input at restart (the rest from	
		*.h5) pellet_rate, pellet_rate_D2, pellet_var_tor, pellet_var, cloud_pel,	
		pellet_mix, cauchy_fraction	
abl_fac	1.0	factor to multiply ablation rate from predefined formula	
$n\_control\_type$	-1	-1: no density control	
		0: old density control algorithm (not recommended)	
		1: Standard PID control with the following parameters:	
$n\_control\_p$	0	proportional feedback constant	
$n\_control\_i$	0	integral feedback constant	
n_control_d	0	derivative feedback constant	
igaussian_heat_source	0	1: include Gaussian heat fource	
ghs_x	0	R coordinate of Gaussian heat source	
ghs_z	0	Z coordinate of Gaussian heat source	
ghs_rate	0	amplitude of Gaussian heat source	
ghs_var	0	variance of Gaussian heat source	
ghs_phi	0	phi coordinate of Gaussian heat source	
ghs_var_tor	0	toroidal variance of GAussian heat source	
ionization	0	1: include neutral ionization source	
ionization_rate	0	Ionization rate coefficient	
ionization_temp 0.01		Ionization energy	
ionization_depth 0.01		Temperature scale-length of neutral burn-out	
current drive source	$\psi = \ldots + \eta$	$(\Delta^*\psi - J_{CD})$	
	$J_{CD} = J_0  \epsilon$	$\exp - \left[ (R - R_{0cd})^2 + (Z - Z_{0cd})^2 \right] / W_{cd} - \Delta_{cd}$	
icd_source 0	1: include current drive		
J_0cd 0	$J_0$ : Magnitude of Gaussian		
$R_{-}$ Ocd $0$	$R_{0cd}$ : R coordinate of maximum		
Z_0cd 0	$Z_{0cd}$ : Z coordinate of maximum		
$W_{-}cd$ 0	$W_{cd}$ : Width of maximum		
delta_cd 0	$\Delta_{cd}$ : shift of Gaussian		
isink 0	number of density sinks: 0,1,or 2		
sink1_x 0	R coordinate of first density sink		
sink1_z 0	Z coordinate of first density sink		
sink1_rate 0	rate of first density sink		
sink1_var 1	variance of first density sink		
sink2_x 0	R coordinate of second density sink		
sink2_z 0	Z coordinate of second density sink		
sink2_rate 0	rate of second density sink		
sink2_var 1	variance of second density sink		
idenfloor 0	1: density in vacuum pegged to den_edge		
alphadenfloor 0	multiplier of (den_edge - den). Must be .lt. 1/DT		
ipforce 0	1: include	poloidal momentum source	
	$F = f(\tilde{\psi}) \nabla$	$ abla \psi  imes  abla \phi$	
	$f(\tilde{\psi}) = \alpha($	$(1- ilde{\psi})^N rac{\delta^2}{( ilde{arphi})^2+\delta^2}$	
	2: Include	Luca Guzzatto form of momentum source	

Option	Default	Description
dforce	0	δ
xforce	0	$\psi_0$
nforce	0	N
aforce	0	$\alpha$
iheat_sink	0	1: special heat sink for itaylor=27
coolrate	0	$S = coolrate^*(pedge-p)$ for iheat_sink = 1
iarc_source	0	1: density source due to halo current
arc_source_alpha	0	parameter for arc_source
$arc_source_eta$	.01	parameter for arc_source

# 10.17 Resistive Wall

Option	Default	Description
eta_vac	1	resistivity of vacuum region
eta_wall	.001	resistivity of conducting wall regions
eta_wallRZ	.001	poloidal resistivity of wall region (if different from eta_wall)
iwall_break	0	number of wall breaks
eta_break	1	resistivity of wall break (array)
wall_break_phimax	0	max phi coordinate for break (array)
wall_break_phimin	0	min phi coordinate for break (array)
wall_break_xmax	0	max R coordinate for break (array)
wall_break_xmin	0	min R coordinate for break (array)
wall_break_zmax	0	max Z coordinate for break (array)
wall_break_zmin	0	min Z coordinate for break (array)
iwall_regions	0	number of resistive wall regions
<pre>wall_region_eta()</pre>	1.e-3	resistivity of each wall region
$wall_region_etaRZ()$	1.e-3	poloidal restivity (if different from wall_region_eta)
<pre>wall_region_filename()</pre>		file name will wall contour points
eta_zone(i)	eta_wall	Resistivity of mesh region i. Only applies if $zone_type(i) = 2$ (conductor).
etaRZ_zone(i)	$eta_zone(i)$	Poloidal resistivity of mesh region i. Only applies if $zone_type(i) = 2$ (conductor).
eta_rekc	0	resistivity of runaway electron killer coil (REKC)
ntor_rekc	0	toroidal mode number of REKC
mpol_rekc	0	poloidal mode number of REKC
phi_rekc	0	toroidal angle of fixed point of REKC
theta_rekc	0	poloidal angle of fixed point of REKC
rzero_rekc	0	R0 for computing theta of REKC
zzero_rekc	0	Z0 for computing theta of REKC
isym_rekc	0	if non-zero, coil is double helix with $(+,-)$ mpol_rekc

Option	Default	Description	
( gam	5/3	ratio of sepcific heatx	
db	0	ion skin depth (overrides db_fac	
db_fac	0	factor multiplying physical value of ion skin depth	
mass_ratio	0	ratio of ion to electron mass	
lambdae	0	lambdae	
z_ion	1	Z-effective	
ion_mass	1	ion mass in units of m_p	
lambda_coulomb	17	Couloumb logarithm	
thermal_force_coeff	0	coefficient of thermal force	
ntor	0	toroidal mode number for 3D linear (complex)	
mpol	0	poloidal mode number for certain test problem initializations	

# 10.18 Miscellaneous

# 10.19 Deprecated

Option	Default	Description
ipartitioned	0	
igs_method	-1	
ibform	-1	
delta_wall	1.	

# 10.20 Trilinos Options

Option	Default	Description
drop_tolerance	0	ILU drop tolerance
graph_fill	0	graph fill level
ilu_fill_level	1	ILU fill level
ilu_omega	1	relaxation parameter for rILU
krylov_solver	gmres	Krylov solver
poly_ord	1	polynomial order for certain perconditioners
preconditioner	$dom_decomp$	preconditioner
<pre>sub_dom_solver</pre>	ilu	subdomain solver in preconditioner
${\tt subdomain\_overlap}$	1	subdomain overlap

# 10.21 Simple Radiation Model

Option	Default	Description	
iprad	0	1: call Prad radiation module with one impurity species	
		$P_{rad} = n_e n_D L_D(T_e) + n_e n_Z L_Z(T_e)$	
		Cooling rate of deuterium is $L_D = 5.35 \times 10^{-37} T_e^{1/2} [keV]W - m$	
		$L_Z(T_e)$ taken from Post, et al, Atomic data and nuclear data tables,	
		<b>20</b> pp. 397-439 (1977)	
prad_fz	1	density of impurity species as fraction of $n_e$	
prad_z	1	Z of impurity species: $Z=6(C)$ , $18(Argon)$ , $26(Fe)$ are available	
${\tt iread\_prad}$	0	1: Read impurity density from profile_nz (units of $10^{20}/m^3$ )	

# 10.22 KPRAD Radiation Model

Option	Default	Description
ikprad	0	1: KPRAD module with one impurity species
ikprad_z	1	Z of impurity species in KPRAD module. Presently available: 2(He),
		4(Be), 6 (C),10(Ne), 18 (Ar)
kprad_fz	0	density of neutrals as fraction of ne
kprad_nz	0	Density of neutral impurities
kprad_nemin	$10^{-12}$	Minimum normalized electron density for KPRAD evolution
kprad_temin	$2 \times 10^{-7}$	Minimum normalized electron temperature for KPRAD evolution
ikprad_max_dt	0	Set max time step for KPRAD ionization
		0: MHD time step dt
		1: <b>RECOMMENDED:</b> $dt/(kprad_z + 1)$ : ensures evolution
		through all charge states.
ikprad_evolve_internal	0	Update local temperature during KPRAD subcycling:
		0: Te fixed before subcycling
		1: <b>RECOMMENDED</b> Local ne and Te used for KPRAD ioniza-
		tion/radiation updated during subcycling each KPRAD time step
		due to density and energy changes
ikprad_evolve_neutrals	0	Determines how KPRAD neutrals evolve spatially:
		0: Neither advect nor diffuse
		1: <b>RECOMMENDED</b> Advect and diffuse like other charge states
		2: Diffuse but do not advect
ikprad_min_option	1	Determine how KPRAD behaves below minimum den-
		sity/temperature (kprad_nemin, kprad_temin)
		1: No radiation/ionization/recombination (based on ne/te before
		subcycling)
		2: <b>RECOMMENDED</b> Recombination but no radiation/ionization
		(based on ne/Te during subcycling)
		3: No radiation/ionization/recombination (based on ne/Te dur-
		ing subcycling). NOTE: 1 and 3 behave the same if
		$ikprad_evolve_internal = 0$
iread_lp_source	0	Read impurity source from Lagrangian Particle code cloud.txt (UN-
		DER DEVELOPMENT)

# 10.23 Stellarator Geometry

Option	Default	Description			
type_ext_field	-1	1 External fields to be read in			
		0: Axisymmetric only. For RMP and error fields.			
		1: Free-boundary stellarator (itaylor=41) data.			
		2: Free-boundary stellarator (itaylor=41) data with			
		extsubtract=1.			
file_ext_field		(string) Vacuum/external field to be subtracted for itaylor=41.			
		Currently supported: FIELDLINES, MGRID.			
		Must start with 'fieldlines' or 'mgrid'.			
		Must have type_ext_field=2.			
file_total_field		(string) Total field for free-boundary stellarator (itaylor=41).			
		Currently supported: FIELDLINES, MGRID.			
		Must start with 'fieldlines' or 'mgrid'.			
		Must have type_ext_field=1,2.			
iread_vmec	0	1: read VMEC file to determine geometry. Must be 1 for stellarator.			
vmec_filename		(string) VMEC output .nc file			
bloat_factor	0	Free boundary only: Scale factor to bloat computational boundary			
		using input geometry.			
bloat_distance	0	Free boundary only: Distance to expand computational boundary			
		from input geometry.			
igeometry	0	1: must be set to use stellarator version			
nperiods	1	Number of field periods (stellarator geometry; VMEC nfp must be			
		divisible by nperiods). Note nplanes should be equal to at least $2 \text{ x}$			
		(number of toroidal modes per field period) x nperiods			
ifull_torus	0	0: Solve on one field period			
		1: Solve on full torus			
$nzer_factor$	-1	(integer) Scale factor for resolution of Zernike polynomial (used for			
		interpolation of VMEC)			
		-1: $n_{\text{zeri}} = 2 \text{ x mpol}$ (fixed boundary) and $1 \text{ x mpol}$ (free boundary).			
		Otherwise: $n_z = n_z er_factor \times mpol where mpol is poloidal reso-$			
		lution in VMEC			
nzer_manual	-1	(int) Resolution of Zernike polynomial (mainly for testing)			

# References

- D. A. Dunavant. High degree efficient symmetrical gaussian quadrature for the triange. Int'l J. for Num. Methods in Eng., 21:1129–1148, 1985.
- [2] S. C. Jardin. A triangular finite element with first-derivative continuity applied to fusion MHD applications. J. Comp. Phys., 200:133–152, 2004.

# A IDL Postprocessor

# A.1 Introduction

The IDL routines described here have been created for the purpose of reading, displaying, and manipulating data written by  $M3D-C^1$  to an HDF5 output file. These routines are contained within files stored in the subdirectory trunk/unstructured/idl/ in the M3D- $C^1$  SVN repository.

Invoking IDL

First, IDL must be given access to the postprocessor routines. This may be accomplished either by copying the **\*.pro** files in the repository to the working directory where IDL will be run, or to set the environment variable IDL\_PATH to include the directory where these files are located.

The IDL module must first be loaded to run IDL on portal.pppl.gov. This is done by

module load idl

IDL may then be invoked by

idl

Once IDL is running, the postprocessor routines must be compiled. This is done by

.run plot\_routines
.run power\_spectrum
.run read\_h5

The functions and procedures described in the following section are now ready to use.

Help In IDL

IDL has an on-line help system which may be invoked from the IDL command prompt by

?

Information about a specific intrinsic IDL routine, <routine>, may be obtained by

? <routine>

Note on function/procedure descriptions Functions and procedures in IDL are differentiated by whether a value is returned (as with functions) or not (as with procedures). Both functions and procedures may take "arguments" and "keywords" as command line parameters. In the following descriptions, the arguments of the function or procedure are shown with the command, and keywords are listed separately. Optional arguments are enclosed by square brackets. Keywords are always optional. (This notation differs from that in the on-line IDL help only in that keywords are not listed with the command.) Keywords are specified on the command line by

<keywordname>=<value>

Writing

/<keywordname>

has the same effect as

<keywordname>=1

and is therefore useful for boolean options. Some common examples are given below.

#### Common Examples

To display the field "psi" at time slice 1 of file "C1.h5", sampled on a regular  $100 \times 100$  grid, with an overlay of the LCFS and the mesh,

plot\_field, 'psi', 1, filename='C1.h5', /iso, /lcfs, /mesh, points=100

To plot the time series of the kinetic energy of file "C1.h5" in domain 0 < t < 100,

plot\_scalar, 'ke', filename='C1.h5', xrange=[0,100]

To plot the flux-averaged temperature profiles of files "1/C1.h5" and "2/C1.h5" at time slice 1, versus the normalized poloidal flux,

plot\_flux\_average, 'T', 1, filename=['1/C1.h5', '2/C1.h5'], /norm

### A.2 Functions/Procedures in plot\_routines.pro

## A.2.1 Procedure contour\_and\_legend

```
contour\_and\_legend, z [,x, y]
```

Description

This procedure draws a two-dimensional color contour plot of the data z, with horizontal and vertical coordinates x and y. nt frames are drawn.

Argumen	ts		
Name	I/O	Type	Description
$\overline{z}$	Ι	$\texttt{float}[\mathit{nt}, \mathit{nx}, \mathit{nz}]$	The values of the field to plot
x	Ι	float[nx]	The values of the $x$ -coordinate
y	Ι	float[nz]	The values of the $y$ -coordinate

### Keywords

Name	I/O	Type	Description
label	Ι	$\mathtt{string}[nt]$	The IDL-formatted label of the color
			bar for each frame
title	Ι	$\mathtt{string}[nt]$	The IDL-formatted title of the plot for
			each frame
range	Ι	float[2, nt]	An array of the ranges of $z$ to plot in
			each frame
nlevels	Ι	$\mathtt{int}[nt]$	The number of contour levels to plot
			for each frame
lines	Ι	$\mathtt{bool}[nt]$	Whether to draw contour lines for
			each frame
zlog	Ι	$\mathtt{bool}[nt]$	Whether to draw $z$ on a log scale, for
			each frame
jpeg	Ι	string	Write the resulting .jpeg image to the
			file <i>jpeg</i>
isotropic	Ι	bool	Use an isotropic aspect ratio when
			plotting
$color\_table$	Ι	int	Which intrinsic IDL color table to use
			when plotting

# A.3 Functions/Procedures in read\_h5.pro

## A.3.1 Function read\_field

 $field = read_field(name, r, z, t)$ 

### Description

This function reads the raw field data associated with *name* in the specified output file. The data is interpolated onto a uniform retangular grid, and returned in an array.

### Return Value

float[nt, points, points] field containing the value of the field at points  $\times$  points spatial sampling points at nt time slices.

### Arguments

Name	I/O	Type	e e	Description
name	Ι	stri	ng	The name of the field to read
r	0	floa	t[points]	The values of the $r$ -coordinate
z	0	floa	t[points]	The values of the $z$ -coordinate
t	Ο	floa	t[nt]	The values of the $t$ -coordinate
Keyword	ls			
Name		I/O	Type	Description
filenam	ne	Ι	string	The name of the HDF5 file to read
points		Ι	int	Number of sampling points per spatial dimension
rrange		Ι	float[2]	Range of $r$ -coordinate to read
zrange		Ι	float[2]	Range of $z$ -coordinate to read
$h_{-}symm$	netry	Ι	$\langle$ 1 $ $ -1 $\rangle$	Return only left-right $\langle$ symmetric $ $ anti-symmetric $\rangle$ part of field
v_symn	netry	Ι	$\langle$ 1 $ $ -1 $\rangle$	Return only up-down ( symmetric   anti-symmetric ) part of field
$di\!f\!f$		Ι	bool	If set, <i>name</i> should be an array with two filenames. The difference of the fields from the two files is re- turned.
slices		0	$\langle$ int $ $ int	$[2] \rangle$ $\langle$ Time slice   range of time slices $\rangle$ to read
mesh		Ο		

# A.3.2 Function flux\_average

 $result = \texttt{flux\_average}(field, slice)$ 

# Description

This function finds the flux-surface average of the field *field*.

# Return Value

<code>float[bins]</code> result contains the value of the flux-averaged field for a range of values of flux.

# Arguments

Name	I/O	Type	Description
field	Ι	$\langle \texttt{string}   \texttt{float}[1, points, points] \rangle$	If <i>field</i> is type <b>string</b> , then
			read field associated with
			name <i>field</i> . Otherwise, <i>field</i> is
			taken to contain the field data.
slice	Ι	int	If <i>field</i> is type <b>string</b> , then
			this is the time slice to read.
			Otherwise, <i>slice</i> is ignored.

Keywords

flux\_average takes all of the optional arguments for read\_field. In addition,

Name	I/O	Type	Description	
filename	Ι	string	The name of the HDF5 file to read	-
bins	Ι	int	Number of bins to subdivide the flux	
$psi^*$	I/O	float[1, points, points]	The flux field at the given time slice	
x	I/O	float[points]	<i>r</i> -coordinate values	
z	I/O	float[points]	z-coordinate values	* If $psi, x, z$ , and
t	I/O	float[points]	<i>t</i> -coordinate values	
flux	Ο	<pre>float[bins]</pre>	The value of the flux for each bin	
title	Ο	string	The IDL-formatted title of the field	
symbol	Ο	string	The IDL-formatted symbol of the field	
units	0	string	The IDL-formatted units of the field	
		. <u> </u>		

t are all provided as input, the flux\_average will not read the flux field itself.

## A.3.3 Function read\_scalar

 $result = read_field(name)$ 

# Description

This function reads the scalar quantity associated with name in the specified output file. The data is returned as an array .

# Return Value

float[nt] result contains the value of the scalar at each time step.

### Arguments

Name	I/O	Type	Description
name	Ι	string	The name of the scalar to read

Keywords

Name	I/O	Type	Description
filename	Ι	string	The name of the HDF5 file to read
time	Ο	float[nt]	The $t$ -coordinate of each element
title	Ο	string	The IDL-formatted title of the scalar
symbol	Ο	string	The IDL-formatted symbol of the scalar
units	Ο	string	The IDL-formatted units of the scalar

# A.3.4 Procedure plot\_field

plot\_field, name, slice, r, z, t

# Arguments

Name	I/O	Type	Description
name	Ι	string	The name of the field to read
slice	Ι	string	The time slice to read
r	Ο	float[points]	The values of the $r$ -coordinate
z	Ο	float[points]	The values of the $z$ -coordinate
t	Ο	float[nt]	The values of the $t$ -coordinate

#### Keywords

plot\_field takes all of the optional arguments for read\_field and contour\_and\_legend. In addition,

Name	I/O	Type	Description
xrange	Ι	float[2]	Range of $r$ -coordinate to plot
yrange	Ι	float[2]	Range of $z$ -coordinate to plot
lcfs	Ι	bool	Plot LCFS
mesh	Ι	bool	Plot mesh

# A.3.5 Procedure plot\_flux\_average

 $\verb"plot_flux_average", field, slice$ 

# Description

Plots the flux-surface average of a field as a function of poloidal flux. Data from multiple files or multiple times may be plotted at once.

#### Arguments

Name	I/O	Type	Description
field	Ι	$\langle \texttt{string}   \texttt{float}[1, points, points] \rangle$	If <i>field</i> is type <b>string</b> , then
slice	Ι	<pre>int[nt]</pre>	read field associated with name <i>field</i> . Otherwise, <i>field</i> is taken to contain the field data. If <i>field</i> is type <b>string</b> , then this is the time slice(s) to read.
			Otherwise, since is ignored.

Keywords			
Name	I/O	Type	Description
filename	Ι	string[nfiles]	Names of HDF5 file(s) to read
overplot	Ι	bool	Plot over previous plot
lcfs	Ι	bool	Plot LCFS
$minor\_radius$	Ι	bool	Plot against flux-average minor radius
$normalized\_flux$	Ι	bool	Plot against normalized flux
smooth	Ι	int	Boxcar-average final data over neigh-
			boring <i>smooth</i> bins

# A.3.6 Procedure plot\_timings

# plot\_timings

### Description

This function plots the relative time spent in various subroutines of M3D- $C^1$ . This data is only available if **itimer** = 1 was specified in the input file.

### Keywords

plot\_flux\_average takes all of the optional arguments for flux\_average. In addition,

Name	I/O	Type	Description
filename	Ι	string	The name of the HDF5 file to read
overplot	Ι	bool	Draws data over the

# A.3.7 Procedure write\_geqdsk

write\_geqdsk

Description This procedure writes equilibrium data do disk in the geqdsk format.

Arguments None.

Keywords

Name	I/O	Type	Description
eqfile	Ι	string	Filename to output geqdsk data
b0	Ι	float	Normalization of magnetic field, in Gauss
l0	Ι	float	Normalization of length scale, in cm

write\_geqdsk takes all of the optional arguments for read\_field and read\_parameter. In addition,

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	input1-391.vtk.pvtu	File	
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Figure 14: To visualize a mesh, select .pvtu file from Open File menu

# **B** Mesh Visualization

Paraview is program created by Kitware, Inc. which can visualize meshes and fields on meshes. It is the program of choice for viewing meshes created by the PUMI libraries. The API  $m3dc1\_mesh\_write$  writes a mesh either in "smb" or "vtk".

```
// filename: output file name
// option: 0 vtk file with field; 1 smb file
m3dc1_mesh_write(char* filename, int *option)
```

If the first argument is "output" and the second argument is 0, filename is output, it creates the files output.pvtu. Opening the output.pvtu file in Paraview will show users the mesh. Figures 14 and 15 illustrate "Open File" window and a mesh rendered in "Surface" mode by default.

Changing "Surface" to "Surface with Edges" will outline each visible element. Figures 16 and 17 illustrate how to make the decomposition visible.

Also, the mesh by default is rendered in one "Solid Color". There should be other options corresponding to the fields and numberings that were on this mesh at the time of file writing. There is an "apf\_part" alternative for files written by APF, which allows users to see the parallel partitioning of the mesh in color.

Mesh generation program provides "gface" alternative, which allows users to see the geometric face of the mesh in color. Figure 18 depicts a Paraview window with a mesh rendered with "gface".

Figure 19 illustrates the "Color Palette" tab of "Settings" menu which allows the users to change the colors of "Render View" such as edges, background, etc..



Figure 15: Initial Paraview window with a mesh



Figure 16: To render mesh decomposition, choose "Surface With Edge"



Figure 17: A mesh with "Surface With Edge"



Figure 18: A mesh rendered with "gface"

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Sear	rch (use Esc to clear text)	4
Colo	or used when solid coloring surfaces and faces.	
0	Surface	
Colo	or used for rendering elements like wireframes, points.	
0	Foreground	
Colo	or used for the edges when using 'Surface With Edges' representation.	
۲	Edges	
Colo	or used for text and other annotations.	
0	Text	
Sele	ection	
Cole	or used for showing selected cells/points.	
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Colo	or used for showing interactive selection.	
-	Interactive Selection	
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Cold Bac Cold	or used for showing interactive widget. Interactive Widget : <b>kground</b> Use Gradient Background: Enable gradient background. or used as background for the view. Background	

Figure 19: "Color Palette" tab of "Settings" menu"